Kinematics of point motion along curves of the second order

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Abstract

The motion of a material point along curves of the second order is represented by a kinematic equation. Formulas for the dependence of acceleration and radius, speed and radius are derived. The direction of the velocity and acceleration vectors is determined. The conditions for the conservation of Kepler's laws when a material point moves along second-order curves are shown.

Keywords

Kepler's laws, ellipse, speed, acceleration, radius

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If simple equations of speed and acceleration are sufficient to describe rectilinear motion: V = S/t, $a = S/t^2$, then differential equations of motion are needed to solve problems on the curvilinear motion of material points and their systems. *"The way we derive these equations doesn't matter"*: [1,§11, π .3].

I. Formulas for the dependence of acceleration and radius, speed and radius

Point C moves in an ellipse relative to the focus, Figure 1.



Figure 1

There is a system of equations for a parametric pendulum (1)

The parameter is time (t).

$$\begin{cases} x = r(\varphi(t)) \cdot \cos(\varphi(t)) \\ y = r(\varphi(t)) \cdot \sin(\varphi(t)) \end{cases}$$
(1.1)

Let us substitute into system (1) the radius of the ellipse with respect to the focus:

$$r(\varphi(t)) = \frac{b^2}{a(1 - e * cos(\varphi(t)))}$$
(1.2)

$$\begin{cases} x = \frac{b^2}{a(1 - e * cos(\varphi(t)))} \cdot cos(\varphi(t)) \\ y = \frac{b^2}{a(1 - e * cos(\varphi(t)))} \cdot sin(\varphi(t)) \end{cases}$$
(1.3)

Let's differentiate twice. We get the coordinates of speed and acceleration:

$$\dot{x} = \frac{d}{dt} \left(r(\varphi(t)) \cos(\varphi(t)) \right) = -\frac{b^2 \ast \dot{\varphi} \ast \sin(\varphi(t))}{a(e \ast \cos(\varphi(t)) - 1)^2} = \frac{r^2 \ast \dot{\varphi} \ast \sin(\varphi(t))}{e \ast \cos(\varphi(t)) - 1}$$
(1.4)

$$\dot{y} = \frac{d}{dt} \left(\frac{p}{1 - e \cdot \cos(\varphi(t))} \sin(\varphi(t)) \right) = \frac{b^2 \cdot \dot{\varphi} \cdot (-e + \cos(\varphi(t)))}{a(e \cdot \cos(\varphi(t)) - 1)^2} = \frac{r^2 \cdot \varphi \cdot (-e + \cos(\varphi(t)))}{1 - e \cdot \cos(\varphi(t))}$$
(1.5)

$$\ddot{x} = \frac{b^2 \left(\left(-e * \cos(\varphi(t)) * \sin(\varphi(t)) + \sin(\varphi(t)) \right) \ddot{\varphi} + \dot{\varphi}^2 \left(e * \cos(\varphi(t))^2 - 2e + \cos(\varphi(t)) \right) \right)}{a \left(e * \cos(\varphi(t)) - 1 \right)^3}$$
(1.6)

$$\ddot{y} = \frac{-b^2 \left((-\cos(\varphi(t))(e * \cos(\varphi(t)) - 1) + e) \ddot{\varphi} + 2\dot{\varphi}^2 \left(e^2 - \frac{e * \cos(\varphi(t)) + 1}{2} \right) \sin(\varphi(t)) \right)}{a (e * \cos(\varphi(t)) - 1)^3}$$
(1.7)

Velocity
$$v = \sqrt{\dot{x}^2 + \dot{y}^2} = \frac{b^2 \cdot \dot{\varphi} \cdot \sqrt{1 + e^2 - 2e \cdot \cos \varphi(t)}}{a(-1 + e \cdot \cos \varphi(t))^2} = \frac{r \cdot \dot{\varphi} \cdot \sqrt{1 + e^2 - 2e \cdot \cos \varphi(t)}}{(1 - e \cdot \cos \varphi(t))}$$
 (1.8)

Acceleration $\dot{v} = \sqrt{\ddot{x}^2 + \ddot{y}^2} =$

$$b^{2}\begin{pmatrix}\frac{\sqrt{(e^{2}-2e*cos(\varphi(t))+1)(e*cos(\varphi(t))-1)^{2}*\dot{\varphi}^{2}}}{a(e*cos(\varphi(t))-1)^{3}}+\\\frac{\sqrt{4(e^{2}-\frac{3*e*cos(\varphi(t))+1}{2})\dot{\varphi}^{2}(e*cos(\varphi(t))\sin(\varphi(t))-1)\ddot{\varphi}}}{a(e*cos(\varphi(t))-1)^{3}}-\\\frac{\sqrt{4\dot{\varphi}^{4}(-cos(\varphi(t))^{3}e^{3}+(e^{4}-\frac{e^{2}}{4})cos(\varphi(t))^{2}+(e^{3}+\frac{e}{2})cos(\varphi(t))-e^{4}-\frac{1}{4})}}{a(e*cos(\varphi(t))-1)^{3}}\end{pmatrix}$$
(1.9)

We form a system of equations from (1.6), (1.7) and solve for φ [°]. We obtain the kinematic equation of motion of a point relative to the focus along second-order curves:

$$\ddot{\varphi} = \frac{2 \cdot e \cdot \sin(\varphi) \cdot \dot{\varphi}^2}{1 - e \cdot \cos(\varphi)} \tag{1.10}$$

At different values of eccentricity, the shape of the curve will change.

We substitute (1.10) into (1.9), and simplify:

$$\dot{\nu} = \frac{b^2 \dot{\varphi}^2}{a(1 - e * \cos(\varphi))^2} = \frac{r * \dot{\varphi}^2}{1 - e * \cos(\varphi)}$$
(1.11)

The sector speed is constant:

$$k = r_p^2 * \dot{\phi}_p = r_i^2 * \dot{\phi}_i = r_a^2 * \dot{\phi}_a = const,$$
(1.12)

$$\dot{\varphi} = \frac{k}{r^2} \tag{1.13}$$

where r_p is the perifocal distance, r_a is the apofocal distance

We substitute (1.13) into (1.11):

$$\dot{\nu} = \frac{k^2}{r^3(1 - e * \cos(\varphi))}$$
(1.14)

The acceleration v is recalculated using formula (14). Results (9) and (14) are compared, Figure 2.



