

Mapping Penrose-Rindler Null Tetrads to the Advanced and Retarded Wheeler-Feynman-Aharonov Destiny & History Null Tetrads

Jack Sarfatti

Abstract

This is a short mathematical note clarifying the use of Cramer's Transactional Interpretation in the Spinor Qubit Pre-Geometry of Wheeler's IT FROM BIT.

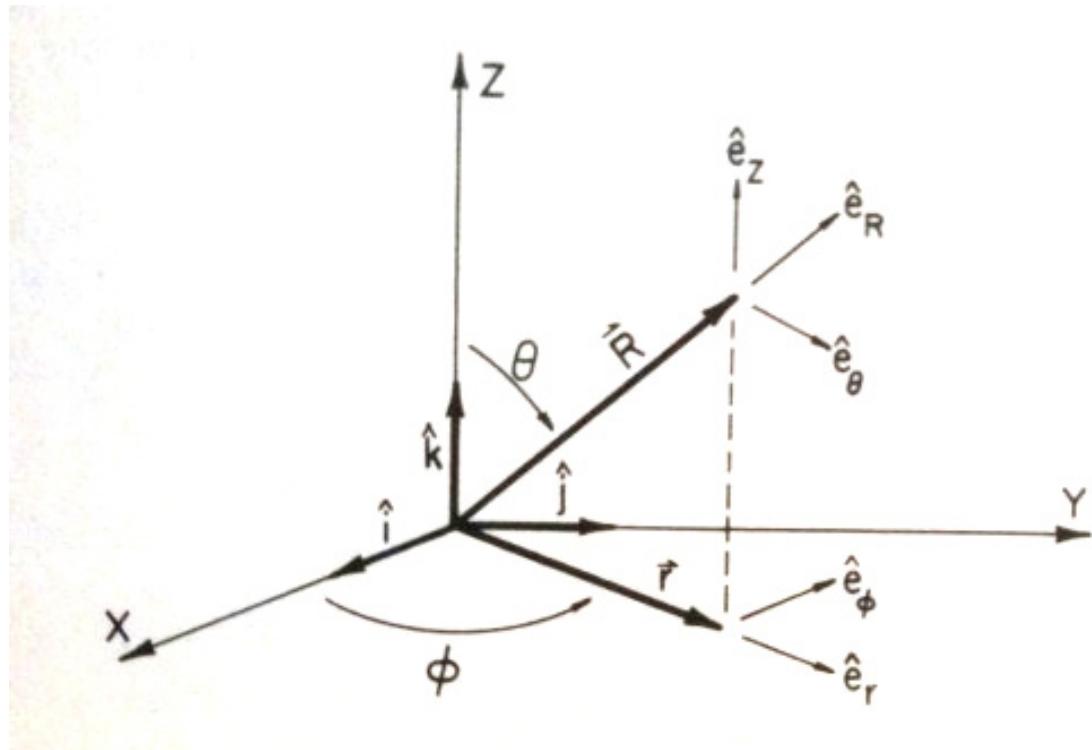
The mapping of the Penrose-Rindler Cartesian null tetrad of orthonormal base vectors to the Wheeler-Feynman tetrad of advanced and retarded 2D spherical wave front orthonormal base vectors is independent of the area/entropy of the wave fronts in the sense of the small 4D regions of LIFs in curved space-time.

$$\begin{aligned}\hat{e}_R &\equiv \hat{R} \equiv \sin \theta \cos \phi \hat{i} + \sin \theta \sin \phi \hat{j} + \cos \theta \hat{k} \\ \hat{e}_\theta &\equiv \hat{\theta} \equiv \cos \theta \cos \phi \hat{i} + \cos \theta \sin \phi \hat{j} + \sin \theta \hat{k} \\ \hat{e}_\phi &\equiv \hat{\phi} \equiv -\sin \phi \hat{i} + \cos \phi \hat{j}\end{aligned}\tag{1.1}$$

Using Roger Penrose's "abstract index notation" the Cartesian tetrad world vectors are mapped to the Einstein-Podolsky-Rosen entangled Bell pair quantum states of 2 qubit strings.

