

# Symmetry model E9CS

by Mag.rer.nat. Kronberger Reinhard  
 Email: support@kro4pro.com  
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## Motivation 1:

Why do we consider the E9 group (more specifically the Coxeter element of this group)?

- 1) E9 is an affine group and thus has something to do with extension.
- 2) The extension is flat as the universe.
- 3) The key Coxeter element of the group produces symmetries involving our current standard model.

The fundamentals here:

- [https://en.wikipedia.org/wiki/Coxeter\\_group](https://en.wikipedia.org/wiki/Coxeter_group)
- <https://de.wikipedia.org/wiki/Wurzelsystem>
- <http://home.mathematik.uni-freiburg.de/soergel/Skripten/XXSPIEG.pdf>

## Symmetries which arise from the Coxeter element of the E9.

$$\mathbf{E9CS} = \mathbf{SU(5)} \times \mathbf{SU(3)} \times \mathbf{SU(2)} \times \mathbf{U(1)} \times \mathbf{U(1)} \quad (\text{SU}(n) = \text{Special unitary group, U}(1) \text{ unitary group})$$

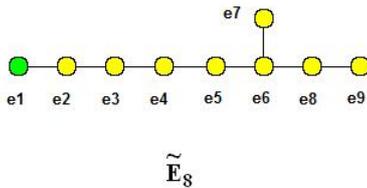
Pronounced E9Coxeter-Symmetry

$$\text{evidently } \mathbf{SU(5)} \times \mathbf{SU(3)} \times \mathbf{SU(2)} \times \mathbf{U(1)} \times \mathbf{U(1)} \supset \mathbf{SU(3)}_c \times \mathbf{SU(2)}_L \times \mathbf{U(1)}_Y \quad (\text{Color charge, isospin, Hyper charge})$$

Write the symmetry in order to:

$$\mathbf{E9CS} = \mathbf{SU(5)}_s \times \mathbf{U(1)}_{Y2} \times \mathbf{U(1)}_{Y1} \times \mathbf{SU(2)}_L \times \mathbf{SU(3)}_c \quad (= \text{Expansion} \times \text{actual Standard Model})$$

Dynkin Diagram E9 (affine one point extension of group E8):



## Derivative of the symmetries of E9CS from the invariants of the Coxeter elements E9:

The Coxeter element is the product of the generating reflections of E9.

$$\mathbf{Coxeterelement} = e1.e2.e3.e4.e5.e6.e7.e8.e9$$

The Coxeterpolynomial is the characteristic polynomial of Coxeter elements and has the form:

$$E_9(x) = \frac{x^5 - 1}{x - 1} \cdot \frac{x^3 - 1}{x - 1} \cdot \frac{x^2 - 1}{x - 1} \cdot (x - 1)^2$$

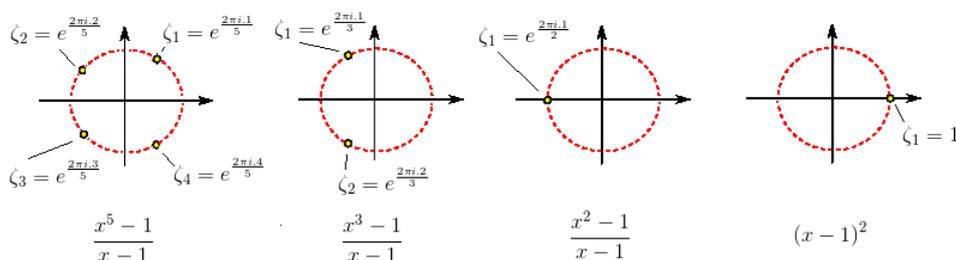
$$E9CS = SU(5) \times SU(3) \times SU(2) \times U(1)^2$$

$E_9(x)$  ... characteristic polynomial of the coxeterelement of E9

$E_9(x)$  is a polynomial with terms of cyclotomic factors  $Z_n = \frac{x^n - 1}{x - 1}$  for  $n > 1$  and  $(x - 1)$  for  $n = 1$ .  
 The cyclotomic factors are the characteristic polynomial of the  $A_{n-1}$  (which is the Dynkin diagram for the  $SU(n)$  Liegroup. See more here: [https://en.wikipedia.org/wiki/Special\\_unitary\\_group](https://en.wikipedia.org/wiki/Special_unitary_group)).

So finally the symmetry space of the Coxeterelement is  $\mathbf{SU(5)} \times \mathbf{SU(3)} \times \mathbf{SU(2)} \times \mathbf{U(1)} \times \mathbf{U(1)}$

## Eigenvalues of the Coxeterpolynomial



Eigenspace of the Coxeterpolynomial

$$\mathbb{C}^4 \times \mathbb{C}^2 \times \mathbb{C}^1 \times \mathbb{C}^2$$

## Motivation 2:

What bring us the additional symmetries?

- (1) These have the potential to describe new particles.
- (2) These have the potential to describe the space and time.
- (3) These have the potential to describe gravity.

Wish to analogously represent Graviton to the photon as a blend (Weinberg angle see <8>).

## <1> The Idea

Light and gravitation just like photon and graviton have something in common.  
Both are massless and propagate with the speed of light.

We know that light by the symmetry breaking 1:  $SU(2) \times U(1) \rightarrow U(1)$  is described as a mixture.  
So light is a part of the **electro-weak interactions**.

we consider analog gravity as a result of a further symmetry breaking

**Symmetry breaking 2:  $SU(5) \times U(1) \times U(1) \rightarrow U(1)$**

Our extended standard model allows us this.

We will now like to assign our relevant  $SU(n)$ 's to algebras division (real numbers, complex numbers, ...).

$$SU(1) \leftrightarrow \mathbb{R}$$

$$SU(2) \leftrightarrow \mathbb{C}$$

$$SU(3) \leftrightarrow \mathbb{H}$$

$$SU(5) \leftrightarrow \mathbb{O}$$

This 4 division algebras (real numbers, complex numbers, quaternions and octonions) develop through the doubling process

see more at <https://de.wikipedia.org/wiki/Verdopplungsverfahren>

Considering the dimensions of the  $SU(2) = 1, SU(3) = 2, SU(5) = 4$  then this is double as well.

There appears to be a connection between the division algebras and the  $SU(n)$ 's ( $n = 2, 3, 5$ ) which I hope is known in analytic geometry or another area.  
I assume this connection warrants as simply as given.

Notes but no clear allocation can be found in this direction at Corinne A. Manogue and Tevian Dray, John Baez, etc.

Therefore, we rely analogously on the **Higgsfield** (2 x complex = doublet)

$$\phi = \begin{bmatrix} \phi^+ \\ \phi^0 \end{bmatrix} = \begin{bmatrix} \phi_1^+ + i.\phi_2^+ \\ \phi_1^0 + i.\phi_2^0 \end{bmatrix}$$

## <2> the Oktoquintenfield (5 x Oktonions= Quintett).

$$\phi = \begin{bmatrix} \phi^G \\ \phi^R \\ \phi^F \\ \phi^S \\ \phi^O \end{bmatrix} = \begin{bmatrix} \phi_0^G + i_1.\phi_1^G + i_2.\phi_2^G + i_3.\phi_3^G + i_4.\phi_4^G + i_5.\phi_5^G + i_6.\phi_6^G + i_7.\phi_7^G \\ \phi_0^R + i_1.\phi_1^R + i_2.\phi_2^R + i_3.\phi_3^R + i_4.\phi_4^R + i_5.\phi_5^R + i_6.\phi_6^R + i_7.\phi_7^R \\ \phi_0^F + i_1.\phi_1^F + i_2.\phi_2^F + i_3.\phi_3^F + i_4.\phi_4^F + i_5.\phi_5^F + i_6.\phi_6^F + i_7.\phi_7^F \\ \phi_0^S + i_1.\phi_1^S + i_2.\phi_2^S + i_3.\phi_3^S + i_4.\phi_4^S + i_5.\phi_5^S + i_6.\phi_6^S + i_7.\phi_7^S \\ \phi_0^O + i_1.\phi_1^O + i_2.\phi_2^O + i_3.\phi_3^O + i_4.\phi_4^O + i_5.\phi_5^O + i_6.\phi_6^O + i_7.\phi_7^O \end{bmatrix}$$

or written otherwise so that the equivalence to the Higgs field is clear (where  $i_4$  is pulled from)

$$\phi = \begin{bmatrix} \phi^G \\ \phi^R \\ \phi^F \\ \phi^S \\ \phi^O \end{bmatrix} = \begin{bmatrix} \phi_0^G + i_1.\phi_1^G + i_2.\phi_2^G + i_3.\phi_3^G + i_4.(\phi_4^G + i_1.\phi_5^G + i_2.\phi_6^G + i_3.\phi_7^G) \\ \phi_0^R + i_1.\phi_1^R + i_2.\phi_2^R + i_3.\phi_3^R + i_4.(\phi_4^R + i_1.\phi_5^R + i_2.\phi_6^R + i_3.\phi_7^R) \\ \phi_0^F + i_1.\phi_1^F + i_2.\phi_2^F + i_3.\phi_3^F + i_4.(\phi_4^F + i_1.\phi_5^F + i_2.\phi_6^F + i_3.\phi_7^F) \\ \phi_0^S + i_1.\phi_1^S + i_2.\phi_2^S + i_3.\phi_3^S + i_4.(\phi_4^S + i_1.\phi_5^S + i_2.\phi_6^S + i_3.\phi_7^S) \\ \phi_0^O + i_1.\phi_1^O + i_2.\phi_2^O + i_3.\phi_3^O + i_4.(\phi_4^O + i_1.\phi_5^O + i_2.\phi_6^O + i_3.\phi_7^O) \end{bmatrix}$$

**Comparison  
Oktoquintenfield &  
Higgsfield**

$$\phi = \begin{bmatrix} \phi^+ \\ \phi^0 \end{bmatrix} = \begin{bmatrix} \phi_1^+ + i.\phi_2^+ \\ \phi_1^0 + i.\phi_2^0 \end{bmatrix}$$

This provides 40 degrees of freedom.

24 of which will be "spent" for our  $SU(5)$  tensor bosons for the 5th longitudinal spin degree of freedom (24 Goldstone bosons swallowed over gauge transformation) thus remain 16 left.

The S, F, R, G and H charges are the 5 charges of the  $SU(5)$  analogous to the 3 color charges of  $SU(3)$  and the 2 charges (+ -) of  $SU(2)$ .

The letters stand for S = See, F = feeling, R=smelling G = Taste and H = Hear

Calling therefore the charges of the  $SU(5)$  sense charges.

Note: These charges have (such as the color charges of quarks with color) nothing to do with the senses, but to give a name to the child for reference only.

We now want to look at the 16 ( $40 - 24 = 16$ ) remaining degrees of freedom.

Make the following division for the 40 field components of the Oktoquinten field as a physical approach:

Take care that the division is not unique because for the left half 4 gray fields we can use  $4.3.2.1 = 24$  Permutations of them in the orange area.  
And for the left 4 charges we have five over  $4 = 5$  Permutations.