Figure 2

We substitute (1.13) into (1.8):



Figure 3

Formulas (1.14, 1.15) do not give any advantage for calculating the modulus of speed and acceleration. First, to calculate the sector constant k, you need to calculate the angular velocity once. Secondly, in order for the motion of a point to comply with Kepler's laws, the angle must change according to elliptic equations. The value of these formulas is in the logical definition of the dependence of speed and acceleration on the radius.

II. Velocity and acceleration vectors

Let's consider two variants of point movement, Figure 4: 1) - movement relative to the center 2) - movement relative to the focus.



v - speed, a - acceleration, dx, dy, ddx, ddy - first and second derivatives along the coordinate axes.

Figure 4

Note the property of collinear vectors on the plane - rectangles built on vectors, Figure 5, should be similar:

$$\frac{BD}{AD} = \frac{B_1 D_1}{A_1 D_1} \tag{2.1}$$





1. movement relative to focus

Let's compare the ratio of the coordinates of the radius and acceleration:

$$\frac{x}{y} = \frac{\cos\varphi}{\sin\varphi} \tag{2.2}$$

$$\frac{\ddot{x}}{\ddot{y}} = \frac{(-2e^2\cos^2\varphi + 3e^2 - 1)\cos(\varphi)}{\sin(\varphi)(e^2 - 1)(2e^2\cos^2\varphi + 1)}$$
(2.3)

If e = 0 we get a circle and $\frac{\ddot{x}}{\ddot{y}} = \frac{x}{y}$, a special case of an ellipse, figure 5.

In Figures 5 - 7 they are marked with red lines for speed, green for acceleration.

$$\frac{d^2}{dt^2}\varphi(t) = 0,$$
рисунок 6 (2.5)

Coordinates of the beginning of the velocity and acceleration vectors, points of the initial ellipse (x, y). The coordinates of the end of the velocity vector (dx+x, dy+y). Acceleration vector end coordinates (ddx+x, ddy+y).

Velocity, Acceleration, a = 0.5000, b = 0.5000, days = 80.00



Figure 6

If $e \neq 0$, then $\frac{\ddot{x}}{\ddot{y}} \neq \frac{x}{y}$, Figure 7

(2.6)

(2.4)



Figure 7

2. movement relative to the center, figure 4.1)

$$r(\varphi(t)) = \frac{b}{\sqrt{1 - e^2 \cos^2 \varphi(t)}}$$
(2.7)

To derive the kinematic equation of motion of a point relative to the center, we will replace the radius formula (2) with (13) in the system of equations (1.1).

Let's differentiate twice. We get the coordinates of speed and acceleration:

$$\dot{x} = \frac{d}{dt} \left(\frac{b * \cos(\varphi(t))}{\sqrt{1 - e^2 * \cos(\varphi(t))^2}} \right) = -\frac{b * \sin(\varphi)}{(1 - e^2 * \cos^2 \varphi)^{3/2}}$$
(2.8)

$$\dot{y} = \frac{d}{dt} \left(\frac{b * sin(\varphi(t))}{\sqrt{1 - e^2 * cos(\varphi(t))^2}} \right) = \frac{b(1 - e^2) cos(\varphi)}{(1 - e^2 * cos^2 \varphi)^{3/2}}$$
(2.9)

$$\ddot{x} = \frac{d^2}{dt^2} \left(\frac{b * \cos(\varphi(t))}{\sqrt{1 - e^2 * \cos(\varphi(t))^2}} \right) = -\frac{b * \cos(\varphi) (2e^2 \cos^2 \varphi - 3e^2 + 1)}{(1 - e^2 * \cos^2 \varphi)^{5/2}}$$
(2.10)

$$\ddot{y} = \frac{d^2}{dt^2} \left(\frac{b * sin(\varphi(t))}{\sqrt{1 - e^2 * cos(\varphi(t))^2}} \right) = \frac{b * sin(\varphi)(e^2 - 1)(2e^2 cos^2 \varphi + 1)}{(1 - e^2 * cos^2 \varphi)^{5/2}}$$
(2.11)

$$\nu = \sqrt{\dot{x}^2 + \dot{y}^2} = \sqrt{\frac{b^2 \dot{\varphi}^2 \left(1 - 2e^2 \cos(\varphi(t))^2 + e^4 \cos(\varphi(t))^2\right)}{\left(1 - e^2 \cos(\varphi(t))^2\right)^3}}$$
(2.12)

We solve for $\ddot{\varphi}$. We obtain the kinematic equation of motion of a point relative to the center along second-order curves:

$$\ddot{\varphi} = \frac{2 * e^2 * \cos(\varphi) * \sin(\varphi) * \dot{\varphi}^2}{1 - e^2 * \cos(\varphi)^2}$$
(2.13)

Let's compare the ratio of the coordinates of the radius and acceleration:

$$\frac{x}{y} = \frac{\cos\varphi}{\sin\varphi}$$
(2.14)

$$\frac{\ddot{x}}{\ddot{y}} = \frac{(-2e^2\cos^2\varphi + 3e^2 - 1)\cos(\varphi)}{\sin(\varphi)(e^2 - 1)(2e^2\cos^2\varphi + 1)}$$
(2.15)

If
$$e = 0$$
 we get a circle and $\frac{\ddot{x}}{\ddot{y}} = \frac{x}{y}$, (2.16)
a special case of an ellipse, Figure 5.

Eccentricity e = 0. Substitute in equation (2.15)

$$\frac{d^2}{dt^2}\varphi(t) = 0, \text{ Figure 6}$$
If $e \neq 0$, then $\frac{\ddot{x}}{\ddot{y}} \neq \frac{x}{\dot{y}}$, Figure 8
(2.17)
Velocity, Acceleration, a = 0.8000, b = 0.7000, days = 80.00



Figure 8

V. Trammel of Archimedes



Figure 9

Any point on the ellipsograph ruler moves along an elliptical path around the center.

In order not to refer the reader to the sources, we present the derivation of the formulas necessary for calculating velocities, accelerations, and rotation angles.