$$\begin{aligned}\hat{i} &\rightarrow \frac{1}{\sqrt{2}}(o^A i^{A'} + i^A o^{A'}) \\ \hat{j} &\rightarrow \frac{1}{\sqrt{2}}(o^A i^{A'} - i^A o^{A'}) \\ \hat{k} &\rightarrow \frac{1}{\sqrt{2}}(o^A o^{A'} - i^A i^{A'}) \\ \hat{t} &\rightarrow \frac{1}{\sqrt{2}}(o^A o^{A'} + i^A i^{A'})\end{aligned}\tag{1.2}$$

The mapping to the spherical wave advanced and retarded Wheeler-Feynman tetrads comes from the orthogonal transformation of spherical trigonometry. Hence



$$\begin{aligned}
\ell_{adv} &\equiv \frac{1}{\sqrt{2}}(\hat{t} + \hat{R}) \\
&\rightarrow \frac{1}{2} \left(\begin{array}{l} (o^A o^{A'} + i^A t^{A'}) + \sin \theta \cos \phi (o^A t^{A'} + i^A o^{A'}) \\ + \sin \theta \sin \phi (o^A t^{A'} - i^A o^{A'}) + \cos \theta (o^A o^{A'} - i^A t^{A'}) \end{array} \right) \\
&= \frac{1}{2} \left(\begin{array}{l} (1 + \cos \theta) o^A o^{A'} + (1 - \cos \theta) i^A t^{A'} + (\sin \theta \cos \phi + \sin \theta \sin \phi) o^A t^{A'} \\ + (\sin \theta \cos \phi - \sin \theta \sin \phi) i^A o^{A'} \end{array} \right) \\
n_{ret} &\equiv \frac{1}{\sqrt{2}}(\hat{t} - \hat{R}) \\
&\rightarrow \frac{1}{2} \left(\begin{array}{l} (o^A o^{A'} + i^A t^{A'}) - \sin \theta \cos \phi (o^A t^{A'} + i^A o^{A'}) \\ - \sin \theta \sin \phi (o^A t^{A'} - i^A o^{A'}) - \cos \theta (o^A o^{A'} - i^A t^{A'}) \end{array} \right) \\
&= \frac{1}{2} \left(\begin{array}{l} (1 - \cos \theta) o^A o^{A'} + (1 + \cos \theta) i^A t^{A'} \\ - (\sin \theta \cos \phi + \sin \theta \sin \phi) o^A t^{A'} - (\sin \theta \cos \phi - \sin \theta \sin \phi) i^A o^{A'} \end{array} \right) \\
m_{WF} &\equiv \frac{1}{\sqrt{2}}(\hat{\theta} + i\hat{\phi}) \\
&\rightarrow \frac{1}{2} \left(\begin{array}{l} \cos \theta \cos \phi (o^A t^{A'} + i^A o^{A'}) + \cos \theta \sin \phi (o^A t^{A'} - i^A o^{A'}) \\ + \sin \theta (o^A o^{A'} - i^A t^{A'}) + i(-\sin \phi (o^A t^{A'} + i^A o^{A'}) + \cos \phi (o^A t^{A'} - i^A o^{A'})) \end{array} \right) \\
&= \frac{1}{2} \left(\begin{array}{l} \sin \theta (o^A o^{A'} - i^A t^{A'}) + (\cos \theta \cos \phi + \cos \theta \sin \phi - i \sin \phi + i \cos \phi) o^A t^{A'} \\ + (\cos \theta \cos \phi - \cos \theta \sin \phi - i \sin \phi - i \cos \phi) i^A o^{A'} \end{array} \right)
\end{aligned} \tag{1.3}$$

Note in the limit $\theta \rightarrow 0$ we have

$$\begin{aligned}
\ell_{adv} &\equiv \frac{1}{\sqrt{2}}(\hat{t} + \hat{R}) \xrightarrow{\theta \rightarrow 0} o^A o^{A'} \\
n_{ret} &\equiv \frac{1}{\sqrt{2}}(\hat{t} - \hat{R}) \rightarrow i^A t^{A'} \\
m_{WF} &\equiv \frac{1}{\sqrt{2}}(\hat{\theta} + i\hat{\phi}) \\
&= \frac{1}{2} \left(\begin{array}{l} (\cos \phi + \sin \phi - i \sin \phi + i \cos \phi) o^A t^{A'} \\ + (\cos \phi - \sin \phi - i \sin \phi - i \cos \phi) i^A o^{A'} \end{array} \right)
\end{aligned} \tag{1.4}$$