Some Properties of Conformal β -Change

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Abstract: We have considered the conformal β -change of the Finsler metric given by

$$L(x,y) \to \overline{L}(x,y) = e^{\sigma(x)} f(L(x,y),\beta(x,y)),$$

where $\sigma(x)$ is a function of x, $\beta(x,y) = b_i(x)y^i$ is a 1-form on the underlying manifold M^n , and $f(L(x,y), \beta(x,y))$ is a homogeneous function of degree one in L and β . We have studied quasi-C-reducibility, C-reducibility and semi-C-reducibility of the Finsler space with this metric. We have also calculated V-curvature tensor and T-tensor of the space with this changed metric in terms of v-curvature tensor and T-tensor respectively of the space with the original metric.

Key Words: Conformal change, β -change, Finsler space, quasi-C-reducibility, C-reducibility, semi-C-reducibility, V-curvature tensor, T-tensor.

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§1. Introduction

Let $F^n = (M^n, L)$ be an n-dimensional Finsler space on the differentialble manifold M^n equipped with the fundamental function L(x,y).B.N.Prasad and Bindu Kumari and C. Shibata [1,2] have studied the general case of β -change,that is, $L^*(x,y) = f(L,\beta)$,where f is positively homogeneous function of degree one in L and β , and β given by $\beta(x,y) = b_i(x)y^i$ is a one- form on M^n . The β -change of special Finsler spaces has been studied by H.S.Shukla, O.P.Pandey and Khageshwar Mandal [7].

The conformal theory of Finsler space was initiated by M.S. Knebelman [12] in 1929 and has been investigated in detail by many authors (Hashiguchi [8], Izumi[4,5] and Kitayama [9]). The conformal change is defined as $L^*(x, y) = e^{\sigma(x)}L(x, y)$, where $\sigma(x)$ is a function of position only and known as conformal factor. In 2008, Abed [15,16] introduced the change $\bar{L}(x, y) = e^{\sigma(x)}L(x, y) + \beta(x, y)$, which he called a β -conformal change, and in 2009 and 2010, Nabil L.Youssef, S.H.Abed and S.G. Elgendi [13,14] introduced the transformation $\bar{L}(x, y) = f(e^{\sigma}L, \beta)$, which is β -change of conformally changed Finsler metric L. They have not only established the relationships between some important tensors of (M^n, L) and the corresponding tensors of (M^n, \bar{L}) , but have also studied several properties of this change.

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We have changed the order of combination of the above two changes in our paper [6], where we have applied β -change first and conformal change afterwards, i.e.,

$$\bar{L}(x,y) = e^{\sigma(x)} f(L(x,y), \beta(x,y)), \qquad (1.1)$$

where $\sigma(x)$ is a function of x, $\beta(x, y) = b_i(x)y^i$ is a 1-form. We have called this change as conformal β -change of Finsler metric. In this paper we have investigated the condition under which a conformal β -change of Finsler metric leads a Douglas space into a Douglas space. We have also found the necessary and sufficient conditions for this change to be a projective change.

In the present paper, we investigate some properties of conformal β -change. The Finsler space equipped with the metric \overline{L} given by (1.1) will be denoted by $\overline{F^n}$. Throughout the paper the quantities corresponding to $\overline{F^n}$ will be denoted by putting bar on the top of them. We shall denote the partial derivatives with respect to x^i and y^i by ∂_i and $\dot{\partial}_i$ respectively. The Fundamental quantities of F^n are given by

$$g_{ij} = \dot{\partial}_i \dot{\partial}_j \frac{L^2}{2} = h_{ij} + l_i l_j, \quad l_i = \dot{\partial}_i L.$$

Homogeneity of f gives

$$Lf_1 + \beta f_2 = f, \tag{1.2}$$

where subscripts 1 and 2 denote the partial derivatives with respect to L and β respectively. Differentiating above equations with respect to L and β respectively, we get

$$Lf_{12} + \beta f_{22} = 0$$
 and $Lf_{11} + \beta f_{21} = 0.$ (1.3)

Hence we have

$$f_{11}/\beta^2 = (-f_{12})/L\beta = f_{22}/L^2,$$
 (1.4)

which gives

$$f_{11} = \beta^2 \omega, f_{12} = -L\beta \omega, f_{22} = L^2 \omega, \tag{1.5}$$

where Weierstrass function ω is positively homogeneous of degree -3 in L and β . Therefore

$$L\omega_1 + \beta\omega_2 + 3\omega = 0, \tag{1.6}$$

where ω_1 and ω_2 are positively homogeneous of degree -4 in L and β . Throughout the paper we frequently use the above equations without quoting them. Also we have assumed that f is not linear function of L and β so that $\omega \neq 0$.

