On eikonal and black hole entropy

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We formulate the eikonal equation in (3+1)-dimensional spherically symmetric curved space-time using Clebsch variables. We assume that the mass is the mass of a black hole and it is related to a black hole entropy through its area.

Keywords: geometrical optics, eikonal equation, null geodesic, Schwarzschild metric, entropy, black hole area.

The static spherically symmetric curved space-time (gravitational field) described by the Schwarzschild metric¹, an interval, can be written as

$$ds^{2} = -\left(1 - \frac{r_{s}}{r}\right)c^{2}dt^{2} + \left(1 - \frac{r_{s}}{r}\right)^{-1}dr^{2} + r^{2} d\theta^{2} + r^{2} \sin^{2}\theta d\phi$$
(1)

where r is the spatial (radial) coordinate, the distance from the centre of a mass of the massive body (such as a black hole, a massive star) to the points (these points have the same distance from the centre) where the Schwarzschid metric is being evaluated, and

$$r_s = \frac{2GM}{c^2} \tag{2}$$

is the Schwarzschild radius, M is the mass of the spherically symmetric body, G is the gravitational constant, and c is the speed of light in flat space-time. Equation (1) is also known as the Schwarzschild solution² (to Einstein field equation). The Schwarzschild solution holds outside the surface of the massive body that is producing the gravitational field, where there is no matter².

In our previous work³, we found that the eikonal equation in (3+1)-dimensional spherically symmetric curved space-time can be written below

$$\left|\partial_{\mu}\left\{\frac{c}{f_{\theta}}\int f \ \partial_{\nu}q \ f^* \ dx^{\nu} + ct\right\}\right| = \left(1 - \frac{2GM}{c^2r}\right)^{-1}(3)$$

where $f \partial_{\nu} q = \vec{A}_{\nu}$ is the U(1) gauge potential, f, q are Clebsch variables, f^* is the conjugate complex of f, and f_{θ} is angular frequency. Eq.(3) shows explicitly the relation between the mass of the spherically symmetric body that is producing the spherically symmetric gravitational field and the U(1) gauge potential which represents the existence of the light ray.

The relation between a black hole entropy and its area can be written as $\!\!\!^4$

$$S_{bh} = \frac{\eta kA}{l_p^2} \tag{4}$$

where S_{bh} is the entropy of a black hole, A is the area of a black hole, k is the Boltzmann's constant, η is a constant

number of order unity, l_p is the Planck length

$$l_p = \left(\frac{\hbar G}{c^3}\right)^{1/2} \tag{5}$$

 $\hbar = h/2\pi$, h is the Planck constant.

The surface area of a spherically symmetric black hole can be written as

$$A = 4\pi r_s^2 \tag{6}$$

By substituting eq.(2) to (6) we obtain

$$A = 4\pi \left(\frac{2GM}{c^2}\right)^2 \tag{7}$$

By substituting eq.(7) into (4), we obtain

$$S_{bh} = \frac{4\pi\eta k}{l_p^2} \left(\frac{2GM}{c^2}\right)^2 \tag{8}$$

By substituting (5) into (8), we obtain

$$S_{bh} = \frac{16\pi\eta k}{\hbar c} \ GM^2 \tag{9}$$

or

$$M = \sqrt{\frac{\hbar c}{16\pi\eta kG} S_{bh}} \tag{10}$$

By substituting (10) into (2), we obtain

$$r_s = l_p \sqrt{\frac{S_{bh}}{4\pi\eta k}} \tag{11}$$

Eq.(11) relates the Schwarzschild radius to a black hole entropy.

If we replace eq.(2) with (11) then we can write eq.(3) as

$$\left| \partial_{\mu} \left\{ \frac{c}{f_{\theta}} \int f \ \partial_{\nu} q \ f^* \ dx^{\nu} + ct \right\} \right|$$

= $\left(1 - \frac{l_p}{r} \sqrt{\frac{S_{bh}}{4\pi\eta k}} \right)^{-1}$ (12)

Eq.(12) is the eikonal equation that describes the light ray propagation in static spherically symmetric spacetime with the distance r from the centre related to the entropy of a black hole. It relates the U(1) gauge potential to the entropy of a black hole. Does it mean that the U(1) gauge potential shares similar properties with the entropy?

If we take $r = r_s$ then eqs.(1), (12) become

$$ds^2 = \infty \tag{13}$$

and

$$\left|\partial_{\mu}\left\{\frac{c}{f_{\theta}}\int f \ \partial_{\nu}q \ f^{*} \ dx^{\nu} + ct\right\}\right| = \infty \qquad (14)$$

respectively. Eqs.(13), (14) show us that there exist the

singularity.

Thank to Juwita Armilia and Aliya Syauqina Hadi for much love. Al Fatihah for his Ibunda and Ayahanda. May Allah bless them Jannatul Firdaus.

This research is supported fully by self-funding.

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