Ruler *AB* moves from horizontal to vertical position, Figure 9. Point *C* describes $\frac{1}{4}$ of the ellipse. The direction of the instantaneous rotation of the ruler *AB* around *P*_{AB} is clockwise in accordance with the direction of the known velocity vector of point *A*.

Speeds of points *B* and *C*:

$$\omega_{AB} = \frac{v_A}{AP_{AB}} \tag{3.1}$$

$$\boldsymbol{v}_B = \omega_{AB} * BP_{AB} = \boldsymbol{v}_A \frac{BP_{AB}}{AP_{AB}} \tag{3.2}$$

Vector v_c is directed perpendicular to *CP*.

$$\boldsymbol{v}_{C} = \omega_{AB} * CP_{AB} = v_{A} \frac{CP_{AB}}{AP_{AB}}$$
(3.3)

The directions of the velocities of the points $\overrightarrow{v_B}$ and $\overrightarrow{v_C}$ are determined by the instantaneous rotation of the ruler *AB* around the instantaneous center of velocities P_{AB} .

Determination of accelerations of points B and C

Let's use the theorem - acceleration of points of a flat figure. Point A will be a pole, since the acceleration of point A is known.

The vector equation for the acceleration of point *B* has the form:

$$\overrightarrow{a_B} = \overrightarrow{a_A} + \overrightarrow{a_{BA}^r} + \overrightarrow{a_{BA}^c}$$
(3.4)

where $\overrightarrow{a_A}$ – is the acceleration of the pole A (given);

 $\overline{a_{BA}^r}$ and $\overline{a_{BA}^c}$ – are the rotational and centripetal accelerations of the point B in the rotation of the ruler around the pole A. In this case:

$$a_{BA}^c = \omega_{AB}^2 * BA \tag{3.5}$$

The vector $\overrightarrow{a_B}$ is located perpendicular to the ruler *AB*, its direction is unknown, since the direction of the angular acceleration ε_{AB} is unknown.

In equation (3.4) there are two unknowns: accelerations $\overrightarrow{a_A}$ and $\overrightarrow{a_{BA}}$, which can be determined from the equations of vector equality projections onto the directions of axes AX and AY:

$$\begin{cases} a_{Bx} = a_{Ax} + a_{BAx}^{r} + a_{BAx}^{c} \\ a_{By} = a_{Ay} + a_{BAy}^{r} + a_{BAy}^{c} \end{cases}$$
(3.7)

The direction of the vectors $\overrightarrow{a_B}$ and $\overrightarrow{a_{BA}^r}$ is chosen arbitrarily. The solution of system (3.7) allows one to find the numerical value $\overrightarrow{a_B}$ and $\overrightarrow{a_{BA}^r}$ with a plus or minus sign. A positive value indicates the correctness of the chosen direction of the vectors $\overrightarrow{a_B}$ and $\overrightarrow{a_{BA}^r}$ a negative value indicates the need to change their direction.

$$a_{A} = \sqrt{(a_{Ax})^{2} + (a_{Ay})^{2}}, \ a_{AB}^{r} = \sqrt{(a_{ABx}^{r})^{2} + (a_{ABy}^{r})^{2}}$$
(3.8)

Ruler angular acceleration:

$$\boldsymbol{\varepsilon}_{AB} = \frac{a_{BA}^r}{BA} \tag{3.9}$$

The acceleration of point C is determined by the equation:

$$\overrightarrow{a_C} = \overrightarrow{a_A} + \overrightarrow{a_{CA}^r} + \overrightarrow{a_{CA}^c}$$
(3.10)



Figure 10

where $(\overrightarrow{a_{CA}^r} \text{ and } \overrightarrow{a_{CA}^c})$ are, respectively, the rotational and centripetal accelerations of the point *C* relative to the pole *A*:

$$\boldsymbol{a}_{CA}^c = \boldsymbol{\omega}_{AB}^2 * AC \tag{3.11}$$

$$\boldsymbol{a}_{CB}^{r} = \boldsymbol{\varepsilon}_{AB} * AC \tag{3.12}$$

Vector $\overrightarrow{a_{CA}^c}$ is located on *CA* and is directed from point *C* to pole *A*. Vector $\overrightarrow{a_{CA}^r}$ is perpendicular to *CA* and directed in the same direction as $\overrightarrow{a_{BA}^r}$, Figure 10.

Equation (3.10) can be represented in projections on the axes Ax and Ay:

$$\begin{cases} a_{Cx} = a_{Ax} + a_{CAx}^{c} + a_{CAx}^{r} \\ a_{Cy} = a_{Ay} + a_{CAy}^{c} + a_{CAy}^{r} \end{cases}$$
(3.13)

The acceleration projections of point *C* are determined from (3.10). The direction of the vector $\vec{a_C}$ is determined by the signs of the projections \boldsymbol{a}_{Cx} and \boldsymbol{a}_{Cy} .

Vector modulus:

$$\boldsymbol{a}_{C} = \sqrt{(a_{Cx})^{2} + (a_{Cy})^{2}}$$
(3.14)

Let's take a look at the different travel options

T is the period specified by arbitrary units of time. AB = a + b, $A(0, y_A) B(x_B, 0)$. Initial coordinates of points: A(0,0), B(a + b, 0), C(a, 0). Initial speed $v_{A0} = 0$.

Uniform movement

Given: point *C* divides *AB* into segments *a* and *b*, $A(0, y_A)$, $B(x_B, 0)$, initial A(0, 0), B(AB, 0). *A* moves uniformly from $O \rightarrow Y$. Accelerations $a_A = 0$, $a_B = 0$, speed

$$\boldsymbol{v}_A = \frac{AB*4}{T} \tag{3.15}$$

Find: y_{A_i} , x_{C_i} , y_{C_i} , v_{C_i} , a_{C_i} , ϕ_i

Solution

Coordinates $A(0, y_{A_i})$:

$$y_{A_i} = \boldsymbol{v}_A * i \tag{3.16}$$

Further, according to equations (3.4) - (3.14)

