De Broglie’s wavelength, elementary particles, the wavefunction and relativity

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Summary
This paper explores how the ring current model of an electron – or of matter-particles in general – relates to Louis de Broglie’s $\lambda = h/p$ relation and rephrases the theory in terms of the wavefunction as well as Schrödinger’s wave equation for an electron in free space. In the process, we argue that the latter is relativistically correct because the $\frac{1}{2}$ factor does not stem from the $\frac{1}{2}$ factor in the non-relativistic kinetic energy equation but from the concept of the (relativistic) effective mass of the pointlike charge inside of the electron, which it acquires from zipping around near or at lightspeed.

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De Broglie’s wavelength and the Compton radius

De Broglie’s ideas on the matter-wave are oft-quoted and are usually expressed in de Broglie’s $\lambda = h/p$ relation. However, there is remarkably little geometric or physical interpretation: what is that wavelength, exactly? The relation itself is easy enough to read: $\lambda$ goes to infinity as $p$ goes to zero. In contrast, for $p = mv$ going to $p = mc$, this length becomes the Compton wavelength $\lambda = h/p = h/mc$. This must mean something, obviously, but what exactly?

Mainstream theory does not answer this question because the Heisenberg interpretation of quantum mechanics essentially refuses to look into a geometric or physical interpretation of de Broglie’s relation and/or the underlying concept of the matter-wave or wavefunction which, lest we forget, must somehow represent the particle itself. In contrast, we will request the reader to think of the (elementary) wavefunction as representing a current ring.

To be precise, we request the reader to think of the (elementary) wavefunction $r = \psi = a \cdot e^{ik} a$ as representing the physical position of a pointlike elementary charge – pointlike but not dimensionless\(^1\) – moving at the speed of light around the center of its motion in a space that is defined by the electron’s Compton radius $a = h/mc$. This radius – which effectively doubles up as the amplitude of the wavefunction – can easily be derived from (1) Einstein’s mass-energy equivalence relation, (2) the Planck-Einstein relation, and (3) the formula for a tangential velocity, as shown below:

$$\begin{align*}
E &= mc^2 \\
E &= h\omega \\
c &= a\omega \iff a = \frac{c}{\omega} \iff \omega = \frac{c}{a}
\end{align*}$$

$$\begin{align*}
ma^2\omega^2 &= h\omega \implies m\frac{c^2}{\omega^2}\omega^2 = h\frac{c}{a} \iff a = \frac{h}{mc}
\end{align*}$$

This easy derivation\(^2\) already gives a more precise explanation of Prof. Dr. Patrick R. LeClair’s interpretation of the Compton wavelength as “the scale above which the particle can be localized in a particle-like sense”\(^3\), but we may usefully further elaborate the details by visualizing the model (Figure 1) and exploring how it fits de Broglie’s intuitions in regard to the matter-wave, which is what we set out

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\(^1\) The non-zero dimension of the elementary charge explains the small anomaly in the magnetic moment which is, therefore, not anomalous at all. For more details, see our paper on the electron model.

\(^2\) It is a derivation one can also use to derive a theoretical radius for the proton (or for any elementary particle, really). It works perfectly well for the muon, for example. However, for the proton, an additional assumption in regard to the proton’s angular momentum and magnetic moment is needed to ensure it fits the experimentally established radius. We shared the derivation with Prof. Dr. Randolph Pohl and the PRad team but we did not receive any substantial comments so far, except for the PRad spokesman (Prof. Dr. Ashot Gasparan) confirming the Standard Model does not have any explanation for the proton radius from first principles and, therefore, encouraging us to continue our theoretical research. In contrast, Prof. Dr. Randolph Pohl suggested the concise calculations come across as numerological only. We hope this paper might help to make him change his mind!