So at all we have  $5 \times 24 = 120$  possible permutations.

On the Higgsfield we have  $2 \times 1 = 2$  permutations.

$\Phi G$	1	i1	i2	i3	i4	i5	i6	i7	
$\Phi R$	1	i1	i2	i3	i4	i5	i6	i7	
$\Phi F$	1	i1	i2	i3	i4	i5	i6	i7	
$\Phi S$	1	i1	i2	i3	i4	i5	i6	i7	
$\Phi O$	1	i1	i2	i3	i4	i5	i6	i7	Charge 0

	Energy-Time
	Momentum-Space
	Shearstress/Energyflux
	for the 4 neutral Bosons
	for the 20 charged Bosons

Analogous to the Higgspotential we declare a Potential on the Oktoquintenfeld

### <3> Potential over the Oktoquintenfeld

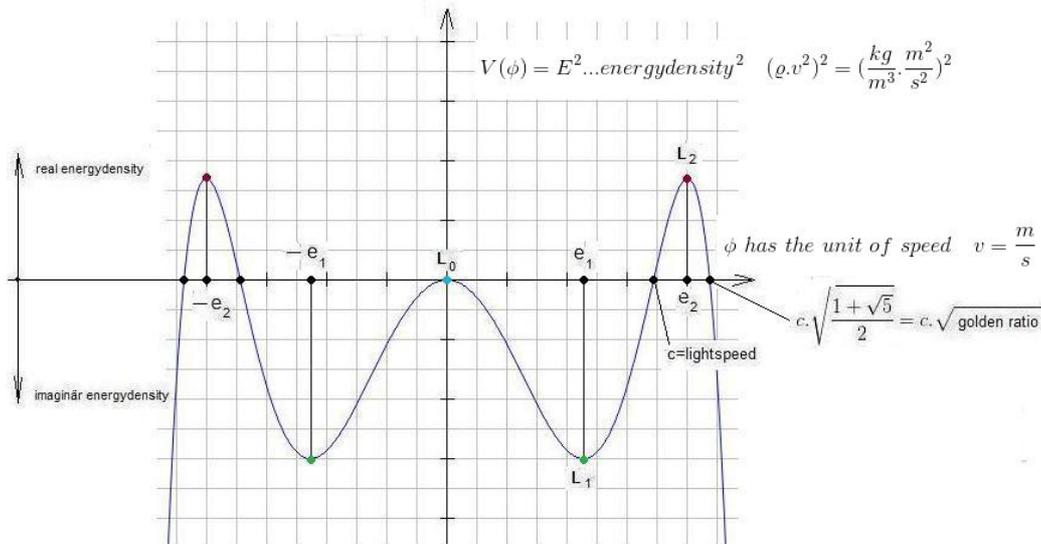
$$V(\phi) = \frac{\gamma^2}{2} |\phi|^2 + \frac{\mu^2}{4} |\phi|^4 + \frac{\lambda^2}{8} |\phi|^8 \quad \text{with } \phi \in \mathbb{O}^5$$

$\gamma, \lambda \in i.\mathbb{R}$  (imaginaer) and  $\mu \in \mathbb{R}$

$$\frac{\gamma^2}{2} \dots \text{momentumdensity}^2 \quad (\rho.v)^2 = \left(\frac{kg}{m^3} \cdot \frac{m}{s}\right)^2$$

$$\frac{\mu^2}{4} \dots \text{massdensity}^2 \quad (\rho)^2 = \left(\frac{kg}{m^3}\right)^2$$

$$\frac{\lambda^2}{8} \dots \text{spindensity}^2 \quad \left(\frac{\rho}{v^2}\right)^2 = \left(\frac{kg}{\frac{m^2}{s^2}}\right)^2$$



The coefficients of the potential comes from selfinteractions.

Therefore we make the assumption that we have the following relation :

$$C := \frac{-\mu^2}{\lambda^2} = \left(4 \cdot \frac{\gamma^2}{\mu^2}\right)^2 \quad \gamma^2, \lambda^2 < 0 \quad \mu^2 > 0$$

Then it follows by exact calculation that

$$C = c^4 \cdot \varphi^2$$

$c$ ...speed of light

$\varphi$ ...golden ratio = 1,6180...

The first mixingangle which comes from the minimum of the Oktoquintenpotential is appr. equal to the

WEINBERG - ANGLE  $\approx 28,89^\circ$

see <8>

For aesthetic reasons we want keep in mind that for the coming formulars  $\phi = \text{absolut}(\phi)$ .

We want that the second part of the Oktoqintenpotential is our quadratic vacuumenergydensity.

$$\frac{\mu^2}{4} |c|^4 = \frac{1}{4} \left( \frac{\Lambda \cdot c^4}{8\pi G} \right)^2 = \frac{1}{4} (\rho_{\text{vacuum}} \cdot c^2)^2$$

$\Lambda$ ...cosmological constant

$\rho_{\text{vacuum}}$ ...vacuum massdensity

then with the relation  $c^4 \cdot \varphi^2 = \frac{-\mu^2}{\lambda^2} = \left(4 \cdot \frac{\gamma^2}{\mu^2}\right)^2$  we get the potential as

EINSTEIN – FORM

$$V(\phi) = \left( \frac{\Lambda \cdot c^4}{8\pi G} \right)^2 \cdot \frac{1}{8 \cdot \varphi^2} \cdot \left( -\varphi^3 \cdot \left(\frac{\phi}{c}\right)^2 + 2 \cdot \varphi^2 \cdot \left(\frac{\phi}{c}\right)^4 - \left(\frac{\phi}{c}\right)^8 \right)$$

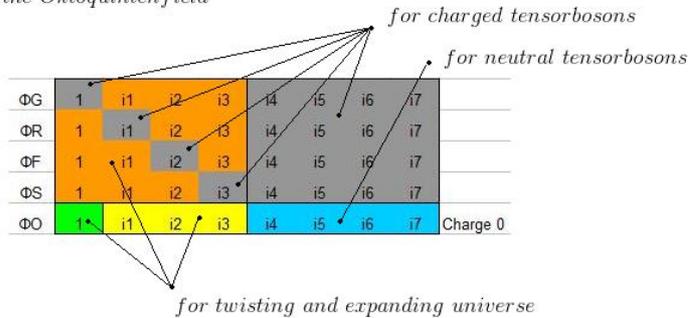
## <4> Lagrangedensity of the Oktoqintenfield/Oktoqintenpotential

Hint:

I do the same steps as shown in this cooking recipe for the Higgsfield.

[https://www.lsw.uni-heidelberg.de/users/mcamenzi/HD\\_Higgs.pdf](https://www.lsw.uni-heidelberg.de/users/mcamenzi/HD_Higgs.pdf)

splitting of the Oktoqintenfield



Analogous to the **electroweak** theory, we want to talk about a gravito super weak theory here. The **electroweak** theory brings together the electrical with the weak interactions. The **gravito-sensecharge** theory brings together the gravitational with the dark interactions.

Similar to the **SU(3) Vectorbosons** which are named **Gluons**

we name our **SU(5) Tensorbosons Repelions**.

The force between different charged Repelions is repulsive because they are tensorbosons (2nd – order).

Similar to the Higgsfield we assign our Repelions to the Oktoqintenfield by the following scheme.

The numbers are the sense charges (see < 2 >).

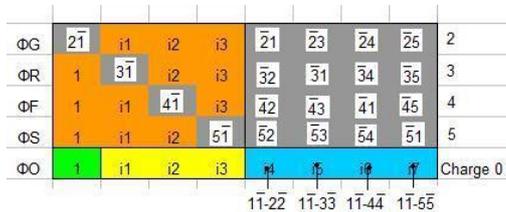
5 = See

4 = Feeling

3 = Smelling

2 = Taste

1 = Hear



similar to the the higgsfield where the vacuumexpectation is

$$\phi_{\text{vac}} = \begin{pmatrix} 0 \\ v \end{pmatrix}$$

the vacuumexpectation of the Oktoqintenfield is (green, yellow, orange)

$$\phi_{\text{vac}} = v \cdot \begin{pmatrix} 0 + i_1 + i_2 + i_3 \\ 1 + 0 + i_2 + i_3 \\ 1 + i_1 + 0 + i_3 \\ 1 + i_1 + i_2 + 0 \\ 1 + i_1 + i_2 + i_3 \end{pmatrix}$$

where  $v = e_1$  ist the minimum of the Oktoqintenpotential and  $i_1, i_2$  and  $i_3$  the imaginaer quaternions.

As mentioned in < 6 > point 3) we assume that the left 4 charged bosons decomposed to electrons, neutrinos and quarks and couple to the higgs field instead. Then the vacuum expectation changes to :

$$\phi_{vac} = v \cdot \begin{pmatrix} 1 + i_1 + i_2 + i_3 \\ 1 + i_1 + i_2 + i_3 \end{pmatrix}$$

### STEP 1: Lorentzinvariant Lagrangedensity for the Oktoquintenfield

$$\mathcal{L}_\phi = (D^\mu \phi)^\dagger (D_\mu \phi) - V(\phi^\dagger \phi)$$

with  $\phi \in \mathbb{O}^5$  Octonions<sup>5</sup>

The potential  $V$  is shown in < 3 >

$T_{ij}$	$i \rightarrow$	Generators of the $SU(5)$				
$\downarrow$	$j$	$\begin{pmatrix} 0 & 1 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \end{pmatrix}$	$\begin{pmatrix} 0 & -i & 0 & 0 & 0 \\ i & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \end{pmatrix}$	$\begin{pmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \end{pmatrix}$	$\begin{pmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \end{pmatrix}$	$\begin{pmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 \end{pmatrix}$
		$\begin{pmatrix} 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \end{pmatrix}$	$\begin{pmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & -i & 0 & 0 \\ 0 & i & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \end{pmatrix}$	$\begin{pmatrix} 0 & 0 & -i & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ i & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \end{pmatrix}$	$\begin{pmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \end{pmatrix}$	$\begin{pmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \end{pmatrix}$
		$\begin{pmatrix} 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \end{pmatrix}$	$\begin{pmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -i & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & i & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \end{pmatrix}$	$\begin{pmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -i & 0 \\ 0 & 0 & i & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \end{pmatrix}$	$\begin{pmatrix} 0 & 0 & 0 & -i & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ i & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \end{pmatrix}$	$\begin{pmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 \\ i & 0 & 0 & 0 & 0 \end{pmatrix}$
		$\begin{pmatrix} 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 \end{pmatrix}$	$\begin{pmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & -i \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & i & 0 & 0 & 0 \end{pmatrix}$	$\begin{pmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -i & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & i & 0 & 0 \end{pmatrix}$	$\begin{pmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -i & 0 \\ 0 & 0 & 0 & i & 0 \end{pmatrix}$	$\begin{pmatrix} 0 & 0 & 0 & 0 & -i \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ i & 0 & 0 & 0 & 0 \end{pmatrix}$
		$\begin{pmatrix} 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & -1 \end{pmatrix}$	$\begin{pmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & -1 \end{pmatrix}$	$\begin{pmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & -1 \end{pmatrix}$	$\begin{pmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & -1 \end{pmatrix}$	