Coordinates $B(x_{B_i}, 0)$:

$$\sin \alpha = \frac{y_{A_i}}{AB}, \alpha = \operatorname{asin} \frac{y_{A_i}}{AB}$$
(3.17)

$$x_{B_i} = \cos\alpha * \text{AB}, y_{B_i} = 0 \tag{3.18}$$

$$\boldsymbol{\omega}_{AB} = \frac{v_A}{AP_{AB}} = \frac{v_A}{x_{B_i}} \tag{3.19}$$

$$\boldsymbol{\nu}_B = \boldsymbol{\omega}_{AB} * B \boldsymbol{P}_{AB} = \boldsymbol{\omega}_{AB} * \boldsymbol{y}_{A_i} \tag{3.20}$$

From equation (5) $\boldsymbol{a}_{BA}^{c} = \boldsymbol{\omega}_{AB}^{2} * BA$

$$\begin{cases} \boldsymbol{a}_{Bx} = \boldsymbol{a}_{BA}^{c} * \cos \alpha + \boldsymbol{a}_{BA}^{r} * \sin \alpha \\ \boldsymbol{0} = \boldsymbol{a}_{Ay} + \boldsymbol{a}_{BA}^{c} * \sin \alpha + \boldsymbol{a}_{BA}^{r} * \cos \alpha \end{cases}$$
(3.21)

Solving the resulting equations, we find a_B ,

$$\boldsymbol{a}_{BA}^{r} = \frac{-a_{Ay} - a_{BA}^{c} \sin \alpha}{\cos \alpha} = \frac{-a_{BA}^{c} \sin \alpha}{\cos \alpha}$$
(3.22)

$$\boldsymbol{\varepsilon}_{AB} = \frac{a_{BA}^r}{AB} \tag{3.23}$$

Coordinates $P_{AB}(x_{B_i}, y_{A_i})$

Coordinates $C(x_{C_i}, y_{C_i})$

$$\frac{a}{AB} = \frac{x_{C_i}}{x_{B_i}}, \frac{b}{AB} = \frac{y_{C_i}}{y_{A_i}}$$
(3.24)

$$x_{C_{i}} = \frac{a}{AB} * x_{B_{i}}, y_{C_{i}} = \frac{b}{AB} * y_{A_{i}}$$
(3.25)

$$CP_{AB} = \sqrt{x_{B_i}^2 + a^2 - 2(a * x_{B_i})\cos\alpha}$$
(3.26)

$$\boldsymbol{\nu}_{C} = \boldsymbol{\omega}_{AB} * CP_{AB} = \boldsymbol{\omega}_{AB} * \sqrt{x_{B_{i}}^{2} + a^{2} - 2 * (a * x_{B_{i}}) * \cos \alpha}$$
(3.27)

$$\varphi = \operatorname{atan} \frac{y_{C_i}}{x_{C_i}} \tag{3.28}$$

The acceleration of point *C* is determined by equation (3.10): $\overrightarrow{a_{C}} = \overrightarrow{a_{A}} + \overrightarrow{a_{CA}^{r}} + \overrightarrow{a_{CA}^{c}}$

$$\boldsymbol{a}_{CA}^{c} = \boldsymbol{\omega}_{AB}^{2} * AC = \boldsymbol{\omega}_{AB}^{2} * a \tag{3.29}$$

$$\boldsymbol{a}_{CA}^{r} = \boldsymbol{\varepsilon}_{AB} * AC = \boldsymbol{\varepsilon}_{AB} * a \tag{3.30}$$

$$\begin{cases} a_{Cx} = a_{Ax} + a_{CAx}^{r} + a_{CAx}^{c} \\ a_{Cy} = a_{Ay} + a_{CAy}^{r} + a_{CAy}^{c} \end{cases}$$
(3.31)

$$\begin{cases} \boldsymbol{a}_{Cx} = 0 + \boldsymbol{a}_{CA}^r * \sin \alpha + \boldsymbol{a}_{CA}^c * \cos \alpha \\ \boldsymbol{a}_{Cy} = 0 + \boldsymbol{a}_{CA}^r * \cos \alpha + \boldsymbol{a}_{CA}^c * \sin \alpha \end{cases}$$
(3.32)

$$\boldsymbol{a}_{C} = \sqrt{\boldsymbol{a}_{Cx}^{2} + \boldsymbol{a}_{Cy}^{2}} \tag{3.33}$$

Uniformly accelerated motion

Given: point *C* divides *AB* into segments *a* and *b*, *A* moves uniformly accelerated from $O \rightarrow Y$, $A(0, y_A) B(x_B, 0)$, initial A(0, 0), B(AB, 0), $\mathbf{a}_{A_i} = const$, $\mathbf{v}_{A_0} = 0$.

Find: y_{A_i} , (x_{C_i}, y_{C_i}) , v_{C_i} , a_{C_i} , φ_i

Solution

$$v_{A_i} = \frac{a_A * i^2}{2}; \ i = 1 \dots n = \frac{T}{4}$$
 (3.34)

$$AB = v_{An} = \frac{a_{A} * n^2}{2}$$
(3.35)

$$\boldsymbol{a}_{A_i} = \boldsymbol{a}_A = \frac{2AB}{n^2} \tag{3.36}$$

Coordinates $A(0, y_{A_i})$

$$y_{A_i} = \frac{a_A * i^2}{2} \tag{3.37}$$

Further, according to equations (3.4) - (3.14)

Coordinates $B(x_{B_i}, 0)$:

$$x_{B_i} = \sqrt{AB^2 - y_{A_i}^2} \tag{3.38}$$

Coordinates $C(x_{C_i}, y_{C_i})$:

$$\frac{a}{AB} = \frac{x_{C_i}}{x_{B_i}}, \frac{b}{AB} = \frac{y_{C_i}}{y_{A_i}}$$
(3.39)

$$x_{C_i} = \frac{a}{AB} * x_{B_i}, y_{C_i} = \frac{b}{AB} * y_{A_i}$$
(3.40)

$$\boldsymbol{\omega}_{AB} = \frac{\boldsymbol{v}_{A_i}}{A \boldsymbol{P}_{AB}} = \frac{\boldsymbol{v}_{A_i}}{\boldsymbol{x}_{B_i}} \tag{3.41}$$

$$\boldsymbol{a}_{BA}^{c} = \boldsymbol{\omega}_{AB}^{2} * AB \tag{3.42}$$

$$\boldsymbol{a}_{BA}^r = \boldsymbol{\varepsilon}_{AB} * BA \tag{3.43}$$

The vector $\overrightarrow{a_{BA}^r}$ is located perpendicular to the ruler *AB*, its direction is unknown, since the direction of the angular acceleration $\boldsymbol{\varepsilon}_{AB}$ is unknown.