\(^3\) Prof. Dr. Patrick LeClair, Introduction to Modern Physics, Course Notes (PH253), 3 February 2019, p. 10.
to do in this paper.\textsuperscript{4}

Figure 1: The ring current model of an electron

Of course, the reader will, most likely, not be familiar with the ring current model or – using the term Erwin Schrödinger coined for it – the \textit{Zitterbewegung} model and we should, therefore, probably quote an unlikely authority on it so as to establish some early credentials\textsuperscript{5}:

“The variables [of Dirac’s wave equation] give rise to some rather unexpected phenomena concerning the motion of the electron. These have been fully worked out by Schrödinger. It is found that an electron which seems to us to be moving slowly, must actually have a very high frequency oscillatory motion of small amplitude superposed on the regular motion which appears to us. As a result of this oscillatory motion, the velocity of the electron at any time equals the velocity of light. This is a prediction which cannot be directly verified by experiment, since the frequency of the oscillatory motion is so high and its amplitude is so small. But one must believe in this consequence of the theory, since other consequences of the theory which are inseparably bound up with this one, such as the law of scattering of light by an electron, are confirmed by experiment.” (Paul A.M. Dirac, \textit{Theory of Electrons and Positrons}, Nobel Lecture, December 12, 1933)

Indeed, the dual radius of the electron (Thomson versus Compton radius) and the \textit{Zitterbewegung} model combine to explain the wave-particle duality of the electron and, therefore, diffraction and/or interference as well as Compton scattering itself. We will not dwell on these aspects of the ring current electron model because we have covered them in (too) lengthy papers before. Indeed, we will want to stay focused on the prime objective of this paper, which is a geometric or physical interpretation of the matter-wave.

Before we proceed, we must note the momentum of the pointlike charge – which we denote by $p$ in the illustration – must be distinguished from the momentum of the electron as a whole. The momentum of

\textsuperscript{4} We will analyze de Broglie’s views based on his paper for the 1927 Solvay Conference: Louis de Broglie, \textit{La Nouvelle Dynamique des Quanta} (the new quantum dynamics), 5th Solvay Conference, 1927. This paper has the advantage of being concise and complete at the same time. Indeed, its thirty pages were written well after the publication of his thesis on the new \textit{mécanique ondulatoire} (1924) but the presentation helped him to secure the necessary fame which would then lead to him getting the 1929 Nobel Prize for Physics.

\textsuperscript{5} For an overview of other eminent views, we refer to our paper on the 1921 and 1927 Solvay Conferences.
the pointlike charge will always be equal to \( p = mc \).\(^6\) The rest mass of the pointlike charge must, therefore, be zero. However, its velocity gives it an effective mass which one can calculate to be equal to \( m_{\text{eff}} = \frac{m_e}{2} \).\(^7\)

Let us now review de Broglie’s youthful and/or not-so-youthful intuitions.\(^8\)

**De Broglie’s dissipating wavepacket**

The ring current model of an electron incorporates the wavelike nature of an electron: the frequency of the oscillation is the frequency of the circulatory or oscillatory motion (Zitterbewegung) of the pointlike electric charge. Hence, the intuition of the young Louis de Broglie that an electron must have a frequency was, effectively, a stroke of genius. However, as the magnetic properties of an electron were, by then, not well established and this may explain why Louis de Broglie is either not aware of it or refuses to further build on it.\(^9\)

Let us have a closer look at his paper for the 1927 Solvay Conference, titled *La Nouvelle Dynamique des Quanta*, which we may translate as *The New Quantum Dynamics*. The logic is, by now, well known: we think of the particle as a wave packet composed of waves of slightly different frequencies \( \nu_i \).\(^10\) This leads to a necessary distinction between the group and phase velocities of the wave. The group velocity corresponds to the classical velocity \( v \) of the particle, which is often expressed as a fraction or relative velocity.

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\(^6\) We consciously use a vector notation to draw attention to the rather particular direction of \( p \) and \( c \): they must be analyzed as tangential vectors in this model.

\(^7\) We may refer to one of our previous papers here (Jean Louis Van Belle, *An Explanation of the Electron and Its Wavefunction*, 26 March 2020). The calculations involve a relativistically correct analysis of an oscillation in two independent directions: we effectively interpret circular motion as a two-dimensional oscillation. Such oscillation is, mathematically speaking, self-evident (Euler’s function is geometric sum of a sine and a cosine) but its physical interpretation is, obviously, not self-evident at all!