The concept of concurrent vector field has been given by Matsumoto and K. Eguchi [11] and S. Tachibana [17], which is defind as follows:

The vector field b_i is said to be a concurrent vector field if

$$b_{i|j} = -g_{ij}$$
 $b_i|_j = 0,$ (1.7)

where small and long solidus denote the h- and v-covariant derivatives respectively. It has been

proved by Matsumoto that b_i and its contravariant components b^i are functions of coordinates alone. Therefore from the second equation of (1.7), we have $C_{ijk}b^i = 0$.

The aim of this paper is to study some special Finsler spaces arising from conformal β -change of Finsler metric,viz., quasi-C-reducible, C-reducible and semi-C-reducible Finsler spaces. Further, we shall obtain v-curvature tensor and T-tensor of this space and connect them with v-curvature tensor and T-tensor respectively of the original space.

§2. Metric Tensor and Angular Metric Tensor of \bar{F}^n

Differentiating equation (1.1) with respect to y^i we have

$$\bar{l}_i = e^{\sigma} (f_1 l_i + f_2 b_i). \tag{2.1}$$

Differentiating (2.1) with respect to y^j , we get

$$\bar{h}_{ij} = e^{2\sigma} \left(\frac{ff_1}{L} h_{ij} + fL^2 \omega m_i m_j \right), \qquad (2.2)$$

where $m_i = b_i - \frac{\beta}{L}L_i$.

From (2.1) and (2.2) we get the following relation between metric tensors of F^n and \bar{F}^n :

$$\bar{g}_{ij} = e^{2\sigma} \left[\frac{ff_1}{L} g_{ij} - \frac{p\beta}{L} l_i l_j + (fL^2\omega + f_2^2) b_i b_j + p(b_i l_j + b_j l_i) \right],$$
(2.3)

where $p = f_1 f_2 - f \beta L \omega$.

The contravariant components \bar{g}^{ij} of the metric tensor of \bar{F}^n , obtainable from $\bar{g}^{ij}\bar{g}_{jk} = \delta^i_k$, are as follows:

$$\bar{g}^{ij} = e^{-2\sigma} \left[\frac{L}{ff_1} g^{ij} + \frac{pL^3}{f^3 f_1 t} \left(\frac{f\beta}{L^2} - \Delta f_2 \right) l^i l^j - \frac{L^4 \omega}{ff_1 t} b^i b^j - \frac{pL^2}{f^2 f_1 t} (l^i b^j + l^j b^i) \right],$$
(2.4)

where $l^i = g^{ij}l_j$, $b^2 = b_i b^i$, $b^i = g^{ij}b_j$, g^{ij} is the reciprocal tensor of g_{ij} of F^n , and

$$t = f_1 + L^3 \omega \Delta, \Delta = b^2 - \frac{\beta^2}{L^2}.$$
 (2.5)

$$(a) \dot{\partial}_{i}f = e^{\sigma} \left(\frac{f}{L}l_{i} + f_{2}m_{i}\right), \quad (b) \dot{\partial}_{i}f_{1} = -e^{\sigma}\beta L\omega m_{i},$$

$$(c)\dot{\partial}_{i}f_{2} = e^{\sigma}L^{2}\omega m_{i}, \quad (d) \dot{\partial}_{i}p = -\beta qLm_{i},$$

$$(e) \dot{\partial}_{i}\omega = -\frac{3\omega}{L}l_{i} + \omega_{2}m_{i}, \quad (f)\dot{\partial}_{i}b^{2} = -2C_{..i},$$

$$(g) \dot{\partial}_{i}\Delta = -2C_{..i} - \frac{2\beta}{L^{2}}m_{i}, \qquad (2.6)$$

(a)
$$\dot{\partial}_i q = -\frac{3q}{L} l_i$$
, (b) $\dot{\partial}_i t = -2L^3 \omega C_{..i} + [L^3 \Delta \omega_2 - 3\beta L \omega] m_i$,
(c) $\dot{\partial}_i q = -\frac{3q}{L} l_i + (4f_2 \omega_2 + 3\omega^2 L^2 + f\omega_{22}) m_i$. (2.7)