We project the vector equation (3.4) on the coordinate axis:

$$\begin{cases} \boldsymbol{a}_{Bx} = \boldsymbol{a}_{BA}^c * \cos \alpha + \boldsymbol{a}_{BA}^r * \sin \alpha \\ 0 = \boldsymbol{a}_{Ay} + \boldsymbol{a}_{BA}^c * \sin \alpha + \boldsymbol{a}_{BA}^r * \cos \alpha \end{cases}$$
(3.44)

Solving the resulting equations, we find a_B :

$$\boldsymbol{a}_{BA}^{r} = \frac{-\boldsymbol{a}_{Ay} - \boldsymbol{a}_{BA}^{c} * \sin \alpha}{\cos \alpha} \tag{3.45}$$

$$\boldsymbol{\varepsilon}_{AB} = \frac{a_{BA}^r}{AB} \tag{3.46}$$

Equation (10) can be represented in projections on the axes Ax and Ay:

$$\begin{cases} a_{Cx} = a_{Ax} + a_{CAx}^{r} + a_{CAx}^{c} \\ a_{Cy} = a_{Ay} + a_{CAy}^{r} + a_{CAy}^{c} \end{cases}$$
(3.49)

$$\begin{cases} \boldsymbol{a}_{Cx} = 0 + \boldsymbol{a}_{CA}^r * \sin \alpha + \boldsymbol{a}_{CA}^c * \cos \alpha \\ \boldsymbol{a}_{Cy} = \boldsymbol{a}_A + \boldsymbol{a}_{CA}^r * \cos \alpha + \boldsymbol{a}_{CA}^c * \sin \alpha \end{cases}$$
(3.50)

$$a_{C} = \sqrt{a_{Cx}^{2} + a_{Cy}^{2}} \tag{3.51}$$

Elliptical

The movement of the points of the ruler along the ellipse relative to the center,

$$\ddot{\varphi} = \frac{2 * e^2 * \cos(\varphi) * \sin(\varphi) * \dot{\varphi}^2}{1 - e^2 * \cos(\varphi)^2}$$
(2.13)

Given: point *C* divides *AB* into segments *a* and *b*, *A* moves elliptically according to the formula (2.13), from $O \rightarrow Y$, $A(0, y_A)$, $B(x_B, 0)$, initial A(0, 0), B(AB, 0), $v_{A_0} = 0$.

Find: $y_{A_i}, x_{C_i}, y_{C_i}, v_{C_i}, a_{C_i}$.

Solution

Equation (2.13) calculates φ_i , x_{C_i} , y_{C_i}

$$\alpha = \arcsin \frac{y_C}{b} \tag{3.52}$$

$$\beta = \frac{\pi}{2} - \varphi_i \tag{3.53}$$

$$\gamma = \arcsin\left(\frac{r_i * \sin\beta}{a}\right) \tag{3.54}$$

$$\psi = \pi - \gamma - \beta \tag{3.55}$$

$$y_{A_i} = \frac{y_{C_i} + a \cdot \sin \alpha}{b} \tag{3.60}$$

$$v_{A_i} = y_{A_i} - y_{A_{i-1}} \tag{3.61}$$

$$a_{A_i} = v_{A_i} - v_{A_{i-1}} \tag{3.62}$$

Further, according to equations (3.4) - (3.14)

Coordinates $B(x_{B_i}, 0)$:

$$x_{B_i} = \sqrt{AB^2 - y_{A_i}^2} \tag{3.63}$$

Find the coordinates $C(x_{C_i}, y_{C_i})$ again:

$$\frac{a}{AB} = \frac{x_{Ci}}{x_{Bi}}, \frac{b}{AB} = \frac{y_{Ci}}{y_{Ai}}$$
(3.64)

$$x_{C_i} = \frac{a}{AB} * x_{B_i}, y_{C_i} = \frac{b}{AB} * y_{A_i}$$
(3.65)

$$\boldsymbol{\omega}_{AB} = \frac{\boldsymbol{v}_{Ai}}{A\boldsymbol{P}_{AB}} = \frac{\boldsymbol{v}_{Ai}}{\boldsymbol{x}_{Bi}} \tag{3.66}$$

$$\boldsymbol{a}_{BA}^c = \boldsymbol{\omega}_{AB}^2 * AB \tag{3.67}$$

$$\boldsymbol{a}_{BA}^r = \boldsymbol{\varepsilon}_{AB} * BA \tag{3.68}$$

The vector $\overrightarrow{a_{BA}^r}$ is located perpendicular to the ruler *AB*, its direction is unknown, since the direction of the angular acceleration ε_{AB} is unknown.

We project the vector equation (3.4) on the coordinate axis:

$$\begin{cases} \boldsymbol{a}_{Bx} = \boldsymbol{a}_{BA}^{c} * \cos \alpha + \boldsymbol{a}_{BA}^{r} * \sin \alpha \\ 0 = \boldsymbol{a}_{Ay} + \boldsymbol{a}_{BA}^{c} * \sin \alpha + \boldsymbol{a}_{BA}^{r} * \cos \alpha \end{cases}$$
(3.69)

Solving the resulting equations, we find \boldsymbol{a}_{B} ,

$$\boldsymbol{a}_{BA}^{r} = \frac{-\boldsymbol{a}_{AY} - \boldsymbol{a}_{BA}^{r} \sin \alpha}{\cos \alpha} \tag{3.70}$$

$$\boldsymbol{\varepsilon}_{AB} = \frac{a_{BA}^r}{AB} \tag{3.71}$$

The acceleration of point C is determined by equation (3.10): $\overrightarrow{a_C} = \overrightarrow{a_A} + \overrightarrow{a_{CA}^r} + \overrightarrow{a_{CA}^c}$

$$\boldsymbol{a}_{CA}^{c} = \boldsymbol{\omega}_{AB}^{2} * AC = \boldsymbol{\omega}_{AB}^{2} * a \tag{3.72}$$

$$\boldsymbol{a}_{CA}^{r} = \boldsymbol{\varepsilon}_{AB} * AC = \boldsymbol{\varepsilon}_{AB} * a \tag{3.73}$$