\(^8\) We must qualify this remark. Louis de Broglie was, obviously, a youthful genius but he does trace his own ideas on the matter-wave back to the time of writing of his PhD thesis, which is 1923-1924. Hence, he was 32 years old at the time, not nineteen! The reader will also know that, after WW II, Louis de Broglie would distance him from modern interpretations of his own theory and modern quantum physics by developing a realist interpretation of quantum physics itself. This interpretation would culminate in the *de Broglie-Bohm* theory of the pilot wave. We do not think there is any need for such alternative theories: we should just go back to where de Broglie went wrong and connect the dots.

\(^9\) The papers and interventions by Ernest Rutherford at the 1921 Conference do, however, highlight the magnetic dipole property of the electron. It should also be noted that Arthur Compton would highlight in his famous paper on Compton scattering, which he published in 1923 and was an active participant in the 1927 Conference itself. Louis de Broglie had extraordinary exposure to all of the new ideas, as his elder brother – Maurice Duc de Broglie – had already engaged him scientific secretary for the very first Solvay Conference in 1911, when Louis de Broglie was just 19 years. More historical research may reveal why Louis de Broglie did not connect the dots. As mentioned, he must have been very much aware of the limited but substantial knowledge on the magnetic moment of an electron as highlighted by Ernest Rutherford and others at the occasion of the 1921 Solvay Conference.

\(^10\) We invite the reader to check our exposé against de Broglie’s original 1927 paper in the *Solvay Conference proceedings*. We will try to stick closely to the symbols that are used in this paper, such as the \( \nu \) (\( \nu \)) symbol for the frequency.
velocity $\beta = \nu/c$.

The assumption is then that we know how the phase frequencies $\nu_i$ are related to wavelengths $\lambda_i$. This is modeled by a so-called dispersion relation, which is usually written in terms of the angular frequencies $\omega_i = 2\pi\nu_i$ and the wave numbers $k_i = 2\pi/\lambda_i$. The relation between the frequencies $\nu_i$ and the wavelengths $\lambda_i$ (or between angular frequencies $\omega_i$ and wavenumbers $k_i$) is referred to as the dispersion relation because it effectively determines if and how the wave packet will disperse or dissipate. Indeed, wave packets have a rather nasty property: they dissipate away. A real-life electron does not.

Prof. H. Pleijel, then Chairman of the Nobel Committee for Physics of the Royal Swedish Academy of Sciences, dutifully notes this rather inconvenient property in the ceremonial speech for the 1933 Nobel Prize, which was awarded to Heisenberg for nothing less than “the creation of quantum mechanics”:

“Matter is formed or represented by a great number of this kind of waves which have somewhat different velocities of propagation and such phase that they combine at the point in question. Such a system of waves forms a crest which propagates itself with quite a different velocity from that of its component waves, this velocity being the so-called group velocity. Such a wave crest represents a material point which is thus either formed by it or connected with it, and is called a wave packet. [...] As a result of this theory on is forced to the conclusion to conceive of matter as not being durable, or that it can have definite extension in space. The waves, which form the matter, travel, in fact, with different velocity and must, therefore, sooner or later separate. Matter changes form and extent in space. The picture which has been created, of matter being composed of unchangeable particles, must be modified.”

This should sound very familiar to you. However, it is, obviously, not true: real-life particles – electrons or atoms traveling in space – do not dissipate. Matter may change form and extent in space a little bit – such as, for example, when we are forcing them through one or two slits but not fundamentally so!