§3. Cartan's C-Tensor and C-Vectors of \bar{F}^n

Cartan's covariant C-tensor C_{ijk} of F^n is defined by

$$\bar{C}_{ijk} = \frac{1}{4} \dot{\partial}_i \dot{\partial}_j \dot{\partial}_k L^2 = \dot{\partial}_k g_{ij}$$

and Cartan's C-vectors are defined as follows:

$$C_i = C_{ijk} g^{jk}, C^i = C^i_{jk} g^{jk}.$$
 (3.1)

We shall write $C^2 = C^i C_i$. Under the conformal β -chang (1.1) we get the following relation between Cartan's C-tensors of F^n and \bar{F}^n :

$$\bar{C}_{ijk} = e^{2\sigma} \left[\frac{ff_1}{L} C_{ijk} + \frac{p}{2L} (h_{ij}m_k + h_{jk}m_i + h_{ki}m_j) + \frac{qL^2}{2} m_i m_j m_k \right].$$
(3.2)

We have

(a)
$$m_i l^i = 0,$$

(b) $m_i b^i = b^2 - \frac{\beta^2}{L^2} = \Delta = b_i m^i,$
(c) $g_{ij} m^i = h_{ij} m^i = m_j.$
(3.3)

From (2.1), (2.3), (2.4) and (3.2), we get

$$\bar{C}_{ij}^{h} = C_{ij}^{h} + \frac{p}{2ff_{1}}(h_{ij}m^{h} + h_{j}^{h}m_{i} + h_{i}^{h}m_{j}) + \frac{qL^{3}}{2ff_{1}}m_{j}m_{k}m^{h} - \frac{L}{ft}C_{.jk}n^{h} - \frac{pL\Delta}{2f^{2}f_{1}t}h_{jk}n^{h} - \frac{2pL + qL^{4}\Delta}{2f^{2}f_{1}t}m_{j}m_{k}n^{h}, \qquad (3.4)$$

where $n^h = fL^2\omega b^h + pl^h$ and $h^i_j = g^{il}h_{lj}, C_{.ij} = C_{rij}b^r, C_{..i} = C_{rji}b^rb^j$ and so on.

Proposition 3.1 Let $\bar{F}^n = (M^n, \bar{L})$ be an n-dimensional Finsler space obtained from the conformal β -change of the Finsler space $F^n = (M^n, L)$, then the normalized supporting element \bar{l}_i , angular metric tensor \bar{h}_{ij} , fundamental metric tensor \bar{g}_{ij} and (h)hv-torsion tensor \bar{C}_{ijk} of \bar{F}^n are given by (2.1), (2.2), (2.3) and (3.2), respectively.

From (2.4),(3.1),(3.2) and (3.4) we get the following relations between the C-vectors of of F^n and \bar{F}^n and their magnitudes

$$\bar{C}_i = C_i - L^3 \omega C_{i..} + \mu m_i, \qquad (3.5)$$

where

$$\mu = \frac{p(n+1)}{2ff_1} - \frac{3pL^3\omega\Delta}{2ff_1} + \frac{qL^3\Delta(1-L^3\omega\Delta)}{2ff_1};$$

$$\bar{C}^i = \frac{e^{-2\sigma}L}{ff_1}C^i + M^i,$$
 (3.6)

where

$$M^{i} = \frac{\mu e^{-2\sigma}L}{ff_{1}}m^{i} - \frac{L^{4}\omega}{ff_{1}}C^{i}_{..} - \left(C_{i} - e^{2\sigma}L^{3}\omega C_{i..} + \mu\Delta\right)\left(\frac{L^{3}\omega}{ff_{1}}b^{i} + \frac{L}{ft}y^{i}\right)$$

and

where

$$\bar{C}^2 = \frac{e^{-2\sigma}}{p}C^2 + \lambda, \qquad (3.7)$$

$$\lambda = \left(\frac{e^{-2\sigma}L}{ff_1} - L^3\omega\Delta\right)\mu^2\Delta + \frac{2\mu e^{-2\sigma}L}{ff_1}C_{...}$$
$$- \left(1 + 2\mu\Delta\right)L^3\omega + \left(1 - 3\mu + e^{2\sigma}L^2\omega ff_1C_{..}\right)L^3\omega C_{...}$$
$$+ L^3\omega C_{..r}\left(\left(e^{4\sigma}L\omega f^2f_1^2C_{i..} - \mu\Delta\right)L^3\omega b^r - e^{2\sigma}L^2\omega ff_1C_{..}^r - 2C^r\right).$$

§4. Special Cases of \bar{F}^n

In this section, following Matsumoto [10], we shall investigate special cases of \bar{F}^n which is conformally β -changed Finsler space obtained from F^n .