$W - Boson$  scheme

$$\begin{pmatrix} W^{11} & W^{12} & W^{13} & W^{14} & W^{15} \\ W^{21} & W^{22} & W^{23} & W^{24} & W^{25} \\ W^{31} & W^{32} & W^{33} & W^{34} & W^{35} \\ W^{41} & W^{42} & W^{43} & W^{44} & W^{45} \\ W^{51} & W^{52} & W^{53} & W^{54} & \end{pmatrix}$$

hint :  $W^{ij} = W_\mu^{ij}$

We take a look on the symmetry

$$SU(5) \times U(1) \times U(1)$$

$$W^{ij} \quad B^0 \quad B^1$$

calculate covariant derivation

$$D_\mu \phi = (\partial_\mu + \frac{i.g}{2} \cdot \tau_{ij} \cdot W_\mu^{ij} + \frac{i.g'}{2} \cdot Id^0 B_\mu^0 + \frac{i.g''}{2} \cdot Id^1 B_\mu^1) \cdot \phi$$

$$\tau_{ij} \cdot W_\mu^{ij} = \begin{pmatrix} W^{51} & W^{11} - i.W^{12} & W^{21} - i.W^{23} & W^{31} - i.W^{34} & W^{41} - i.W^{45} \\ W^{11} + i.W^{12} & W^{52} & W^{22} - i.W^{13} & W^{14} - i.W^{32} & W^{15} - i.W^{42} \\ W^{21} + i.W^{23} & W^{22} + i.W^{13} & W^{53} & W^{24} - i.W^{33} & W^{25} - i.W^{43} \\ W^{31} + i.W^{34} & W^{14} + i.W^{32} & W^{24} + i.W^{33} & W^{54} & W^{35} - i.W^{44} \\ W^{41} + i.W^{45} & W^{15} + i.W^{42} & W^{25} + i.W^{43} & W^{35} + i.W^{44} & -(W^{51} + W^{52} + W^{53} + W^{54}) \end{pmatrix}$$

$$Id^0 = Id^1 = \begin{pmatrix} 1 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 1 \end{pmatrix}$$

and for example

$$W_{12} = \frac{W^{11} - i.W^{12}}{\sqrt{2}}$$

The boson which changes the charge from 1 (hear) to 2 (taste).

Then

$$D_\mu \phi_{vac} = \frac{\phi_{vac} \cdot i}{2} \cdot g \cdot \begin{pmatrix} W^{51} & \sqrt{2} W_{12} & \sqrt{2} W_{13} & \sqrt{2} W_{14} & \sqrt{2} W_{15} \\ \sqrt{2} W_{21} & W^{52} & \sqrt{2} W_{23} & \sqrt{2} W_{24} & \sqrt{2} W_{25} \\ \sqrt{2} W_{31} & \sqrt{2} W_{32} & W^{53} & \sqrt{2} W_{34} & \sqrt{2} W_{35} \\ \sqrt{2} W_{41} & \sqrt{2} W_{42} & \sqrt{2} W_{43} & W^{54} & \sqrt{2} W_{45} \\ \sqrt{2} W_{51} & \sqrt{2} W_{52} & \sqrt{2} W_{53} & \sqrt{2} W_{54} & -(W^{51} + W^{52} + W^{53} + W^{54}) \end{pmatrix} + \begin{pmatrix} g' \cdot B^0 + g'' \cdot B^1 & 0 & 0 & 0 & 0 \\ 0 & g' \cdot B^0 + g'' \cdot B^1 & 0 & 0 & 0 \\ 0 & 0 & g' \cdot B^0 + g'' \cdot B^1 & 0 & 0 \\ 0 & 0 & 0 & g' \cdot B^0 + g'' \cdot B^1 & 0 \\ 0 & 0 & 0 & 0 & g' \cdot B^0 + g'' \cdot B^1 \end{pmatrix} \cdot \begin{pmatrix} 1 + i_1 + i_2 + i_3 \\ 1 + i_1 + i_2 + i_3 \end{pmatrix}$$

then

$$(D^\mu \phi_{vac})^\dagger (D_\mu \phi_{vac}) = \frac{\phi_{vac}^2}{4} \cdot [ 4 \cdot (gW^{51} + g'B^0 + g''B^1)^2 + 8g^2(W_{12}W_{21} + W_{13}W_{31} + W_{14}W_{41} + W_{15}W_{51}) + \text{something} + 4 \cdot (gW^{52} + g'B^0 + g''B^1)^2 + 8g^2(W_{12}W_{21} + W_{23}W_{32} + W_{24}W_{42} + W_{25}W_{52}) + \text{something} + 4 \cdot (gW^{53} + g'B^0 + g''B^1)^2 + 8g^2(W_{13}W_{31} + W_{23}W_{32} + W_{34}W_{43} + W_{35}W_{53}) + \text{something} + 4 \cdot (gW^{54} + g'B^0 + g''B^1)^2 + 8g^2(W_{14}W_{41} + W_{24}W_{42} + W_{34}W_{43} + W_{45}W_{54}) + \text{something} + 4 \cdot (-g(W^{51} + W^{52} + W^{53} + W^{54}) + g'B^0 + g''B^1)^2 + 8g^2(W_{15}W_{51} + W_{25}W_{52} + W_{35}W_{53} + W_{45}W_{54}) + \text{something} ]$$

hint :  $W^{ij} = W_\mu^{ij}$  and  $B^0 = B_\mu^0$  and  $B^1 = B_\mu^1$

like the result of the Higgsfield we expect something like that :

$$(D^\mu \phi_{vac})^\dagger (D_\mu \phi_{vac}) = \frac{v^2}{8} \cdot (g^2 \cdot (W^+)^2 + g^2 \cdot (W^-)^2 + (g^2 + g'^2) \cdot Z_\mu \cdot Z^\mu + 0 \cdot A_\mu \cdot A^\mu)$$

We have a lot of summands so we first want to take a look on the diagonal elements of the covariant derivation.

In the Higgsfieldtheory we get as result the massive Z - Bosons and the Photon as a mixing of neutral W and B bosons.

We calculate the expression which is a symmetric bilinear form :

Momentumdensity - Matrix

$$(W^{51} \ W^{52} \ W^{53} \ W^{54} \ B^0 \ B^1) \cdot \begin{pmatrix} 8g^2 & 4g^2 & 4g^2 & 4g^2 & 0 & 0 \\ 4g^2 & 8g^2 & 4g^2 & 4g^2 & 0 & 0 \\ 4g^2 & 4g^2 & 8g^2 & 4g^2 & 0 & 0 \\ 4g^2 & 4g^2 & 4g^2 & 8g^2 & 0 & 0 \\ 0 & 0 & 0 & 0 & 20g'^2 & -20g'g'' \\ 0 & 0 & 0 & 0 & -20g'g'' & 20g''^2 \end{pmatrix} \cdot \begin{pmatrix} W^{51} \\ W^{52} \\ W^{53} \\ W^{54} \\ B^0 \\ B^1 \end{pmatrix}$$

linearly independent      linearly dependent

and compare it with the red area of the dynamic lagrange part.

Someone can easy proof that is identical.

Then with diagonalizing the Momentumdensity - Matrix we get the following result :

$$(D^\mu \phi_{vac})^\dagger (D_\mu \phi_{vac}) = \frac{\phi_{vac}^2}{4} \cdot [\text{green area} + \text{red area} + \text{something}] = 0$$

$$(D^\mu \phi_{vac})^\dagger (D_\mu \phi_{vac}) = \frac{\phi_{vac}^2}{4} \cdot [ 8g^2 \cdot (|W_{12}|^2 + |W_{21}|^2 + |W_{13}|^2 + |W_{31}|^2 + |W_{14}|^2 + |W_{41}|^2 + |W_{23}|^2 + |W_{32}|^2 + |W_{24}|^2 + |W_{42}|^2 + |W_{34}|^2 + |W_{43}|^2 + |W_{15}|^2 + |W_{51}|^2 + |W_{35}|^2 + |W_{53}|^2 + |W_{25}|^2 + |W_{52}|^2 + |W_{45}|^2 + |W_{54}|^2) + 8g^2 \cdot (|Z^0|^2 + |Z^1|^2 + |Z^2|^2 + |Z^3|^2) + 20 \cdot (g'^2 + g''^2) \cdot |\Gamma|^2 + 0 \cdot |G|^2 ]$$

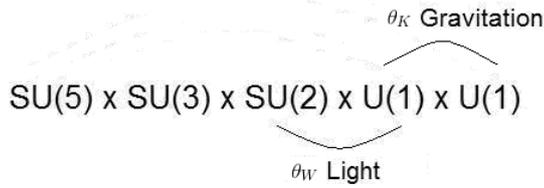
U(1) Gravitation field

eigenvalues of the Momentumdensity matrix

Momentumdensity Quadrats

The coupling angles  $\alpha_1$  and  $\alpha_2$  comes from the extremal values of the Oktoquintenpotential.

As calculated in < 8 > we have :



Mixingangles

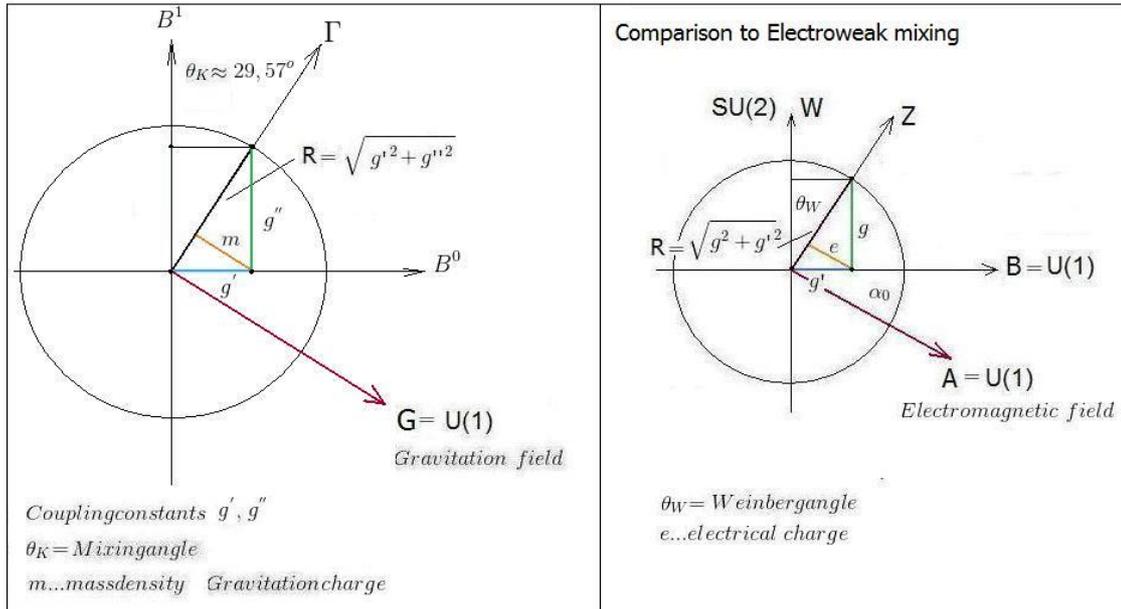
$$\theta_W \approx \alpha_1 \quad (\theta_W = \text{Weinbergangle})$$

$$\theta_K \approx 29,57^\circ$$

$$\alpha_1 \approx 28,89^\circ$$

$$\alpha_2 \approx 121,60^\circ$$

$$\begin{cases} g' = R \cdot \sin(\alpha_1) \\ g'' = R \cdot \sin(\alpha_2) \end{cases}$$



Then the Graviton and the  $\Gamma$  - Boson is a mixing :

$$\begin{pmatrix} \Gamma_\mu \\ G_\mu \end{pmatrix} = \begin{pmatrix} \cos(\theta_K) & -\sin(\theta_K) \\ \sin(\theta_K) & \cos(\theta_K) \end{pmatrix} \cdot \begin{pmatrix} B_\mu^0 \\ B_\mu^1 \end{pmatrix}$$

The Impulsdensity and therefore the massdensity and therefore the mass of the W and Z Bosons are equal because they have the same couplingconstant g.