Equation (3.10) can be represented in projections on the axes Ax and Ay:

$$\begin{cases} a_{Cx} = a_{Ax} + a_{CAx}^{r} + a_{CAx}^{c} \\ a_{Cy} = a_{Ay} + a_{CAy}^{r} + a_{CAy}^{c} \end{cases}$$
(3.74)

$$\begin{cases} \boldsymbol{a}_{Cx} = 0 + \boldsymbol{a}_{CA}^{r} * \sin \alpha + \boldsymbol{a}_{CA}^{c} * \cos \alpha \\ \boldsymbol{a}_{Cy} = \boldsymbol{a}_{A} + \boldsymbol{a}_{CA}^{r} * \cos \alpha + \boldsymbol{a}_{CA}^{c} * \sin \alpha \end{cases}$$
(3.75)

$$a_{C} = \sqrt{a_{Cx}^{2} + a_{Cy}^{2}} \tag{3.76}$$

The obtained motion parameters allow checking the fulfillment of Kepler's laws.

Kepler's second law

Uniform movement

```
Enter char =

if char = "y" then the source data is specified:

y

a = 0.500; b = 0.450; T = 360

Second law of Kepler

Point bypasses 1/4 ellipse counterclockwise in 89 time units

Input 0 - uniform motion OR

Input 1-uniformly accelerated motion OR,

Input 2- elliptical motion):

0

UNIFORM MOTION

Set the start of the first sector <1,..., 89>: 3

Set the end of the first sector < 3< end < 89>:17

Set the start of the second sector <1,..., 89>: 55

first sector: angle(start)= 0.03; angle(end)= 0.17

second secto: angle(start)= 0.61; angle(end)= 0.82

intervals of time t1= 14; t2= 14

Area of the first sector: 0.1767757E-01

IERR: 0

Area of the second sector: 0.2445188E-01

IERR: 0
```

Figure 11

Uniformly accelerated motion

```
Enter char =

if char = "y" then the source data is specified:

y

a = 0.500; b = 0.450; T = 360

Second law of Kepler

Point bypasses 1/4 ellipse counterclockwise in 89 time units

Input 0 - uniform motion OR

Input 1-uniformly accelerated motion OR,

Input 2- elliptical motion):

1

UNIFORMLY ACCELERATEM MOTION

Set the start of the first sector <1,..., 89): 3

Set the end of the first sector <1,..., 89): 55

first sector: angle(start)= 0.00; angle(end)= 0.03

second secto: angle(start)= 0.35; angle(end)= 0.59

intervals of time t1= 14; t2= 14

Area of the first sector: 0.3933465E-02

IERR: 0

Area of the second sector: 0.2803914E-01

IERR: 0
```

Figure 12

Elliptical movement

```
Enter char =

if char = "y" then the source data is specified:

y

a = 0.500; b = 0.450; T = 360

Second law of Kepler

Point bypasses 1/4 ellipse counterclockwise in 89 time units

Input 0 - uniform motion OR

Input 1-uniformly accelerated motion OR,

Input 2- elliptical motion):

2

ELLEPTICAL MOTION

Set the start of the first sector (1,..., 89): 3

Set the end of the first sector (3< end < 89):17

Set the start of the second sector (1,..., 89): 55

first sector: angle(start)= 0.03; angle(end)= 0.25

second secto: angle(start)= 0.89; angle(end)= 1.15

intervals of time t1= 14; t2= 14

Area of the first sector: 0.2748870E-01

IERR: 0

Area of the second sector: 0.2748918E-01

IERR: 0
```



Equality of the areas of sectors is carried out only with elliptical motion.

Graphical results of moving a point along an ellipse at different speeds.

Uniform movement, Figure 14







Uniformly accelerated motion, Figure 15.

Uniformly motion: velocity and acceleration vectors, a = 0.50, b = 0.45, T = 80.00





Elliptical movement, Figure 16



Figure 16

VI. Kepler's laws as properties of kinematic equations of motion of a point along curves of the second order

The equations are solved by computer programs. The calculation results are compared with Kepler's laws. The uniqueness of the orbital velocity for the given parameters of the curve is noted. The orbital velocity is calculated from the kinematic equation and compared with the values of astronomical tables

The sector velocity modulus is a constant for a given ellipse.

$$|v_{\sigma}| = \frac{1}{2} |r| * |v| * \sin(r^{*}v) = const$$
(4.1)

If a point moves along a flat curve and its position is determined by the polar coordinates r and φ , then

$$|v_{\sigma}| = \frac{1}{2} |r|^2 \frac{d\varphi}{dt} = const$$
(4.2)

To illustrate the constancy of the sectoral velocity, a program was written to calculate the sector area in a given time interval. The program, TygeBraheKepler2_focal [A.1], calculates the parameters of the point movement according to equation (8) and shows the equality of the areas of the sectors at equal time intervals, fig. 17 - 19.

```
4-angular velocity, 5-polar radius, 6-linear velocity
Enter char =
if char = "y" then the source data is specified:
  a = 9.00 b = 7.00
          0.000000
                                1.00000005E-03
 dpi =
                          H =
Second law of Kepler H= 1.00000005E-03
The point bypasses the ellipse in 1/H time units (0 < H < 1), counterclockwise.
1/H =
        999
 Set the start of the first sector (i0=1,...,
                                                       999 ) i0 =
 Set the end of the first sector (i0\langle i1\langle 1/H) i1 =
999
 Set the start of the second sector (0<i02<1/H-i1+i0) i02 =
 angle(i0) 0.00; angle(i1) 6.28
 angle(i02)= 0.00; angle(i12) 6.28
Area of the first sector:
                            0.1975210E+03
IERR:
          Θ
Area of the second sector:
                              0.1975210E+03
IERR:
         0
```



Figure 17 shows the program test. The area of the ellipse is π ab. 3.14159*9*7 = 197.92017