Definition 4.1 A Finsler space (M^n, L) with dimension $n \ge 3$ is said to be quasi-C-reducible if the Cartan tensor C_{ijk} satisfies

$$C_{ijk} = Q_{ij}C_k + Q_{jk}C_i + Q_{ki}C_j, \qquad (4.1)$$

where Q_{ij} is a symmetric indicatory tensor.

The equation (3.2) can be put as

$$\bar{C}_{ijk} = e^{2\sigma} \left[\frac{ff_1}{L} C_{ijk} + \frac{1}{6} \pi_{(ijk)} \left\{ \left(\frac{3p}{L} h_{ij} + qL^2 m_i m_j \right) m_k \right\} \right],$$

where $\pi_{(ijk)}$ represents cyclic permutation and sum over the indices i,j and k.

Putting the value of m_k from equation (3.5) in the above equation, we get

$$\bar{C}_{ijk} = e^{2\sigma} \left[\frac{ff_1}{L} C_{ijk} + \frac{1}{6\mu} \pi_{(ijk)} \left\{ \left(\frac{3p}{L} h_{ij} + qL^2 m_i m_j \right) (\bar{C}_k - C_k + L^3 \omega C_{k..} \right) \right\} \right].$$

Rearranging this equation, we get

$$\bar{C}_{ijk} = e^{2\sigma} \left[\frac{ff_1}{L} C_{ijk} + \frac{1}{6\mu} \pi_{(ijk)} \left\{ \left(\frac{3p}{L} h_{ij} + qL^2 m_i m_j \right) \bar{C}_k \right\} + \frac{1}{6\mu} \pi_{(ijk)} \left\{ \left(\frac{3p}{L} h_{ij} + qL^2 m_i m_j \right) \left(L^3 \omega C_{k..} - C_k \right) \right\} \right].$$

Further rearrangement of this equations gives

$$\bar{C}_{ijk} = \pi_{(ijk)}(\bar{H}_{ij}\bar{C}_k) + U_{ijk}, \qquad (4.2)$$

where $\bar{H}_{ij} = \frac{e^{2\sigma}}{6\mu} \{ (\frac{3p}{L}h_{ij} + qL^2m_im_j), \text{ and } \}$

$$U_{ijk} = e^{2\sigma} \left[\frac{ff_1}{L} C_{ijk} + \frac{1}{6\mu} \pi_{(ijk)} \left\{ \left(\frac{3p}{L} h_{ij} + qL^2 m_i m_j \right) \left(L^3 \omega C_{k..} - C_k \right) \right\} \right]$$
(4.3)

Since H_{ij} is a symmetric and indicatory tensor, therefore from equation (4.2) we have the following theorem.

Theorem 4.1 Conformally β -changed Finsler space \overline{F}^n is quasi-C-reducible iff the tensor U_{ijk} of equation (4.3) vanishes identically.

We obtain a generalized form of Matsumoto's result [10] as a corollary of the above theorem.

Corollary 4.1 If F^n is Reimannian space, then the conformally β -changed Finsler space \overline{F}^n is always a quasi-C-reducible Finsler space.

Definition 4.2 A Finsler space (M^n, L) of dimension $n \ge 3$ is called C-reducible if the Cartan tensor C_{ijk} is written in the form

$$C_{ijk} = \frac{1}{n+1} (h_{ij}C_k + h_{ki}C_j + h_{jk}C_i).$$
(4.4)

Define the tensor $G_{ijk} = C_{ijk} - \frac{1}{(n+1)}(h_{ij}C_k + h_{ki}C_j + h_{jk}C_i)$. It is clear that G_{ijk} is symmetric and indicatory. Moreover, G_{ijk} vanishes iff F^n is C-reducible.