The relation of the  $\Gamma$  particlemass to the W - Bosonmass is :

$$\frac{M_\Gamma}{M_W} = \frac{20 \cdot (g'^2 + g''^2)}{8g^2} = \frac{5 \cdot (g'^2 + g''^2)}{2g^2}$$

We developpe around the vacuumexpectation  $\phi_{min}$  to see the interactionterms.

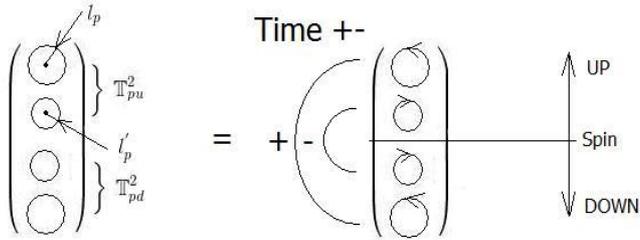
$$\phi = \frac{\phi_{min}}{c} \cdot (c + H) \cdot \begin{pmatrix} 1 + i_1 + i_2 + i_3 \\ 1 + i_1 + i_2 + i_3 \end{pmatrix}$$



Motivated by our results in < 6 > we are thinking about two such Planck – Tori

$\mathbb{T}_{pu}^2$  and  $\mathbb{T}_{pd}^2$

Visual



To generate an 2 – dimensional object with spin = 2 we have to connect the coordinatecycles like in < 6.4 > .

schematic in our second curvature tensor for the double Cliffordtorus

$$W_{Boson} = \left( \begin{array}{c|c|c} \text{Shear - Components} \\ \hline \text{[Diagram of 3x3 grid of circles with colored squares]} \\ \hline \end{array} \right) \subset S_{l_p \cdot \sqrt{\varphi}}^3$$

This construction has 3 different Tori  $\subset S_{l_p \cdot \sqrt{\varphi}}^3$

$\mathbb{T}_{pu}^2$	$\mathbb{T}_{pd}^2$	$\mathbb{T}_{ps}^2$
$\left( \begin{array}{c} \text{[Red square in circle]} \\ \text{[Red square in circle]} \end{array} \right) \quad r = l_p$	$\left( \begin{array}{c} \text{[Blue square in circle]} \\ \text{[Blue square in circle]} \end{array} \right) \quad r = \frac{l_p}{\sqrt{\varphi}}$	$\left( \begin{array}{c} \text{[Green square in circle]} \\ \text{[Green square in circle]} \end{array} \right) \quad r = l_p \cdot \sqrt{\frac{\varphi}{2}}$
$\left( \begin{array}{c} \text{[Red square in circle]} \\ \text{[Red square in circle]} \end{array} \right) \quad r = \frac{l_p}{\sqrt{\varphi}}$	$\left( \begin{array}{c} \text{[Blue square in circle]} \\ \text{[Blue square in circle]} \end{array} \right) \quad r = l_p$	$\left( \begin{array}{c} \text{[Green square in circle]} \\ \text{[Green square in circle]} \end{array} \right) \quad r = l_p \cdot \sqrt{\frac{\varphi}{2}}$

The  $W_{Boson}$  has Planckmass then it follows that the frequency is:

$$\omega_p = \frac{1}{t_p} = \sqrt{\frac{c^5}{h \cdot G}}$$

resting  $W_{Boson}$ :

$$W_{Boson} = \left( \begin{array}{c|c|c|c} l_p \cdot e^{-i\omega_p t} & 0 & 0 & l_p \cdot \sqrt{\frac{\varphi}{2}} \cdot e^{-i\omega_p t} \\ \hline 0 & \frac{l_p}{\sqrt{\varphi}} \cdot e^{i\omega_p t} & l_p \cdot \sqrt{\frac{\varphi}{2}} \cdot e^{-i\omega_p t} & 0 \\ \hline 0 & l_p \cdot \sqrt{\frac{\varphi}{2}} \cdot e^{-i\omega_p t} & \frac{l_p}{\sqrt{\varphi}} \cdot e^{i\omega_p t} & 0 \\ \hline l_p \cdot \sqrt{\frac{\varphi}{2}} \cdot e^{-i\omega_p t} & 0 & 0 & l_p \cdot e^{-i\omega_p t} \end{array} \right)$$

Properties of  $W_{Boson}$

- 1) is a Tensorboson
- 2) lies in a 3 – Sphere
- 3) is build by 3 Clifford Tori
- 4) Is a flat two dimensional object

We set  $l_p = 1$  for easier handling.

The class of all Clifford – Tori in  $S^3_{\sqrt{\varphi}}$  then is

$$C_a := \left\{ (z_1, z_2) \in \mathbb{C}^2 : |z_1| = \sqrt{\varphi} \cdot \sqrt{\frac{1+a}{2}} \quad |z_2| = \sqrt{\varphi} \cdot \sqrt{\frac{1-a}{2}} \right\}$$

for a real parameter  $a \in (-1, 1)$ . This is an embedded surface in  $S^3_{\sqrt{\varphi}}$  with mean curvature constant equal to

$$H_a := \frac{2a}{\sqrt{\varphi} \cdot \sqrt{1-a^2}}$$

Then our three Clifford – Tori  $\mathbb{T}_{pu}^2, \mathbb{T}_{pd}^2$  and  $\mathbb{T}_{ps}^2$  can be written as

$$\mathbb{T}_{pu}^2 = C_a \quad \text{with } a = \frac{2-\varphi}{\varphi} \quad \text{and}$$

$$\mathbb{T}_{pd}^2 = C_{-a}$$

$$\mathbb{T}_{ps}^2 = C_0$$

Then the mean curvature of  $\mathbb{T}_{pu}^2$  is

$$H_a = H_{\frac{2-\varphi}{\varphi}} = \sqrt{\frac{2-\varphi}{\varphi}}$$

and the mean curvature of  $\mathbb{T}_{pd}^2$  is

$$H_{-a} = H_{\frac{\varphi-2}{\varphi}} = -\sqrt{\frac{2-\varphi}{\varphi}} = -H_a$$

and the mean curvature of  $\mathbb{T}_{ps}^2$  is

$$H_0 = 0$$

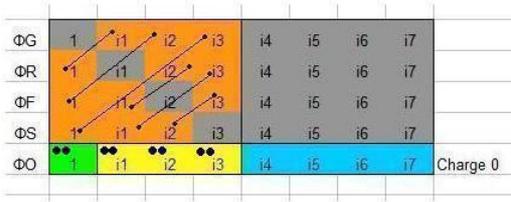
The shape of the W – Boson then is :

$$\text{topology } W_{Boson} = \left\{ C_{\frac{\varphi-2}{\varphi}}, C_0, C_{\frac{2-\varphi}{\varphi}} \right\} \subset S^3_{\sqrt{\varphi}}$$

## <5> Curvaturesensors by the Oktoquintenfield

The construction comes from multiplications (symmetric to the diagonal) by 2 degrees of freedom (complex subspaces). With this construction the tensor is symmetric in the diagonal.

$$\phi = \begin{bmatrix} \phi^G \\ \phi^R \\ \phi^F \\ \phi^S \\ \phi^O \end{bmatrix} = \begin{bmatrix} \phi_0^G + i_1 \cdot \phi_1^G + i_2 \cdot \phi_2^G + i_3 \cdot \phi_3^G + i_4 \cdot \phi_4^G + i_5 \cdot \phi_5^G + i_6 \cdot \phi_6^G + i_7 \cdot \phi_7^G \\ \phi_0^R + i_1 \cdot \phi_1^R + i_2 \cdot \phi_2^R + i_3 \cdot \phi_3^R + i_4 \cdot \phi_4^R + i_5 \cdot \phi_5^R + i_6 \cdot \phi_6^R + i_7 \cdot \phi_7^R \\ \phi_0^F + i_1 \cdot \phi_1^F + i_2 \cdot \phi_2^F + i_3 \cdot \phi_3^F + i_4 \cdot \phi_4^F + i_5 \cdot \phi_5^F + i_6 \cdot \phi_6^F + i_7 \cdot \phi_7^F \\ \phi_0^S + i_1 \cdot \phi_1^S + i_2 \cdot \phi_2^S + i_3 \cdot \phi_3^S + i_4 \cdot \phi_4^S + i_5 \cdot \phi_5^S + i_6 \cdot \phi_6^S + i_7 \cdot \phi_7^S \\ \phi_0^O + i_1 \cdot \phi_1^O + i_2 \cdot \phi_2^O + i_3 \cdot \phi_3^O + i_4 \cdot \phi_4^O + i_5 \cdot \phi_5^O + i_6 \cdot \phi_6^O + i_7 \cdot \phi_7^O \end{bmatrix}$$



$c$ ...speed of light  
 $G$ ...Gravitationconstant  
 $\hbar$ ...Planckconstant

symmetric Curvature Tensor

$$C_{em} = \frac{c}{G \cdot \hbar} \cdot \begin{bmatrix} \phi_0^O \cdot \phi_0^O & i_1 \cdot \phi_0^R \cdot \phi_1^G & i_2 \cdot \phi_0^F \cdot \phi_2^G & i_3 \cdot \phi_0^S \cdot \phi_3^G \\ i_1 \cdot \phi_0^R \cdot \phi_1^G & -\phi_1^O \cdot \phi_1^O & i_3 \cdot \phi_1^F \cdot \phi_2^R & -i_2 \cdot \phi_1^S \cdot \phi_3^R \\ i_2 \cdot \phi_0^F \cdot \phi_2^G & i_3 \cdot \phi_1^F \cdot \phi_2^R & -\phi_2^O \cdot \phi_2^O & i_1 \cdot \phi_2^S \cdot \phi_3^F \\ i_3 \cdot \phi_0^S \cdot \phi_3^G & -i_2 \cdot \phi_1^S \cdot \phi_3^R & i_1 \cdot \phi_2^S \cdot \phi_3^F & -\phi_3^O \cdot \phi_3^O \end{bmatrix}$$