Figure 18 shows equal time intervals at different points in the period.

```
4-angular velocity, 5-polar radius, 6-linear velocity
Enter char =
if char = "y" then the source data is specified:
  a = 9.00 b = 7.00
 dpi = 0.0000000
                         H =
                               1.0000005E-03
Second law of Kepler H= 1.00000005E-03
The point bypasses the ellipse in 1/H time units (0 < H < 1), counterclockwise.
1/H =
        999
 Set the start of the first sector (i0=1,..., 999 ) i0 =
22
 Set the end of the first sector (iO<i1<1/H) i1 =
333
 Set the start of the second sector (0<i02<1/H-i1+i0) i02 =
555
 angle(i0) 0.04; angle(i1) 0.81
 angle(i02)= 4.57; angle(i12)
                                 6.03
Area of the first sector:
                          0.6155315E+02
IERR:
         0
Area of the second sector:
                            0.6155347E+02
IERR:
         0
```



On fig. 19 added precession (dpi = 0.1) to the parameters of fig. eighteen.

```
9.00 b = 7.00
  a =
  (precession 0<=dpi<=pi/10) dpi=
Θ
 1
 0.10000000
 (period = 1, 0 < H < 1) H =
0.001
Second law of Kepler H= 1.00000005E-03
The point bypasses the ellipse in 1/H time units (0 < H < 1), counterclockwise.
1/H =
         999
 Set the start of the first sector (i0=1,...,
                                                       999 ) i0 =
22
 Set the end of the first sector (iO<i1<1/H) i1 =
333
 Set the start of the second sector (0<i02<1/H-i1+i0) i02 =
555
 angle(i0) 0.04; angle(i1) 0.89
 angle(i02)= 4.94; angle(i12)
                                  6.12
Area of the first sector:
                            0.6487499E+02
IERR:
         0
Area of the second sector:
                              0.6487521E+02
IERR:
         0
```

Figure 19

Kepler's third law

At perihelion and aphelion, $sin(\varphi) = 0$, so the acceleration at these points is zero, and the modulo velocity difference is a constant:

$$v_p - v_a = \delta \tag{4.3}$$

Sector velocity according to the law of conservation of momentum is a constant value:

$$\boldsymbol{v}_s = 1/2\,\boldsymbol{v}\mathbf{r} \tag{4.4}$$

Let us express the sector velocity modulo the linear velocity.

Since
$$sin(v_p^{\wedge} r_p) = sin(v_a^{\wedge} r_a)$$
 1, then

$$v_s = 1/2 v_p r_p = 1/2 r_p (v_a + \delta)$$
(4.5)

$$v_s = 1/2 \, v_a r_a \tag{4.6}$$

$$1/2r_p(v_a+\delta) = 1/2r_av_a \tag{4.7}$$

$$v_a = \frac{r_p \delta}{r_a - r_p} \tag{4.8}$$

We substitute (4.8) into (4.6):

$$v_s = \frac{\delta r_p r_a}{2(r_a - r_p)} \tag{4.9}$$

Calculate the area of the ellipse. One side:

$$S_{ellipse} = \pi ab \tag{4.10}$$

where a is the length of the major semiaxis, b is the length of the minor semiaxis of the orbit.

On the other hand

$$S_{ellipse} = v_s T = T \frac{\delta r_p r_a}{2(r_a - r_p)}$$
(4.11)

Consequently,

$$T\frac{\delta r_p r_a}{2(r_a - r_p)} = \pi ab \tag{4.12}$$

For further transformations, we use the geometric properties of the ellipse. We have ratios: $r_a - r_p = 2c, c = ae, r_p r_a = a^2 - c^2 = b^2$.

Substitute into (4.12):

$$T\frac{\delta b^2}{4ae} = \pi ab \tag{4.13}$$

$$T\frac{\delta b}{a^2 e} = 4\pi ; \quad \text{где } T = 1; \tag{4.14}$$

$$\frac{\delta b}{4\pi a^2 e} = 1 \tag{4.15}$$

Kepler's third law:
$$\frac{T^2}{a^3} = 1$$
 (4.16)

$$\frac{\delta b}{4\pi a^2 e} = \frac{T^2}{a^3}; \frac{\delta b}{4\pi e} = \frac{T^2}{a}; T = \frac{1}{2}\sqrt{\frac{\delta ba}{\pi e}} = \frac{1}{2}\sqrt{\frac{(v_p - v_a)ba}{\pi e}}$$
(4.17)

The program Movement of a mat point along an ellipse [A.2], using formulas (4.16) and (4.17), calculates the periods. $\delta = v_p - v_a [au/planet year]$

```
The differential equation of the second order curves
with respect to the focus is calculated.
The data table is displayed in the file ellpi.txt.
Table columns:1-number, 2 - time, 3 - angle,
4-angular speed, 5-polar radius, 6-linear speed
7-angular acceleration, 8-linear acceleration
 Enter 0 or 1 or 2 or 3 or 4
 0 - enter a, b. Select planet 1 - Mercury, 2 - Venus, 3 - Earth, 4 - Mars:
 a =
9
 b =
  a = 9.00 b = 7.00
 orbital points (N):
                              999
                               27.000000
 period(Kepler3 sqrt(a××3)=
 period(sqrt(((v1-v2)*b*a)/(pi*ex))/2) =
                                            26.999981
PAUSE
To resume execution, type go. Other input will terminate the job.
```



```
The differential equation of the second order curves
with respect to the focus is calculated.
The data table is displayed in the file ellpi.txt.
Table columns:1-number, 2 - time, 3 - angle,
4-angular speed, 5-polar radius, 6-linear speed
7-angular acceleration, 8-linear acceleration
Enter 0 or 1 or 2 or 3 or 4
0 - enter a, b. Select planet 1 - Mercury, 2 -Uenus, 3 - Earth, 4 - Mars:
1
a = 0.39 b = 0.38
orbital points (N): 999
period(Kepler3 sqrt(a×3)= 0.24084271
period(sqrt(((u1-u2)×b×a)/(pi×ex))/2) = 0.24084280
PAUSE
To resume execution, type go. Other input will terminate the job.
```



2. Differential equation of motion of a point along an ellipse with respect to the center

Let's move the origin of coordinates to the center of the ellipse, Fig. 22. The radius function (2.7) will change.



M - material point. *Q* is a generalized force acting on a point. *O* - center. *v* - linear speed of the point. $\varphi(t)$ is the angle between the *X*-axis and the point, counterclockwise.

