## Unfringing Interference of Cross-Polarized Slits

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I project a diagonally-polarized (D) monochromatic complex plane-wave<sup>(1)</sup> electric field onto cross-polarized slits. (1) What is the instantaneous behavior of the electric field versus diffraction angle in the far field<sup>(2)</sup>? (2) The time-parametric vector sum of crossedpolarized slits' electric fields in a fixed plane about normal to the slits' rays traces Lissajous curves<sup>(3)</sup> particular to the diffraction<sup>(4)</sup> angle. (3) These uniform-brightness, diffraction-angle-dependent lines, circles, and ellipses constructively constitute unfringing interference of transverse field undulations<sup>(5)</sup> summed in parallel and falsify the first Fresnel-Arago Law<sup>(6,7)</sup> since interference (the 3-d vector sum)<sup>(1)</sup>, somewhat ironically, predicts the outcome. This result essentially retro-extends Fresnel optics (transverse vibrations) upon the Young's Double-Slit via the (parallel and direct) lineages of Maxwell, Heaviside, and Poynting, supporting Young's original assertion of light as a wave<sup>(8.9)</sup>.

In the diffraction plane beyond linearly crossed-polarized optical slits, assuming a plane-wave E-field or wavefunction<sup>(11)</sup>, the parametric vector sum, neglecting the diffraction envelope, is<sup>(10)</sup>:

$$\psi = \begin{pmatrix} e^{i\phi_1} & 0 \\ 0 & e^{i\phi_2} \end{pmatrix} \begin{bmatrix} E_{0x} \\ E_{0y} \end{bmatrix} e^{-i\omega t} \text{ so that } \Re(\psi) = \begin{bmatrix} E_{0x}\cos(-\omega t) \\ E_{0y}\cos(-\omega t + \Delta\phi) \end{bmatrix}$$
(1)

(where  $\phi_1, \phi_2 \equiv \phi_1(\theta, +d), \phi_2(\theta, -d), \Delta \phi = \phi_2 - \phi_1, \omega = 2\pi\lambda/c$ ,  $\theta$  is the diffraction angle, and d is the slit spacing/2). This vector sum traces Lissajous patterns<sup>(3)</sup> which are cyclic versus  $\theta$ . The time-average of

$$\vec{S} = \frac{1}{\mu_0} (\Re \vec{E}) \times (\Re \vec{B}) = \hat{k} \sqrt{\frac{\epsilon_0}{\mu_0}} [E_{0x}^2 \cos^2(\omega t) + E_{0y}^2 \cos^2(\omega t - \Delta \phi)]$$
(2)

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(the instantaneous real irradiance<sup>(12, 13)</sup>), is independent of  $\Delta \phi(\theta)$ . (1) cycles in polarization and

(2) in pulsed versus constant irradiance as functions of  $\Delta \phi(\theta)$ (e.g.,  $\Delta \phi$ 's of 0 and  $\pi$  contain D/A[nti]D with 2x-peak power pulses at  $2\nu$ , and  $\Delta \phi$ 's of  $\pi/2$  and  $3\pi/2$  contain L[eft]C[ircular-polarization]/R[ight]C with 1xconstant power.) For  $E_{0y} = E_{0z}$ , (fixed  $\theta$ ) sums constructing RC/LC change in direction only and sums constructing D/AD change in magnitude only

with t. Comparing the present to parallel-polarized<sup>(8.9)</sup> (fringing) slits gives the two-slit

geometry information encoded in polarization (and time-domain irradiance) in the former and in brightness (fringes) in the latter.

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