11 The concept of an angular frequency (radians per time unit) may be more familiar to you than the concept of a wavenumber (radians per distance unit). Both are related through the velocity of the wave (which is the velocity of the component wave here, so that is the phase velocity $v_p$):

$$v_p = \nu_i \cdot \lambda_i = \frac{\omega_i}{2\pi} \cdot \frac{2\pi}{k_i} = \frac{\omega_i}{k_i}$$

12 To be precise, Heisenberg got a postponed prize from 1932. Erwin Schrödinger and Paul A.M. Dirac jointly got the 1933 prize. Prof. Pleijel acknowledges all three in the introduction of his speech: “This year’s Nobel Prizes for Physics are dedicated to the new atomic physics. The prizes, which the Academy of Sciences has at its disposal, have namely been awarded to those men, Heisenberg, Schrödinger, and Dirac, who have created and developed the basic ideas of modern atomic physics.”

13 The wave-particle duality of the ring current model should easily explain single-electron diffraction and interference (the electromagnetic oscillation which keeps the charge swirling would necessarily interfere with itself when being forced through one or two slits), but we have not had the time to engage in detailed research here.

14 We will slightly nuance this statement later but we will not fundamentally alter it. We think of matter-particles as an electric charge in motion. Hence, as it acts on a charge, the nature of the centripetal force that keeps the particle together must be electromagnetic. Matter-particles, therefore, combine wave-particle duality. Of course, it makes a difference when this electromagnetic oscillation, and the electric charge, move through a slit or in free space. We will come back to this later. The point to note is: matter-particles do not dissipate. Feynman actually notes that at the very beginning of his Lectures on quantum mechanics, when describing the double-slit...
We should let this problem rest for a while. Let us briefly look at a related but somewhat different topic: the wave equation.

The wavefunction, the wave equation and Heisenberg’s uncertainty

With the benefit of hindsight, we now know the 1927 and later Solvay Conferences pretty much settled the battle for ideas in favor of the new physics. At the occasion of the 1948 Solvay Conference, it is only Paul Dirac who seriously challenges the approach based on perturbation theory which, at the occasion, is powerfully presented by Robert Oppenheimer. Dirac makes the following comment:

“All the infinities that are continually bothering us arise when we use a perturbation method, when we try to expand the solution of the wave equation as a power series in the electron charge. Suppose we look at the equations without using a perturbation method, then there is no reason to believe that infinities would occur. The problem, to solve the equations without using perturbation methods, is of course very difficult mathematically, but it can be done in some simple cases. For example, for a single electron by itself one can work out very easily the solutions without using perturbation methods and one gets solutions without infinities. I think it is true also for several electrons, and probably it is true generally: we would not get infinities if we solve the wave equations without using a perturbation method.”

However, Dirac is very much aware of the problem we mentioned above: the wavefunctions that come out as solutions dissipate away. Real-life electrons – any real-life matter-particle, really – do not do that. In fact, we refer to them as being particle-like because of their integrity—an integrity that is modeled by the Planck-Einstein relation in Louis de Broglie’s earliest papers too. Hence, Dirac immediately adds the following, recognizing the problem:

“If we look at the solutions which we obtain in this way, we meet another difficulty: namely we have the run-away electrons appearing. Most of the terms in our wave functions will correspond to electrons which are running away, in the sense we have discussed yesterday and cannot correspond to anything physical. Thus nearly all the terms in the wave functions have to be discarded, according to present ideas. Only a small part of the wave function has a physical meaning.”

In our interpretation of matter-particles, this small part of the wavefunction is, of course, the real electron, and it is the ring current or Zitterbewegung electron! It is the trivial solution that Schrödinger had found, and which Dirac mentioned very prominently in his 1933 Nobel Prize lecture. The other part of the solution(s) is (are), effectively, bizarre oscillations which Dirac here refers to as ‘run-away electrons’. With the benefit of hindsight, one wonders why Dirac did not see what we see now.

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experiment for electrons: “Electrons always arrive in identical lumps.”