Figure 22

Kepler's second law

The TygeBraheKepler2_center [A.1] program calculates the parameters of the point movement according to equations (2.7 - 2.13), and shows the equality of the areas of the sectors at equal time intervals. Figures 23–25.

```
4-angular velocity, 5-polar radius, 6-linear velocity
Enter char =
if char = "y" then the source data is specified:
        9.00 b = 7.00
   a
    0.0000000
                                1.00000005E-03
 dpi
                          H =
Second law of Kepler H= 1.00000005E-03
The point bypasses the ellipse in 1/H time units (0 < H < 1), counterclockwise.
 /H =
         999
 Set the start of the first sector (i0=1,...,
                                                        999 ) i0 =
 Set the end of the first sector (iO<i1<1/H) i1 =
999
  Set the start of the second sector (0<i02<1/H-i1+i0) i02 =
             0.00; angle(i1)
 angle(i0)
                              6.28
 angle(i02)= 0.00; angle(i12)
                                  6.28
Area of the first sector:
                             0.1976214E+03
IERR:
          Θ
Area of the second sector:
                              0.1976214E+03
IERR:
          Θ
```



Figure 23 shows the program test. The area of the ellipse is π ab. 2*3.14159*9*7 = 197.92017

```
4-angular velocity, 5-polar radius, 6-linear velocity
Enter char =
if char = "y" then the source data is specified:
  a = 9.00 b = 7.00
 dpi =
          0.0000000
                         H =
                               1.00000005E-03
Second law of Kepler H= 1.00000005E-03
The point bypasses the ellipse in 1/H time units (0 < H < 1), counterclockwise.
1/H =
        999
 Set the start of the first sector (i0=1,...,
                                                     999 ) i0 =
22
Set the end of the first sector (i0<i1<1/H) i1 =
333
 Set the start of the second sector (0<i02<1/H-i1+i0) i02 =
555
 angle(i0) 0.10; angle(i1) 2.20
 angle(i02)= 3.41; angle(i12)
                                 5.56
Area of the first sector:
                           0.6155317E+02
IERR:
         0
Area of the second sector:
                             0.6155319E+02
IERR:
         0
```

Figure 24

On fig. 24 equal time intervals are given at different moments of the period.

```
semiminor axis b =
   a = 9.00 b = 7.00
  (precession 0<=dpi<=pi/10) dpi=
0.1
 0.10000000
  (period = 1, 0 < H < 1) H =
0.001
Second law of Kepler H= 1.00000005E-03
The point bypasses the ellipse in 1/H time units (0 < H < 1), counterclockwise.
1/H =
         999
 Set the start of the first sector (i0=1,...,
                                                      999 ) i0 =
22
 Set the end of the first sector (i0<i1<1/H) i1 =
333
 Set the start of the second sector (0<i02<1/H-i1+i0) i02 =
555
 angle(i0) 0.10; angle(i1) 2.25
 angle(i02)= 3.47; angle(i12) 5.67
Area of the first sector: 0.6280998E+02
IERR:
         ß
Area of the second sector:
                             0.6280998E+02
IERR:
         0
```

Figure 25

On fig. 25 added precession (dpi = 0.1) to the parameters of fig. 23.

Kepler's third law

The program Movement of a mat point along an ellipse center [A.2], using formulas (4.16 – 4.17), calculates the periods. $\delta = vI - v2$ [*au/planet year*].

In Figures 25 - 27 we see that with an increase in the eccentricity, the difference between the periods increases.

```
The differential equation of the second order curves
with respect to the focus is calculated.
The data table is displayed in the file ellpi.txt.
Table columns:1-number, 2 - time, 3 - angle,
4-angular speed, 5-polar radius, 6-linear speed
7-angular acceleration, 8-linear acceleration
 Enter 0 or 1 or 2 or 3 or 4
 0 - enter a, b. Select planet 1 - Mercury, 2 -Venus, 3 - Earth, 4 - Mars:
 a =
9
 b =
  a = 9.00 b = 7.00
 orbital points (N):
                             999
 period(Kepler3 sqrt(axx3)= 27.000000
 period(sqrt(((u1-u2)*b*a)/(pi*ex))/2) = 21.000002
PAUSE
To resume execution, type go. Other input will terminate the job.
```



```
The differential equation of the second order curves

with respect to the focus is calculated.

The data table is displayed in the file ellpi.txt.

Table columns:1-number, 2 - time, 3 - angle,

4-angular speed, 5-polar radius, 6-linear speed

7-angular acceleration, 8-linear acceleration

Enter 0 or 1 or 2 or 3 or 4

0 - enter a, b. Select planet 1 - Mercury, 2 -Venus, 3 - Earth, 4 - Mars:

1

a = 0.39 b = 0.38

orbital points (N): 999

period(Kepler3 sqrt(a××3)= 0.24084271

period(sqrt(((v1-v2)×b×a)/(pi×ex))/2) = 0.23569536

PAUSE

To resume execution, type go. Other input will terminate the job.
```

```
The differential equation of the second order curves
with respect to the focus is calculated.
The data table is displayed in the file ellpi.txt.
Table columns:1-number, 2 - time, 3 - angle,
4-angular speed, 5-polar radius, 6-linear speed
7-angular acceleration, 8-linear acceleration
 Enter 0 or 1 or 2 or 3 or 4
 0 - enter a, b. Select planet 1 - Mercury, 2 -Venus, 3 - Earth, 4 - Mars:
       0.73 b =
                  0.73
  a =
 orbital points (N):
                              999
 period(Kepler3 sqrt(a**3)=
                             0.62144679
 period(sqrt(((u1-u2)*b*a)/(pi*ex))/2) = 0.62116992
PAUSE
To resume execution, type go. Other input will terminate the job.
```



Conclusion

The kinematic equation (1.10) accurately describes the motion along ideal second-order curves. The real orbits of cosmic bodies have deviations from the ideal curve: precession, periodic asymmetry of the lengths of the radii, and other types of deviation.

Equation (1.10) and the center of mass theorem make it possible to simulate the motion of three or more bodies along second-order curves. Example [A.5], fig. 29, 30.



Figure 29



Figure 30

The kinematic equation (2.13) is applicable for modeling streamlines of liquid and gas particles.

The article used materials from textbooks on mechanics.

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Applications

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