A note on numerical solution of wireless power transmission with magnetic resonance

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Abstract

In the present article we argue that it is possible to find numerical solution with computer algebra package of coupled magnetic resonance equation for describing wireless energy transmission, as discussed recently by Karalis (2006) and Kurs et al. (2007). The proposed approach may be found useful in order to understand the phenomena related to magnetic resonance. Further observation is of course recommended in order to refute or verify this proposition.

Introduction

In recent years there are new interests in methods to transmit energy without wire. While it has been known for quite a long time that this method is possible theoretically (since Maxwell and Hertz), until recently only a few researchers consider this method seriously.

For instance, Karalis et al. [1] and also Kurs et al. [2] have presented such experiments and reported that efficiency of this method remains low. A plausible way to solve this problem is by better understanding of the mechanism of magnetic resonance. Meanwhile, Rangelov et al. [3] have presented an analogous scheme with respect to coupled-mode theory in quantum system.

In the present article we argue that it is possible to find numerical solution of coupled magnetic resonance equation for describing wireless energy transmission, as discussed recently by Karalis et al.(2006) and Kurs et al. (2007). The proposed approach may be found useful in order to understand the phenomena related to magnetic resonance.

Nonetheless, further observation is of course recommended in order to refute or verify this proposition.

Numerical solution of coupled-magnetic resonance equation

Recently, Kurs et al. [2] argue that it is possible to represent the physical system behind wireless energy transmission using coupled-mode theory, as follows:

$$a_m(t) = (i\omega_m - \Gamma_m)a_m(t) + \sum_{n \neq m} i\kappa_{nm}a_n(t) - F_m(t)$$
⁽¹⁾

The simplified version of equation (1) for the system of two resonance objects is given by Karalis et al. [1, p.2]:

$$\frac{da_1(t)}{dt} = -i(\omega_1 - i\Gamma_1)a_1(t) + i\kappa a_2(t),$$
⁽²⁾

$$\frac{da_2(t)}{dt} = -i(\omega_2 - i\Gamma_2)a_2(t) + i\kappa a_1(t),$$
(3)

(4)

These equations can be expressed as linear 1st order ODE as follows:

$$f'(t) = -i \alpha f(t) + i \kappa g(t),$$

And
$$g'(t) = -i\beta g(t) + i\kappa f(t),$$
 (5)

Where
$$\alpha = (\omega_1 - i\Gamma_1)$$
, (6)

And
$$\beta = (\omega_2 - i\Gamma_2).$$
 (7)

Numerical solution of these coupled-ODE equations can be found using Maxima release 5.23.2 as follows. First we find test when parameters (6) and (7) are set up to be 1. The solution is given by:

(%o1) 'diff(f(x),x)+%i*f(x)=%i*b*g(x); (%o2) 'diff(g(x),x)+%i*g(x)=%i*b*f(x); (%o3) desolve([%o1,%o2],[f(x),g(x)]);

The solutions for f(x) and g(x) are given by:

$$f(x) = e^{-it} \left[\frac{[2ig(0)b + 2if(0) - 2if(0)\sin(bx)]}{2b} + f(0)\cos(bx) \right]$$
(8)

$$g(x) = e^{-it} \left[\frac{[2ig(0) + 2if(0)b - 2ig(0)\sin(bx)]}{2b} + g(0)\cos(bx) \right]$$
(9)

Translated back to our equations (2) and (3), the solutions for $\alpha = \beta = 1$ are given by:

$$a_{1}(t) = e^{-it} \left[\frac{[2ia_{2}(0)\kappa + 2ia_{1}(0) - 2ia_{1}(0)\sin(\kappa t)]}{2\kappa} + a_{1}(0)\cos(\kappa t) \right]$$
(10)

$$a_{2}(t) = e^{-it} \left[\frac{[2ia_{1}(0)\kappa + 2ia_{2}(0) - 2ia_{2}(0)\sin(\kappa t)]}{2\kappa} + a_{2}(0)\cos(\kappa t) \right]$$
(11)

Now we will find numerical solution of equations (4) and (5) when $\alpha \neq \beta \neq 1$. Using Maxima 5.23.2, we find:

(%04) 'diff(f(x),x)+%i*a*f(x)=%i*b*g(x); (%05)'diff(g(x),x)+%i*c*g(x)=%i*b*f(x); (%06) desolve([%04,%05],[f(x),g(x)]);

And the solution is found to be quite complicated: these are formulae (12) and (13).

$$f(x) = e^{-(ic+ia)t/2} \begin{bmatrix} \frac{[2if(0)c + 2ig(0)b - f(0)(ic+ia)]\sin(\frac{t}{2}\sqrt{c^2 - 2ac + 4b^2 + a^2})}{\sqrt{c^2 - 2ac + 4b^2 + a^2}} + \\ + f(0)\cos(\frac{t}{2}\sqrt{c^2 - 2ac + 4b^2 + a^2}) \end{bmatrix}$$
(12)
$$f(x) = e^{-(ic+ia)t/2} \begin{bmatrix} \frac{[2if(0)b + 2ig(0)a - g(0)(ic+ia)]\sin(\frac{t}{2}\sqrt{c^2 - 2ac + 4b^2 + a^2})}{\sqrt{c^2 - 2ac + 4b^2 + a^2}} + \\ + g(0)\cos(\frac{t}{2}\sqrt{c^2 - 2ac + 4b^2 + a^2}) \end{bmatrix}$$
(13)

Translated back these results into our equations (2) and (3), the solutions are given by (14) and (15),

$$a_{1}(t) = e^{-(i\beta + i\alpha)t/2} \left| \frac{[2i\alpha_{1}(0)c + 2ia_{2}(0)b - a_{1}(0)(i\beta + i\alpha)]\sin(\frac{t}{2}\xi)}{\xi} + a_{1}(0)\cos(\frac{t}{2}\xi) \right|$$
(14)

$$a_{2}(t) = e^{-(i\beta + i\alpha)t/2} \left[\frac{\left[2i\alpha_{2}(0)c + 2ia_{1}(0)b - a_{2}(0)(i\beta + i\alpha)\right]\sin(\frac{t}{2}\xi)}{\xi} + a_{2}(0)\cos(\frac{t}{2}\xi) \right]$$
(15)

where we can define a new "ratio":

$$\xi = \sqrt{c^2 - 2ac + 4b^2 + a^2} \quad . \tag{16}$$

It is perhaps quite interesting to remark here that there is no "distance" effect in these equations, at least in theory. Nonetheless, further observation is of course recommended in order to refute or verify this proposition.

References:

[1] Karalis A., Joannopoulos J. D., and Soljacic M. (2006) Efficient wireless nonradiative midrange energy transfer. arXiv:physics/0611063.

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[3] Rangelov, A.A., H.Suchowski, Y. Silberberg, N.V.Vitanov (2010) Wireless adiabatic power transfer. arXiv:quant-ph/1011.0697.

[4] Maxima release 5.23.2 from http://maxima.sourceforge.net

[5] Smarandache F. & Christianto, V. (editors) (2007) Hadron models and related new energy issues. InfoLearnQuest, USA. ISBN: 978-1-59973-042-4.