Proposition 4.1 Under the conformal β -change(1.1), the tensor \overline{G}_{ijk} associated with the space \overline{F}^n has the form

$$\bar{G}_{ijk} = e^{2\sigma} \frac{ff_1}{L} G_{ijk} + V_{ijk} \tag{4.5}$$

where

$$V_{ijk} = \frac{1}{(n+1)} \pi_{(ijk)} \{ (e^{2\sigma}(n+1)(\alpha_1 h_{ij} + \alpha_2 m_i m_j)m_k + e^{2\sigma}\omega L^2 m_i m_j C_k + e^{2\sigma}L^2 \omega (ff_1 h_{ij} + L^3 \omega m_i m_j)C_{k..} \},$$
(4.6)

$$\alpha_1 = \frac{e^{2\sigma}p}{2L} - \frac{\mu f f_1 e^{2\sigma}}{L(n+1)}, \quad \alpha_2 = \frac{e^{2\sigma}qL^2}{6} - \frac{\mu e^{2\sigma}\omega L^2}{(n+1)}.$$

From (4.5) we have the following theorem.

Theorem 4.2 Conformally β -changed Finsler space \overline{F}^n is C-reducible iff F^n is C-reducible and the tensor V_{ijk} given by (4.6) vanishes identically.

Definition 4.3 A Finsler space (M^n, L) of dimension $n \ge 3$ is called semi-C-reducible if the Cartan tensor C_{ijk} is expressible in the form:

$$C_{ijk} = \frac{r}{n+1} (h_{ij}C_k + h_{ki}C_j + h_{jk}C_i) + \frac{s}{C^2} C_i C_j C_k,$$
(4.7)

where r and s are scalar functions such that r + s = 1.

Using equations (2.2), (3.5) and (3.7) in equation (3.2), we have

$$\bar{C}_{ijk} = e^{2\sigma} \left[\frac{ff_1}{L} C_{ijk} + \frac{p}{2\mu f f_1} (\bar{h}_{ij}\bar{C}_k + \bar{h}_{ki}\bar{C}_j + \bar{h}_{jk}\bar{C}_i) + \frac{\Delta L(f_1q - 3p\omega)}{2ff_1\mu t\bar{C}^2} \bar{C}_i\bar{C}_j\bar{C}_k \right].$$

If we put

$$r' = \frac{p(n+1)}{2\mu f f_1}, s' = \frac{\Delta L(f_1 q - 3p\omega)}{2f f_1 \mu t}$$

we find that r' + s' = 1 and

$$\bar{C}_{ijk} = e^{2\sigma} \left[\frac{ff_1}{L} C_{ijk} + \frac{r'}{n+1} (\bar{h}_{ij}\bar{C}_k + \bar{h}_{ki}\bar{C}_j + \bar{h}_{jk}\bar{C}_i) + \frac{s'}{\bar{C}^2}\bar{C}_i\bar{C}_j\bar{C}_k \right].$$
(4.8)

From equation (4.8) we infer that \overline{F}^n is semi-C- reducible iff $C_{ijk} = 0$, i.e. iff F^n is a Reimannian space. Thus we have the following theorem.

Theorem 4.3 Conformally β -changed Finsler space \overline{F}^n is semi-C-reducible iff F^n is a Riemannian space.

§5. v-Curvature Tensor of \bar{F}^n

The v-curvature tensor [10] of Finsler space with fundamental function L is given by

$$S_{hijk} = C_{ijr}C^r_{hk} - C_{ikr}C^r_{hj}$$

Therefore the v-curvature tensor of conformally β -changed Finsler space \bar{F}^n will be given by

$$\bar{S}_{hijk} = \bar{C}_{ijr}\bar{C}^r_{hk} - \bar{C}_{ikr}\bar{C}^r_{hj}.$$
(5.1)