10 independent fields.

remark :

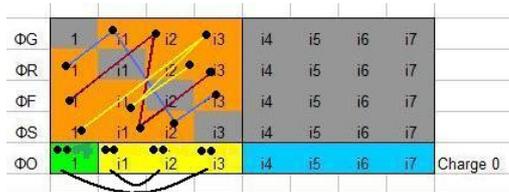
for  $\phi = c$  we get as curvature the plank curvature which is the reciprocal of the plank area.  
 The value of the curvature is :

$$0,34 \times 10^{70} \frac{1}{m^2}$$

**Second CURVATURE TENSOR from the Oktoquintenfield (generates a spinpotential)**

The construction comes from multiplications by 4 degrees of freedom (quaternionic subspaces).  
 With this construction the tensor is symmetric in both diagonals.

$$\phi = \begin{bmatrix} \phi^G \\ \phi^R \\ \phi^F \\ \phi^S \\ \phi^O \end{bmatrix} = \begin{bmatrix} \phi_0^G + i_1 \cdot \phi_1^G + i_2 \cdot \phi_2^G + i_3 \cdot \phi_3^G + i_4 \cdot \phi_4^G + i_5 \cdot \phi_5^G + i_6 \cdot \phi_6^G + i_7 \cdot \phi_7^G \\ \phi_0^R + i_1 \cdot \phi_1^R + i_2 \cdot \phi_2^R + i_3 \cdot \phi_3^R + i_4 \cdot \phi_4^R + i_5 \cdot \phi_5^R + i_6 \cdot \phi_6^R + i_7 \cdot \phi_7^R \\ \phi_0^F + i_1 \cdot \phi_1^F + i_2 \cdot \phi_2^F + i_3 \cdot \phi_3^F + i_4 \cdot \phi_4^F + i_5 \cdot \phi_5^F + i_6 \cdot \phi_6^F + i_7 \cdot \phi_7^F \\ \phi_0^S + i_1 \cdot \phi_1^S + i_2 \cdot \phi_2^S + i_3 \cdot \phi_3^S + i_4 \cdot \phi_4^S + i_5 \cdot \phi_5^S + i_6 \cdot \phi_6^S + i_7 \cdot \phi_7^S \\ \phi_0^O + i_1 \cdot \phi_1^O + i_2 \cdot \phi_2^O + i_3 \cdot \phi_3^O + i_4 \cdot \phi_4^O + i_5 \cdot \phi_5^O + i_6 \cdot \phi_6^O + i_7 \cdot \phi_7^O \end{bmatrix}$$



$$C_{spin} = \frac{1}{G \cdot \hbar \cdot c} \cdot \begin{bmatrix} -\phi_0^O \cdot \phi_0^O \cdot \phi_3^O \cdot \phi_3^O & -C & B & -A \\ -C & \phi_1^O \cdot \phi_1^O \cdot \phi_2^O \cdot \phi_2^O & -A & B \\ B & -A & \phi_1^O \cdot \phi_1^O \cdot \phi_2^O \cdot \phi_2^O & -C \\ -A & B & -C & -\phi_0^O \cdot \phi_0^O \cdot \phi_3^O \cdot \phi_3^O \end{bmatrix}$$

$$A = \phi_0^S \cdot \phi_3^G \cdot \phi_1^F \cdot \phi_2^R \quad B = \phi_0^F \cdot \phi_2^G \cdot \phi_1^S \cdot \phi_3^R \quad C = \phi_0^R \cdot \phi_1^G \cdot \phi_2^S \cdot \phi_3^F$$

5 independent fields A,B,C and two in the diagonal (blue and yellow).

**So finally we get three derivation- or curvaturetensors of the Oktoquintenpotential for twisted spacetime excitation**

0-th Curvaturetensor

$\phi_0^O$	$i_1 \cdot \phi_1^G$	$i_2 \cdot \phi_2^G$	$i_3 \cdot \phi_3^G$
$\phi_0^R$	$i_1 \cdot \phi_1^O$	$i_2 \cdot \phi_2^R$	$i_3 \cdot \phi_3^R$
$\phi_0^F$	$i_1 \cdot \phi_1^F$	$i_2 \cdot \phi_2^O$	$i_3 \cdot \phi_3^F$
$\phi_0^S$	$i_1 \cdot \phi_1^S$	$i_2 \cdot \phi_2^S$	$i_3 \cdot \phi_3^O$

$\phi_i$  unit is speed m/s

1-Curvaturetensor Cem em = energy-momentum

$\phi_0^O \cdot \phi_0^O$	$i_1 \cdot \phi_0^R \cdot \phi_1^G$	$i_2 \cdot \phi_0^F \cdot \phi_2^G$	$i_3 \cdot \phi_0^S \cdot \phi_3^G$
$i_1 \cdot \phi_0^R \cdot \phi_1^G$	$-\phi_1^O \cdot \phi_1^O$	$i_3 \cdot \phi_1^F \cdot \phi_2^R$	$-i_2 \cdot \phi_1^S \cdot \phi_3^R$
$i_2 \cdot \phi_0^F \cdot \phi_2^G$	$i_3 \cdot \phi_1^F \cdot \phi_2^R$	$-\phi_2^O \cdot \phi_2^O$	$i_1 \cdot \phi_2^S \cdot \phi_3^F$
$i_3 \cdot \phi_0^S \cdot \phi_3^G$	$-i_2 \cdot \phi_1^S \cdot \phi_3^R$	$i_1 \cdot \phi_2^S \cdot \phi_3^F$	$-\phi_3^O \cdot \phi_3^O$

This tensor is up to a constant equal to the energy-momentum tensor.

Energydensity :

$$\frac{E_{i,j}}{m^3} = \phi_i \cdot \phi_j \cdot \frac{c}{G \cdot \hbar} \cdot \frac{c^4}{8 \cdot \pi \cdot G} = \phi_i \cdot \phi_j \cdot \frac{c^2}{K \cdot l_p^2}$$

c...speed of light

G...Gravitationconstant

$\hbar$ ...Planckconstant

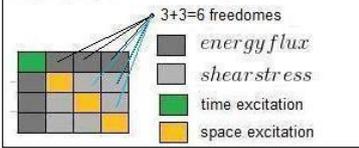
K..Einsteinconstant

$l_p^2$ ...Planckarea

the vacuumexcitation :

$$\phi_i^2 = \Lambda \cdot \frac{G \cdot \hbar}{c} \quad i = 0, 1, 2, 3$$

responsible for SO(1,3) symmetry



2-Curvaturetensor Cspin

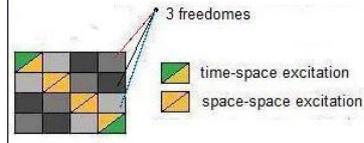
	-C	B	-A
-C		-A	B
B	-A		-C
-A	B	-C	

$$= -\phi_0^O \cdot \phi_0^O \cdot \phi_3^O \cdot \phi_3^O$$

$$= \phi_1^O \cdot \phi_1^O \cdot \phi_2^O \cdot \phi_2^O$$

$$\begin{aligned} A &= \phi_0^S \cdot \phi_3^G \cdot \phi_1^F \cdot \phi_2^R \\ B &= \phi_0^F \cdot \phi_2^G \cdot \phi_1^S \cdot \phi_3^R \\ C &= \phi_0^R \cdot \phi_1^G \cdot \phi_2^S \cdot \phi_3^F \end{aligned}$$

responsible for SU(2) symmetry



**<6> Extension of the ART by the second curvaturetensor**

The Oktoquintenpotential has two symmetric curvaturetensors. This motivates us to extend the Einstein Equation.

$$\text{Vacuumsymmetry} = \begin{pmatrix} SO(1,3) \\ SU(2) \end{pmatrix} = \begin{pmatrix} \text{rotations, boosts} \\ \text{spin} \end{pmatrix} \quad \begin{pmatrix} \text{real energydensity} \\ \text{imaginaer energydensity} \end{pmatrix}$$

I think this shows that the GR (General Relativity) has to be extended by an imaginäry part (spinpart) to be a consistent quantumtheorie.

So finally we expect something like **GR+i.GR<sup>o</sup>** where GR<sup>o</sup> is the spinpart.

with the two curvaturetensors  $C_{em}$  and  $C_{spin}$  we can define following equation :

$$\text{Real}(C_{em}) + \frac{i}{\varphi} \cdot C_{spin} = \frac{8 \cdot \pi \cdot G}{c^4} \cdot (T_{\mu\nu} + \frac{i}{\varphi} \cdot S_{\mu\nu})$$

where the real part is the GR and the imaginaer part is GR<sup>o</sup>

GR...General Relativity

GR<sup>o</sup>...Spinextension of GR

$\varphi$ ...golden ratio

the operator  $Real(A)$  is defined by

$$Real \begin{pmatrix} a_{0,0} & i_1.a_{0,1} & i_2.a_{0,2} & i_3.a_{0,3} \\ i_1.a_{1,0} & a_{1,1} & i_3.a_{1,2} & i_2.a_{1,3} \\ i_2.a_{2,0} & i_3.a_{2,1} & a_{2,2} & i_1.a_{2,3} \\ i_3.a_{3,0} & i_2.a_{3,1} & i_1.a_{3,2} & a_{3,3} \end{pmatrix} = \begin{pmatrix} a_{0,0} & a_{0,1} & a_{0,2} & a_{0,3} \\ a_{1,0} & a_{1,1} & a_{1,2} & a_{1,3} \\ a_{2,0} & a_{2,1} & a_{2,2} & a_{2,3} \\ a_{3,0} & a_{3,1} & a_{3,2} & a_{3,3} \end{pmatrix}$$

the reversing  $Real^{-1}$  is :

$$Real^{-1} \begin{pmatrix} a_{0,0} & a_{0,1} & a_{0,2} & a_{0,3} \\ a_{1,0} & a_{1,1} & a_{1,2} & a_{1,3} \\ a_{2,0} & a_{2,1} & a_{2,2} & a_{2,3} \\ a_{3,0} & a_{3,1} & a_{3,2} & a_{3,3} \end{pmatrix} = \begin{pmatrix} a_{0,0} & i_1.a_{0,1} & i_2.a_{0,2} & i_3.a_{0,3} \\ i_1.a_{1,0} & a_{1,1} & i_3.a_{1,2} & i_2.a_{1,3} \\ i_2.a_{2,0} & i_3.a_{2,1} & a_{2,2} & i_1.a_{2,3} \\ i_3.a_{3,0} & i_2.a_{3,1} & i_1.a_{3,2} & a_{3,3} \end{pmatrix}$$