From equations (3.2) and (3.4), we have

$$\bar{C}_{ijr}\bar{C}_{hk}^{r} = e^{2\sigma} \left[\frac{ff_{1}}{L}C_{ijr}C_{hk}^{r} + \frac{p}{2L}(C_{ijk}m_{h} + C_{ijh}m_{k} + C_{ihk}m_{j} + C_{hjk}m_{i}) + \frac{pf_{1}}{2Lt}(C_{.ij}h_{hk} + C_{hk}h_{ij}) - \frac{ff_{1}L^{2}\omega}{t}C_{.ij}C_{.hk} + C_{hk}h_{ij} \right]$$

$$+\frac{p^{2}\Delta}{4fLt}h_{hk}h_{ij} + \frac{L^{2}(qf_{1}-2p\omega)}{2t}(C_{.ij}m_{k}m_{h} + C_{.hk}m_{i}m_{j}) \\ +\frac{p(p+L^{3}q\Delta)}{4Lft}(h_{ij}m_{h}m_{k} + h_{hk}m_{i}m_{j}) + \frac{p^{2}}{4Lff_{1}}(h_{ij}m_{h}m_{k} + h_{hk}m_{i}m_{j} + h_{hj}m_{i}m_{k} + h_{hi}m_{j}m_{k} + h_{jk}m_{i}m_{h} + h_{ik}m_{h}m_{j}) \\ + \frac{L^{2}(2pqt + (qf_{1}-2p\omega)(2p+L^{3}q\Delta))}{4ff_{1}t}m_{i}m_{j}m_{h}m_{k} \bigg].$$
(5.2)

We get the following relation between v-curvature tensors of (M^n, L) and (M^n, \overline{L}) :

$$\bar{S}_{hijk} = e^{2\sigma} \left[\frac{ff_1}{L} S_{hijk} + d_{hj} d_{ik} - d_{hk} d_{ij} + E_{hk} E_{ij} - E_{hj} E_{ik} \right],$$
(5.3)

where

$$d_{ij} = PC_{ij} - Qh_{ij} + Rm_i m_j, ag{5.4}$$

$$E_{ij} = Sh_{ij} + Tm_i m_j, \tag{5.5}$$

$$P = L\left(\frac{s}{t}\right)^{1/2}, \quad Q = \frac{pg}{2L^2\sqrt{st}}, \\ R = \frac{L\left(2\omega p - f_1q\right)}{2\sqrt{st}}, \\ S = \frac{p}{2L^2\sqrt{f\omega}}, \quad T = \frac{L\left(qf_1 - \omega p\right)}{2f_1\sqrt{f\omega}}.$$

Proposition 5.1 The relation between v-curvature tensors of F^n and \overline{F}^n is given by (5.3).

When b_i in β is a concurrent vector field, then $C_{ij} = 0$. Therefore the value of v-curvature tensor of \bar{F}^n as given by (5.3) is reduced to the extent that $d_{ij} = Rm_im_j - Qh_{ij}$.

§6. The T-Tensor T_{hijk}

The T-tensor of F^n is defined in [3] by

$$T_{hijk} = LC_{hij} \mid_k + C_{hij}l_k + C_{hik}l_j + C_{hjk}l_i + C_{ijk}l_h,$$

$$(6.1)$$

where

$$C_{hij} \mid_{k} = \dot{\partial}_{k} C_{hij} - C_{rij} C_{hk}^{r} - C_{hrj} C_{ik}^{r} - C_{hir} C_{jk}^{r}.$$

$$(6.2)$$

In this section we compute the T-tensor of \bar{F}^n , which is given by

$$\bar{T}_{hijk} = \bar{L}\bar{C}_{hij}\bar{|}_k + \bar{C}_{hij}\bar{l}_k + \bar{C}_{hik}\bar{l}_j + \bar{C}_{hjk}\bar{l}_i + \bar{C}_{ijk}\bar{l}_h,$$
(6.3)

where

$$\bar{C}_{hij}\bar{|}_k = \dot{\partial}_k \bar{C}_{hij} - \bar{C}_{rij}\bar{C}^r_{hk} - \bar{C}_{hrj}\bar{C}^r_{ik} - \bar{C}_{hir}\bar{C}^r_{jk}.$$
(6.4)

The derivatives of m_i and h_{ij} with respect to y^k are given by

$$\dot{\partial}_k m_i = -\frac{\beta}{L^2} h_{ik} - \frac{1}{L} (l_i m_k), \quad \dot{\partial}_k h_{ij} = 2C_{ijk} - \frac{1}{L} (l_i h_{jk} + l_j h_{ki}) \tag{6.5}$$

From equations (3.2) and (6.5), we have

$$\dot{\partial}_{k}\bar{C}_{hij} = e^{2\sigma} \left[\frac{ff_{1}}{L} \partial_{k}C_{hij} + \frac{p}{L} (C_{ijk}m_{h} + C_{ijh}m_{k} + C_{ihk}m_{j} + C_{hjk}m_{i}) \right. \\
\left. - \frac{p\beta}{2L^{3}} (h_{ij}h_{hk} + h_{hj}h_{ik} + h_{ih}h_{jk}) + \frac{p}{2L^{2}} (h_{jk}l_{h}m_{i} + h_{hk}l_{j}m_{i} \\
\left. + h_{hk}l_{i}m_{j} + h_{ik}l_{h}m_{j} + h_{jk}l_{i}m_{h} + h_{jk}l_{h}m_{i} + h_{ij}l_{h}m_{k} + h_{hj}l_{i}m_{k} \\
\left. + h_{ik}l_{j}m_{k} + h_{ij}l_{k}m_{h} + h_{jh}l_{k}m_{i} + h_{hi}l_{k}m_{j}) - \frac{\beta q}{2} (h_{ij}m_{h}m_{k} \\
\left. + h_{hk}m_{i}m_{j} + h_{hj}m_{i}m_{k} + h_{hi}m_{j}m_{k} + h_{jk}m_{i}m_{h} + h_{ik}m_{h}m_{j}) \right. \\
\left. - \frac{qL}{2} (l_{i}m_{j}m_{h}m_{k} + l_{j}m_{i}m_{h}m_{k} + l_{h}m_{i}m_{j}m_{k} + h_{k}m_{i}m_{j}m_{h}) \\
\left. + \frac{L^{2}}{2} (4f_{2}\omega_{2} + 3L^{2}\omega^{2} + f\omega_{22})m_{h}m_{i}m_{j}m_{k} \right].$$
(6.6)

Using equations (6.5) and (5.2) in equation(6.4), we get

$$\begin{split} \bar{C}_{hij}\bar{|}_{k} &= e^{2\sigma}\frac{ff_{1}}{L}C_{hij}|_{k} - \frac{e^{2\sigma}p}{2L}\left(C_{ijk}m_{h} + C_{ijh}m_{k} + C_{ihk}m_{j} + C_{hjk}m_{i}\right) \\ &- pe^{2\sigma}\left(\frac{2f\beta t}{4fL^{3}t} + \frac{L^{2}p\Delta}{4fL^{3}t}\right)\left(h_{ij}h_{hk} + h_{hj}h_{ik} + h_{ih}h_{jk}\right) - e^{2\sigma}\left(\frac{\beta q}{2}\right) \\ &+ \frac{p^{2}f_{1} + pqf_{1}L^{3}\Delta + 3p^{2}}{4Lff_{1}t}\right)\left(h_{ij}m_{h}m_{k} + h_{hk}m_{i}m_{j} + h_{hj}m_{i}m_{k} + h_{hi}m_{j}m_{k}\right) \\ &+ h_{jk}m_{i}m_{h} + h_{ik}m_{h}m_{j}\right) - \frac{e^{2\sigma}p}{2L^{2}}[l_{h}(h_{jk}m_{i} + h_{ij}m_{k} \\ &+ h_{ik}m_{j}) + l_{j}(h_{hk}m_{i} + h_{ik}m_{kh} + h_{ih}m_{k}) + l_{i}(h_{hk}m_{j} + h_{jk}m_{h} \\ &+ h_{hj}m_{k}\right) + l_{k}(h_{ij}m_{h} + h_{jh}m_{i} + h_{hi}m_{j})] - \frac{e^{2\sigma}qL}{2}(l_{im}jm_{h}m_{k} \\ &+ l_{j}m_{i}m_{h}m_{k} + l_{h}m_{i}m_{j}m_{k} + h_{k}m_{i}m_{j}m_{h}\right) - \frac{pf_{1}e^{2\sigma}}{2Lt}\left(C_{.ij}h_{hk} \\ &+ C_{.hj}h_{ik} + C_{.hk}h_{ij} + C_{.ik}h_{h} + C_{.hi}h_{jk} + C_{.jk}h_{hi}\right) + \frac{e^{2\sigma}ff_{1}L^{2}\omega}{t}\left(C_{.ij}C_{.hk} \\ &+ C_{.hj}C_{.ik} + C_{.hi}C_{.jk}\right) - \frac{e^{2\sigma}L^{2}(qf_{1} - 2p\omega)}{2t}\left(C_{.ij}m_{k}m_{h} \\ &+ C_{.hi}m_{j}m_{k} + C_{.jk}m_{h}m_{i}\right) + e^{2\sigma}\left[\frac{L^{2}(4f_{2}\omega_{2} + 3L^{2}\omega^{2} + f\omega_{22})}{2} \\ &- \frac{3L^{2}(2pqt + (qf_{1} - 2p\omega)(2p + L^{3}q\Delta)}{4ff_{1}t}\right]m_{i}m_{j}m_{h}m_{k}. \end{split}$$