$i_1, i_2, i_3 \dots$ imaginaer quaternions

more detailed with the two curvaturetensors of the oktoquintenfield:

$$\frac{8.\pi.G}{c^4} \cdot (T_{\mu\nu} + \frac{i}{\varphi} \cdot S_{\mu\nu}) = Real(C_{em}) + \frac{i}{\varphi} \cdot C_{spin} =$$

$$= \frac{c}{G.h} \left[ \begin{pmatrix} \phi_0^0.\phi_0^0 & \phi_0^R.\phi_1^G & \phi_1^F.\phi_2^G & \phi_2^S.\phi_3^G \\ sym. & -\phi_1^0.\phi_1^0 & \phi_1^F.\phi_2^R & -\phi_1^S.\phi_3^R \\ sym. & sym. & -\phi_2^0.\phi_2^0 & \phi_2^S.\phi_3^F \\ sym. & sym. & sym. & -\phi_3^0.\phi_3^0 \end{pmatrix} + \frac{i}{c^2\varphi} \begin{pmatrix} -\phi_0^0.\phi_0^0.\phi_3^0.\phi_3^0 - \phi_0^R.\phi_1^G.\phi_2^S.\phi_3^F & \phi_0^F.\phi_2^G.\phi_1^S.\phi_3^R - \phi_0^S.\phi_3^G.\phi_1^F.\phi_2^R \\ sym. & \phi_1^0.\phi_1^0.\phi_2^0.\phi_2^0 - \phi_0^S.\phi_3^G.\phi_1^F.\phi_2^R & \phi_0^F.\phi_2^G.\phi_1^S.\phi_3^R - \phi_0^S.\phi_3^G.\phi_1^F.\phi_2^R \\ sym. & sym. & \phi_1^0.\phi_1^0.\phi_2^0.\phi_2^0 - \phi_0^S.\phi_3^G.\phi_1^F.\phi_2^R & -\phi_0^R.\phi_1^G.\phi_2^S.\phi_3^F \\ sym. & sym. & sym. & -\phi_0^0.\phi_0^0.\phi_3^0.\phi_3^0 \end{pmatrix} \right]$$

10 different products
5 different products  
|
sym. and the red products are redundant  
|
 $2 \times S^1 \times S^1 \times SU(2) = 2 \times T^2 \times SU(2)$  with  $T^2 = Torus$

generates Poincare group  $\mathbb{R}^{1,3} \times O(1,3)$ 
generates  $2 \times$  spinning Torus  $T_{\Lambda^0}^2 = S^1 \times S^1$   
|
flat Clifford - Torus  
|
generates flat twisting and expanding vacuum

$T^2 \dots$  Clifford - Torus

This flat torus is a subset of the unit 3 - sphere  $S^3$ .

The Clifford torus divides the 3 - sphere into two congruent solid tori.

The Clifford - Torus embedded in  $S^3$  becomes a minimal surface.

The second curvaturetensor  $C_{spin}$  is determined by the first curvaturetensor  $C_{em}$  because its components are a mix of the components of  $C_{em}$ .

The vacuum part of the extended Einstein equation then is:

$$Vacuum = \Lambda.g + i.\Lambda^o.g^o$$

$\Lambda \dots$ cosmological constant

$\Lambda^o \dots$ second cosmological constant

$\Lambda$  and  $\Lambda^o$  comes from the Oktoquintenpotential (see picture).

$g = \eta$  comes from the first curvaturetensor

$g^o = \eta^o$  comes from the second curvaturetensor

then

$$Vacuum Energydensity = \frac{c^4}{8\pi G} \cdot [\Lambda \cdot \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} + i.\Lambda^o \cdot \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}] \begin{matrix} \begin{pmatrix} \oplus \\ \oplus \\ \oplus \\ \oplus \end{pmatrix} \\ \begin{pmatrix} \ominus \\ \ominus \\ \ominus \\ \ominus \end{pmatrix} \end{matrix} \begin{matrix} T_{up}^2 \\ T_{down}^2 \end{matrix}$$

hyperbolic (1,3)
ultrahyperbolic (2,2)

|
generates flat expanding spacetime
|
generates  $2 \times$  spinning Torus  $T_{\Lambda^0}^2 = S^1 \times S^1$   
|
flat Clifford - Torus  
|
generates flat twisting and expanding vacuum

$$T_{\Lambda^0}^2 = \begin{pmatrix} \oplus \\ -1 \\ \oplus \\ +1 \end{pmatrix} = \begin{pmatrix} C_- \\ C_+ \end{pmatrix}$$

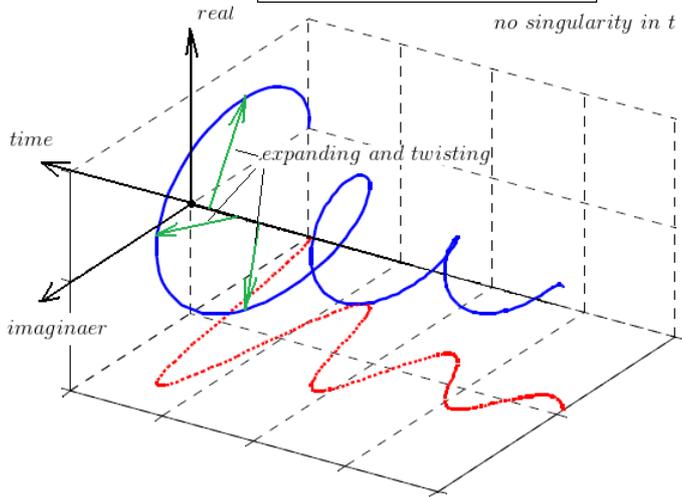
Spinor

Hint :  $\Lambda^o$  comes from the third part of the Oktoquintenpotential.

$$\Lambda^o = \frac{\Lambda}{\varphi} \text{ with } \varphi \dots \text{golden ratio}$$

$$Dirac-Spinor = \begin{pmatrix} \oplus \\ -1 \\ \oplus \\ +1 \end{pmatrix} \quad \text{Weyl-Spinor} = \begin{pmatrix} \oplus \\ -1 \\ \oplus \\ +1 \end{pmatrix} + \begin{pmatrix} \oplus \\ +1 \\ \oplus \\ +1 \end{pmatrix} \quad \text{or} \quad \begin{pmatrix} \oplus \\ +1 \\ \oplus \\ -1 \end{pmatrix} - \begin{pmatrix} \oplus \\ -1 \\ \oplus \\ -1 \end{pmatrix}$$

Universe is similar to  $f(t) = R_o \cdot e^{t+i \cdot t} = R_o \cdot e^t \cdot e^{i \cdot t}$   
 real no singularity in  $t = 0$



The question now is how does the spin of particles act on the second curvature tensor?  
 The second curvature tensor has 5 different exciteable values.  
 $D_o, D_i$  ( $D$  outer,  $D$  inner) in the diagonal and  $A, B, C$  out off the diagonal.  
 $A, B, C$  appears 4 times in the tensor and  $D_o, D_i$  2 times.  
 $16 \text{ components} = 4 \cdot A + 4 \cdot B + 4 \cdot C + 2 \cdot D_o + 2 \cdot D_i$   
 Now we want to assign the fields  $A, B, C, D_o, D_i$  to the different spins of particles  $0, \frac{1}{2}, 1, \frac{3}{2}, 2$ .

$-D_o$	$-C$	$B$	$-A$
$-C$	$D_i$	$-A$	$B$
$B$	$-A$	$D_i$	$-C$
$-A$	$B$	$-C$	$-D_o$

**Example free Diracparticle**

First we want to rearrange the Diracspinor by a matrix.  
 This makes the relation between Diracparticles like electrons and the curvature tensors better visible.

$$\psi_D^G = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 1 & 0 & 0 \end{pmatrix} \cdot \psi_D \quad \text{Dirac-Spinor} = \begin{pmatrix} (-1) \text{ up} \\ (-1) \text{ down} \\ (+1) \text{ up} \\ (+1) \text{ down} \end{pmatrix} \Rightarrow \begin{pmatrix} (-1) \text{ up} \\ (+1) \text{ up} \\ (+1) \text{ down} \\ (-1) \text{ down} \end{pmatrix}$$

**<6.1> The complete curvature Tensor for a free Spin 0 particle:**

$$\frac{8 \cdot \pi \cdot G}{c^4} \cdot (T_{\mu\nu} + \frac{i}{\varphi} \cdot S_{\mu\nu}) = \frac{c}{G \cdot h} \cdot \left( \begin{matrix} \text{Minkowski} \\ \begin{matrix} \blacksquare & \square & \square & \square \\ \square & \square & \square & \square \\ \square & \square & \square & \square \\ \square & \square & \square & \square \end{matrix} \end{matrix} + \frac{i}{c^2 \varphi} \cdot \begin{matrix} \text{double Torus} \\ \begin{matrix} \square & \square & \square & \square \\ \square & \square & \square & \square \\ \square & \square & \square & \square \\ \square & \square & \square & \square \end{matrix} \end{matrix} \right) \begin{matrix} \text{Spin UP} \\ \text{Spin DOWN} \end{matrix}$$

$$= \frac{c}{G \cdot h} \cdot \left[ \begin{pmatrix} \varphi_0^2 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} + \frac{i}{c^2 \varphi} \cdot \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} \right]$$

■ excited  
 $\varphi$ ...golden ratio

**<6.2> The complete curvature Tensor for a free standstill Electron is:**

Curvature by an not moving Electron

$$\frac{8\pi G}{c^4} \cdot (T_{\mu\nu} + \frac{i}{\varphi} \cdot S_{\mu\nu}) = \frac{c}{G \cdot h} \cdot \left( \begin{array}{c} \text{Minkowski} \\ \begin{array}{|c|c|c|c|} \hline \color{red}{+} & & & \\ \hline & \color{yellow}{\square} & & \\ \hline & & & \\ \hline & & & \color{red}{-} \\ \hline \end{array} + \frac{i}{c^2 \varphi} \cdot \begin{array}{c} \text{double Torus} \\ \begin{array}{|c|c|c|c|} \hline \color{red}{-} & & & \\ \hline & \color{yellow}{\square} & & \\ \hline & & & \\ \hline & & & \color{red}{-} \\ \hline \end{array} \end{array} \right) \begin{array}{l} \text{Spin UP} \\ \text{Spin DOWN} \end{array}$$

$$= \frac{c}{G \cdot h} \cdot \left[ \begin{pmatrix} \phi_0^2 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -\phi_3^2 \end{pmatrix} + \frac{i}{c^2 \varphi} \cdot \begin{pmatrix} -\phi_0^2 \cdot \phi_3^2 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -\phi_0^2 \cdot \phi_3^2 \end{pmatrix} \right]$$

■ excited  
 $\varphi$ ...golden ratio  
 pressure on a random spacedirection z

Our experience with electrons say us that they have a attractive gravity.  
 To get a attractive gravity we need a positiv pressure.  
 To get a positiv pressure we are forced to let  $\phi_j^0$  be imaginaer.

$$\phi_j^0 = i \cdot \omega_j^0 \quad \text{and} \quad \phi_0^0 = \omega_0^0$$

for  $j = 1, 2, 3$

we get the Spacetime and Torus curvature by an electron

$$\frac{8\pi G}{c^4} \cdot (T_{\mu\nu} + \frac{i}{\varphi} \cdot S_{\mu\nu}) = \frac{c}{G \cdot h} \cdot \left[ \begin{pmatrix} \omega_0^2 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & \omega_3^2 \end{pmatrix} + \frac{i}{c^2 \varphi} \cdot \begin{pmatrix} \omega_0^2 \cdot \omega_3^2 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & \omega_0^2 \cdot \omega_3^2 \end{pmatrix} \right]$$

positive pressure on a random spacedirection z  
 $\omega$  real

the Spacetime and Torus curvature by a positron is

$$\frac{8\pi G}{c^4} \cdot (T_{\mu\nu} + \frac{i}{\varphi} \cdot S_{\mu\nu}) = \frac{c}{G \cdot h} \cdot \left[ \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & \omega_1^2 & 0 & 0 \\ 0 & 0 & \omega_2^2 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} + \frac{i}{c^2 \varphi} \cdot \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & \omega_1^2 \cdot \omega_2^2 & 0 & 0 \\ 0 & 0 & \omega_1^2 \cdot \omega_2^2 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} \right]$$

positive pressure  
 therefore a positron is falling down on earth

**<6.3> The complete curvature Tensor for a Photon (massless vectorboson) is:**

$$\frac{8\pi G}{c^4} \cdot (T_{\mu\nu} + \frac{i}{\varphi} \cdot S_{\mu\nu}) = \frac{c}{G \cdot h} \cdot \left( \begin{array}{c} \text{Minkowski} \\ \begin{array}{|c|c|c|c|} \hline \color{red}{+} & & & \\ \hline & \color{red}{-} & & \\ \hline & & & \\ \hline & & & \color{red}{-} \\ \hline \end{array} + \frac{i}{c^2 \varphi} \cdot \begin{array}{c} \text{double Torus} \\ \begin{array}{|c|c|c|c|} \hline \color{red}{-} & & & \\ \hline & \color{red}{+} & & \\ \hline & & & \\ \hline & & & \color{red}{+} \\ \hline \end{array} \end{array} \right) \begin{array}{l} \text{Spin +1} \\ \text{Spin -1} \end{array}$$

$$= \frac{c}{G \cdot h} \cdot \left[ \begin{pmatrix} \phi_0^2 & 0 & 0 & 0 \\ 0 & -\phi_1^2 & 0 & 0 \\ 0 & 0 & -\phi_2^2 & 0 \\ 0 & 0 & 0 & -\phi_3^2 \end{pmatrix} + \frac{i}{c^2 \varphi} \cdot \begin{pmatrix} -\phi_0^2 \cdot \phi_3^2 & 0 & 0 & 0 \\ 0 & \phi_1^2 \cdot \phi_2^2 & 0 & 0 \\ 0 & 0 & \phi_1^2 \cdot \phi_2^2 & 0 \\ 0 & 0 & 0 & -\phi_0^2 \cdot \phi_3^2 \end{pmatrix} \right]$$

■ excited  
 $\varphi$ ...golden ratio  
 trace=0

**<6.4> The complete curvature Tensor for a Graviton (massless tensorboson) is:**

$$\frac{8\pi G}{c^4} \cdot (T_{\mu\nu} + \frac{i}{\varphi} \cdot S_{\mu\nu}) = \frac{c}{G \cdot h} \cdot \left( \begin{array}{c} \text{Minkowski} \\ \begin{array}{|c|c|c|c|} \hline \color{red}{+} & & & \\ \hline & \color{red}{+} & & \\ \hline & & & \\ \hline & & & \color{red}{+} \\ \hline \end{array} + \frac{i}{c^2 \varphi} \cdot \begin{array}{c} \text{double Torus} \\ \begin{array}{|c|c|c|c|} \hline \color{red}{-} & & & \\ \hline & \color{red}{+} & & \\ \hline & & & \\ \hline & & & \color{red}{-} \\ \hline \end{array} \end{array} \right) \begin{array}{l} \text{Spin +2, +1} \\ \text{Spin -2, -1} \end{array}$$

$$= \frac{c}{G \cdot h} \cdot \left[ \begin{pmatrix} \phi_0^2 \cdot \phi_3^2 & 0 & 0 & \phi_0^2 \cdot \phi_3^2 \\ 0 & -\phi_1^2 \cdot \phi_1^2 & \phi_1^2 \cdot \phi_2^2 & 0 \\ 0 & \phi_1^2 \cdot \phi_2^2 & -\phi_2^2 \cdot \phi_2^2 & 0 \\ \phi_0^2 \cdot \phi_3^2 & 0 & 0 & -\phi_3^2 \cdot \phi_3^2 \end{pmatrix} + \frac{i}{c^2 \varphi} \cdot \begin{pmatrix} -\phi_0^2 \cdot \phi_0^2 \cdot \phi_3^2 \cdot \phi_3^2 & 0 & 0 & -\phi_0^2 \cdot \phi_3^2 \cdot \phi_1^2 \cdot \phi_2^2 \\ 0 & \phi_1^2 \cdot \phi_1^2 \cdot \phi_2^2 \cdot \phi_2^2 & -\phi_0^2 \cdot \phi_3^2 \cdot \phi_1^2 \cdot \phi_2^2 & 0 \\ 0 & -\phi_0^2 \cdot \phi_3^2 \cdot \phi_1^2 \cdot \phi_2^2 & \phi_1^2 \cdot \phi_1^2 \cdot \phi_2^2 \cdot \phi_2^2 & 0 \\ -\phi_0^2 \cdot \phi_3^2 \cdot \phi_1^2 \cdot \phi_2^2 & 0 & 0 & -\phi_0^2 \cdot \phi_0^2 \cdot \phi_3^2 \cdot \phi_3^2 \end{pmatrix} \right]$$

trace=0

**Getting a closed form for the extended General Relativity EGR.**

we know that our energy – momentum curvature tensor

$$Real(C_{em}) = R_{\mu\nu} - \frac{R}{2} \cdot g_{\mu\nu} + \Lambda \cdot g_{\mu\nu} \quad \text{and that}$$

$C_{spin}$  is defined by multiplication of tensorelements of  $C_{em}$

The question now is how can we express  $C_{spin}$  analogous to  $C_{em}$  above as terms of Riemann – Geometrie?

For that we define the operator for 4 x 4 matrices or tensors:

Matrixoperator Tau

$$\overline{A} = T + \diagup + \diagdown$$

Transpose      in the big diagonal from top right to bottom left.      and then in the small diagonals

$$\begin{pmatrix} a_{0,0} & a_{0,1} & a_{0,2} & a_{0,3} \\ a_{1,0} & a_{1,1} & a_{1,2} & a_{1,3} \\ a_{2,0} & a_{2,1} & a_{2,2} & a_{2,3} \\ a_{3,0} & a_{3,1} & a_{3,2} & a_{3,3} \end{pmatrix} \rightarrow \begin{pmatrix} a_{3,3} & a_{2,3} & a_{1,3} & a_{0,3} \\ a_{3,2} & a_{2,2} & a_{1,2} & a_{0,2} \\ a_{3,1} & a_{2,1} & a_{1,1} & a_{0,1} \\ a_{3,0} & a_{2,0} & a_{1,0} & a_{0,0} \end{pmatrix} \rightarrow \begin{pmatrix} a_{3,3} & a_{2,3} & a_{1,3} & a_{1,2} \\ a_{3,2} & a_{2,2} & a_{0,3} & a_{0,2} \\ a_{3,1} & a_{3,0} & a_{1,1} & a_{0,1} \\ a_{2,1} & a_{2,0} & a_{1,0} & a_{0,0} \end{pmatrix}$$

**A** **A<sup>τ</sup>**

and a

special simple matrices multiplication

$$C = A \cdot B$$

with

$$c_{i,j} = \begin{cases} +a_{i,j} \cdot b_{i,j} & \text{if } i = j \\ -a_{i,j} \cdot b_{i,j} & \text{if } i \neq j \end{cases}$$

then it is easy to see that

$$(A + B)^{\overline{\tau}} = A^{\overline{\tau}} + B^{\overline{\tau}}$$

with

$$C_{em} = \frac{c}{G \cdot \hbar} \cdot \begin{pmatrix} \phi_0^0 \cdot \phi_0^0 & i_1 \cdot \phi_0^R \cdot \phi_1^G & i_2 \cdot \phi_0^F \cdot \phi_2^G & i_3 \cdot \phi_0^S \cdot \phi_3^G \\ \text{sym.} & -\phi_1^0 \cdot \phi_1^0 & i_3 \cdot \phi_1^F \cdot \phi_2^R & -i_2 \cdot \phi_1^S \cdot \phi_3^R \\ \text{sym.} & \text{sym.} & -\phi_2^0 \cdot \phi_2^0 & i_1 \cdot \phi_2^S \cdot \phi_3^F \\ \text{sym.} & \text{sym.} & \text{sym.} & -\phi_3^0 \cdot \phi_3^0 \end{pmatrix}$$

and

$$C_{spin} = \frac{1}{G \cdot \hbar \cdot c} \cdot \begin{pmatrix} -\phi_0^0 \cdot \phi_0^0 \cdot \phi_3^0 \cdot \phi_3^0 - \phi_0^R \cdot \phi_1^G \cdot \phi_2^S \cdot \phi_3^F & \phi_0^F \cdot \phi_2^G \cdot \phi_1^S \cdot \phi_3^R - \phi_0^S \cdot \phi_3^G \cdot \phi_1^F \cdot \phi_2^R \\ \text{sym.} & \phi_1^0 \cdot \phi_1^0 \cdot \phi_2^0 \cdot \phi_2^0 - \phi_0^S \cdot \phi_3^G \cdot \phi_1^F \cdot \phi_2^R & \phi_0^F \cdot \phi_2^G \cdot \phi_1^S \cdot \phi_3^R \\ \text{sym.} & \text{sym.} & \phi_1^0 \cdot \phi_1^0 \cdot \phi_2^0 \cdot \phi_2^0 - \phi_0^R \cdot \phi_1^G \cdot \phi_2^S \cdot \phi_3^F \\ \text{sym.} & \text{sym.} & \text{sym.} & -\phi_0^0 \cdot \phi_0^0 \cdot \phi_3^0 \cdot \phi_3^0 \end{pmatrix}$$

it follows that

$$C_{spin} = \frac{G \cdot \hbar}{c^3} \cdot Real(C_{em}) \cdot Real(C_{em})^{\overline{\tau}} = l_p^2 \cdot Real(C_{em}) \cdot Real(C_{em})^{\overline{\tau}} \quad l_p \dots \text{Plancklength}$$

with

$$Real(C_{em}) = R_{\mu\nu} - \frac{R}{2} \cdot g_{\mu\nu} + \Lambda \cdot g_{\mu\nu} = G_{\mu\nu} + \Lambda \cdot g_{\mu\nu} = K_{\mu\nu}$$

and

$$Real(C_{em}) + \frac{i}{\varphi} \cdot C_{spin} = \frac{8 \cdot \pi \cdot G}{c^4} \cdot (T_{\mu\nu} + \frac{i}{\varphi} \cdot S_{\mu\nu})$$

we get the final compact result for the extension of General Relativity by

$$K_{\mu\nu} + \frac{i}{\varphi} \cdot l_p^2 \cdot K_{\mu\nu} \cdot K_{\mu\nu}^{\mathcal{A}} = \frac{8 \cdot \pi \cdot G}{c^4} \cdot (T_{\mu\nu} + \frac{i}{\varphi} \cdot S_{\mu\nu})$$

with

$$S_{\mu\nu} = l_p^2 \cdot T_{\mu\nu} \cdot T_{\mu\nu}^{\mathcal{A}} \dots \text{Spintensor}$$

$$K_{\mu\nu} = R_{\mu\nu} - \frac{R}{2} \cdot g_{\mu\nu} + \Lambda \cdot g_{\mu\nu}$$

$$K_{\mu\nu}^{\mathcal{A}} = R_{\mu\nu}^{\mathcal{A}} - \frac{R}{2} \cdot g_{\mu\nu}^{\mathcal{A}} + \Lambda \cdot g_{\mu\nu}^{\mathcal{A}}$$

$\varphi$ ...golden ratio

The real part is the known General Relativity.