Using equations (2.1), (3.2) and (6.6) in equation (6.3), we get the following relation

between T-tensors of Finsler spaces F^n and \bar{F}^n :

$$\bar{T}_{hijk} = e^{3\sigma} \left[\frac{f^2 f_1}{L^2} T_{hijk} + \frac{f(f_1 f_2 + f\beta L\omega)}{2L} \left(C_{ijk} m_h + C_{ijh} m_k + C_{ihk} m_j + C_{hjk} m_i \right) + \frac{f^2 f_1 L^2 \omega}{t} \left(C_{.ij} C_{.hk} + C_{.hj} C_{.ik} + C_{.hi} C_{.jk} \right) \\
- \frac{p f_1}{2Lt} \left(C_{.ij} h_{hk} + C_{.hj} h_{ik} + C_{.hk} h_{ij} + C_{.ik} h_h + C_{.hi} h_{jk} + C_{.jk} h_{hi} \right) \\
- \frac{f L^2 (q f_1 - 2p \omega)}{2t} \left(C_{.ij} m_k m_h + C_{.hk} m_{imj} + C_{.hj} m_{imk} + C_{.hj} m_{imk} + C_{.ik} m_j m_h + C_{.hi} m_{jmk} + C_{.hi} m_{jmk$$

Proposition 6.1 The relation between T-tensors of F^n and \overline{F}^n is given by (6.7).

If bi is a concurrent vector field in F^n , then $C_{ij} = 0$. Therefore from (6.8), we have

$$\bar{T}_{hijk} = e^{3\sigma} \left[\frac{f^2 f_1}{L^2} T_{hijk} - \frac{p(2f\beta t + L^2 p\Delta)}{4L^3 t} (h_{ij}h_{hk} + h_{hj}h_{ik} + h_{ih}h_{jk}) - \left(\frac{p^2 f_1 + pqf_1 L^3 \Delta + 3p^2 t}{4Lf_1 t} + \frac{\beta qf}{2} - \frac{pf_2}{L} \right) (h_{ij}m_h m_k + h_{hk}m_i m_j + h_{hj}m_i m_k + h_{hi}m_j m_k + h_{jk}m_i m_h + h_{ik}m_h m_j) + \left[2L^2 f_2 q + \frac{L^2 (4f_2 \omega_2 + 3L^2 \omega^2 + f\omega_{22})}{2} + \frac{3L^2 (qf_1 - 2p\omega)(2p + L^3 q\Delta)}{4Lff_1 t} - \frac{3L^2 2pqt}{4Lff_1 t} \right] m_i m_j m_h m_k \right].$$
(6.9)

If bi is a concurrent vector field in F^n , with vanishing T-tensor then T-tensor of F^n is given by

$$\bar{T}_{hijk} = e^{3\sigma} \left[-\frac{p(2f\beta t + L^2p\Delta)}{4L^3t} (h_{ij}h_{hk} + h_{hj}h_{ik} + h_{ih}h_{jk}) - \left(\frac{p^2f_1 + pqf_1L^3\Delta + 3p^2t}{4Lf_1t} + \frac{\beta qf}{2} - \frac{pf_2}{L} \right) (h_{ij}m_hm_k + h_{hk}m_im_j + h_{hj}m_im_k + h_{hi}m_jm_k + h_{jk}m_im_h + h_{ik}m_hm_j) + \left[\frac{L^2(4f_2\omega_2 + 3L^2\omega^2 + f\omega_{22})}{2} - \frac{3L^22pqt}{4Lf_1t} + \frac{3L^2(qf_1 - 2p\omega)(2p + L^3q\Delta)}{4Lf_1t} + 2L^2f_2q \right] m_im_jm_hm_k \right].$$
(6.10)

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