The imaginaer part is the Spinextension of GR.

Hint : The Energy–Stress tensor is still symmetric with or without Spin!

The question now is what is it good for?

In the same way as energydensity, momentumdensity aso. warps spacetime spin or spindensity warps a  $\mathbb{T}^2$  Torus.

Take care that the multiplication of the tensor is not the normal tensormultiplication or matricesmultiplication. It is the above defined simple multiplication  $C_{ij} = +/- A_{ij} \cdot B_{ij}$ .

## <7> Candidates for the dark matter in the universe.

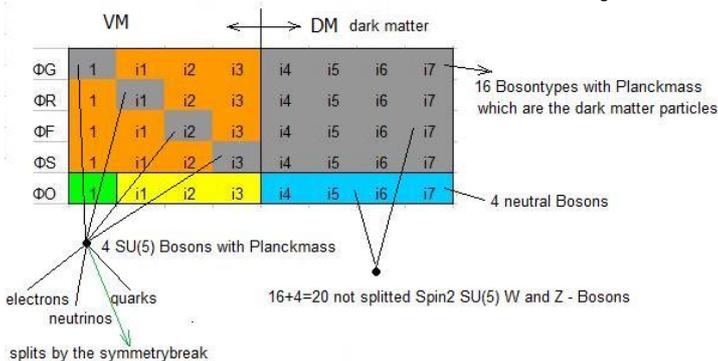
The Oktoquintefield can be divided into 2 areas (left and right).

The left one has 4 SU(5) Bosons which are over the timespace (curvature) fields.

I suppose that the Planckparticles with planckmass come into being by the Symmetriebreak **SU(5) x U(1) x U(1) -->U(1)** and the decomposition (Protons,electrons,...) of it comes from the other Symmetriebreak **SU(2)xU(1) --> U(1)** which only acts on the left half of the Oktoquintefield.

I suppose that the 4 W-Bosons in the left half (4 of the 20 charged SU(5) Bosons) split into protons,electrons and so on and the other SU(5) Bosons in the right half keep planckparticles.

So the left W-Bosons are the reason for the visible matter and the right W- and Z-Bosons are the reason vor the dark matter.



As seen in <4> we get a particle Lambda which is also a additional candidate for dark matter!

## <8> Some important points of the Oktoquintenpotential

To get the maxima, minima and the zeropoints of the potential we have to substitute  $z = \phi^2$  it is enough (because of symmetry) to take a look on the positive  $\phi^2$ s. and solve the cubic equations in the bracket

$$V(\sqrt{z}) = z \cdot \left( \frac{\gamma^2}{2} + \frac{\mu^2}{4} z + \frac{\lambda^2}{8} z^3 \right) \text{ and}$$

$$V'(\sqrt{z}) = \sqrt{z} \cdot (\gamma^2 + \mu^2 \cdot z + \lambda^2 \cdot z^3)$$

We will make it short and write the results.

First the Zeropoints :

$$z_1 = u + v = -\sqrt{\frac{2C}{3}} \cdot \left( \sqrt[3]{\frac{\sqrt{27} - i \cdot \sqrt{5}}{\sqrt{32}}} + \sqrt[3]{\frac{\sqrt{27} + i \cdot \sqrt{5}}{\sqrt{32}}} \right)$$

$$z_2 = \epsilon_1 \cdot u + \epsilon_2 \cdot v$$

$$z_3 = \epsilon_2 \cdot u + \epsilon_1 \cdot v$$

$$\text{Where } \epsilon_1 = -\frac{1}{2} + i \cdot \frac{\sqrt{3}}{2} \text{ and } \epsilon_2 = -\frac{1}{2} - i \cdot \frac{\sqrt{3}}{2}$$

then

$$z_1 = -0,990839414 \times 2 \cdot \sqrt{\frac{2C}{3}}$$

$$z_2 = 0,378466979 \times 2 \cdot \sqrt{\frac{2C}{3}}$$

$$z_3 = 0,612372435 \times 2 \cdot \sqrt{\frac{2C}{3}}$$

then the zeropoints are

$$\phi_1 = -0,995409169 \times \sqrt[4]{\frac{8 \cdot C}{3}}$$

$$\phi_2 = 0,615196699 \times \sqrt[4]{\frac{8 \cdot C}{3}}$$

$$\phi_3 = 0,782542290 \times \sqrt[4]{\frac{8 \cdot C}{3}}$$

$$\phi_2 = c = 0,615196699 \times \sqrt[4]{\frac{8 \cdot C}{3}} = \sin(37,966214178) \times \sqrt[4]{\frac{8 \cdot C}{3}}$$

$$\alpha_{c} = 37,966214178^\circ$$

$$\phi_3 = c \cdot \sqrt{\varphi} = 0,782542290 \times \sqrt[4]{\frac{8 \cdot C}{3}} = \sin(128,506061932) \times \sqrt[4]{\frac{8 \cdot C}{3}}$$

$$\alpha_{c \cdot \sqrt{\varphi}} = 128,506061932^\circ$$

Then the Maxima and the Minima :

$$z_1 = u + v = -\sqrt{\frac{C}{3}} \cdot \left( \sqrt[3]{\frac{\sqrt{37} - i \cdot \sqrt{27}}{\sqrt{64}}} + \sqrt[3]{\frac{\sqrt{37} + i \cdot \sqrt{27}}{\sqrt{64}}} \right)$$

$$z_2 = \epsilon_1 \cdot u + \epsilon_2 \cdot v$$

$$z_3 = \epsilon_2 \cdot u + \epsilon_1 \cdot v$$

Finally we have two positiv results :

$$z_{min} = 0,233475630 \times 2 \cdot \sqrt{\frac{C}{3}} \text{ and}$$

$$z_{max} = 0,725352944 \times 2 \cdot \sqrt{\frac{C}{3}}$$

and one negative

$$z_3 = -(z_{max} + z_{min})$$

Then because of  $z = \phi^2$

$$\phi_{min} = 0,483193160 \times \sqrt[4]{\frac{4 \cdot C}{3}} \text{ and}$$

$$\phi_{max} = 0,8516765489 \times \sqrt[4]{\frac{4 \cdot C}{3}}$$

In cubic equations the real zeropoints comes from the  $\cos(\alpha)$  or from  $\sin(90-\alpha)$  of angles (see [https://en.wikipedia.org/wiki/Cubic\\_function](https://en.wikipedia.org/wiki/Cubic_function)).

Then for  $\phi_{min}$  we get an angle  $\alpha_{min}$  :

$$\phi_{min} = 0,483193160 \times \sqrt[4]{\frac{4 \cdot C}{3}} = \sin(28,894160846) \times \sqrt[4]{\frac{4 \cdot C}{3}}$$

$$\alpha_{min} = 28,894160846 \text{ degrees is very near to the Weinberg angle}$$

$$\sin^2(\alpha_{min}) = \sin^2(28,894160846) = 0,233475630$$

and for  $\phi_{max}$  we get an angle

$$\phi_{max} = 0,8516765489 \times \sqrt[4]{\frac{4 \cdot C}{3}} = \sin(121,605508985) \times \sqrt[4]{\frac{4 \cdot C}{3}}$$

$$\alpha_{max} = 121,605508985 \text{ degrees}$$

$C = c^4 \cdot \varphi^2$   
 $c$ ...speed of light  
 $\varphi$ ...golden ratio

$$\phi_{min} = 0,483193160 \times \sqrt[4]{\frac{4.C}{3}} \quad \text{and}$$

$$\phi_{max} = 0,8516765489 \times \sqrt[4]{\frac{4.C}{3}}$$

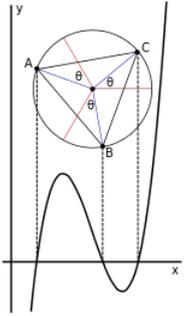
In cubic equations the real zeropoints comes from the  $\cos(\alpha)$  or from  $\sin(90-\alpha)$  of angles (see [https://en.wikipedia.org/wiki/Cubic\\_function](https://en.wikipedia.org/wiki/Cubic_function)).

Our second extreme value  $L_1$  is at  $e_1$

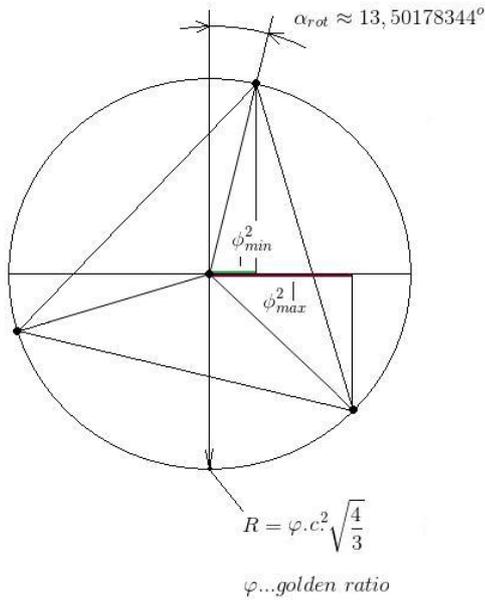
With the relation above we can calculate the third extreme value  $L_2$ .

$$\phi_{max} = 1,762600... \times e_1$$

### Geometric interpretation of the roots (zeropoints) in cubic equations with 3 real zeropoints



graphical zeropoints of the derivation of the  
(radicalized  $\phi^2 = z$ ) Oktoquintenpotential



### Conclusions

**Dark Energy** comes from the Oktoquintenpotential (the second term in the potential).

**Dark Matter** could be the W , Z Bosons of the SU(5) Symmetry and the Lambda Boson from the mixing.