15 This corresponds to wavefunctions dissipating away. The matter-particles they purport to describe obviously do not.

16 See pp. 282-283 of the report of the 1948 Solvay Conference, Discussion du rapport de Mr. Oppenheimer.

17 See the quote from Dirac’s 1933 Nobel Prize speech in this paper.

18 One of our correspondents wrote us this: “Remember these scientists did not have all that much to work with. Their experiments were imprecise – as measured by today’s standards – and tried to guess what is at work. Even
When discussing wave equations, it is always useful to try to imagine what they might be modeling. Indeed, if we try to imagine what the wavefunction might actually be, then we should also associate some (physical) meaning with the wave equation: what could it be? In physics, a wave equation – as opposed to the wavefunctions that are a solution to the wave equation (usually a second-order linear differential equation) – are used to model the properties of the medium through which the waves are traveling. If we are going to associate a physical meaning with the wavefunction, then we may want to assume the medium here would be the same medium as that through which electromagnetic waves are traveling, so that is the vacuum. Needless to say, we already have a set of wave equations here: those that come out of Maxwell’s equations! Should we expect contradictions here?

We hope not, of course—but then we cannot be sure. An obvious candidate for a wave equation for matter-waves in free space is Schrödinger equation without the Coulomb term:

\[
\frac{\partial \psi}{\partial t} = \frac{i \hbar}{2m_{\text{eff}}} \nabla^2 \psi = \frac{\partial \psi}{\partial t} = \frac{i \hbar}{m} \nabla^2 \psi
\]

What is \( m_{\text{eff}} \)? It is the concept of the effective mass of an electron which, in our ring current model, corresponds to the relativistic mass of the electric charge as it zippers around at lightspeed.

This is interesting! The question now is: are we talking one wave or many waves? A wave packet or the elementary wavefunction? Let us first make the analysis for one wave only, assuming that we can write \( \psi \) as some elementary wavefunction \( \psi = a e^{i\theta} = a e^{i(kx - \omega t)} \). Now, two complex numbers \( a + i b \) and \( c + i d \) are equal if, and only if, their real and imaginary parts are the same, and the \( \partial \psi / \partial t = i (\hbar/m) \nabla^2 \psi \) equation amounts to writing something like this: \( a + i b = i (c + i d) \). Remembering that \( i^2 = -1 \), you can then easily figure out that \( i(c + i d) = i c + i^2 d = -d + i c \). The \( \partial \psi / \partial t = i (\hbar/m) \nabla^2 \psi \) wave equation therefore corresponds to the following set of equations:

- \( \text{Re}(\partial \psi / \partial t) = -(\hbar/m) \cdot \text{Im}(\nabla^2 \psi) \Leftrightarrow \omega \cdot \cos(kx - \omega t) = k^2(\hbar/m) \cdot \cos(kx - \omega t) \)
- \( \text{Im}(\partial \psi / \partial t) = (\hbar/m) \cdot \text{Re}(\nabla^2 \psi) \Leftrightarrow \omega \cdot \sin(kx - \omega t) = k^2(\hbar/m) \cdot \sin(kx - \omega t) \)

It is, therefore, easy to see that \( \omega \) and \( k \) must be related through the following dispersion relation:

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my physics professor in 1979 believed Schrödinger’s equation yielded the exact solution (electron orbitals) for hydrogen.” Hence, perhaps we should not be surprised. In light of the caliber of these men, however, we are.

19 For Schrödinger’s equation in free space or the same equation with the Coulomb potential see Chapters 16 and 19 of Feynman’s Lectures on Quantum Mechanics respectively. Note that we moved the imaginary unit to the right-hand side, as a result of which the usual \( \text{minus} \) sign disappears: \( 1/i = -i \).

20 See Dirac’s description of Schrödinger’s Zitterbewegung of the electron for an explanation of the lightspeed motion of the charge. For a derivation of the \( m = 2m_{\text{eff}} \) formula, we refer the reader to our paper on the ring current model of an electron, where we write the effective mass as \( m_{\text{eff}} = m_{\gamma} \). The \( \gamma \) symbol \( (\gamma) \) refers to the photon-like character of the charge as it zips around some center at lightspeed. However, unlike a photon, a charge carries charge. Photons do not.

21 We invite the reader to double-check our calculations. If needed, we provide some more detail in one of our physics blog posts on the geometry of the wavefunction.

22 If you google this (check out the Wikipedia article on the dispersion relation, for example), you will find this relation is referred to as a non-relativistic limit of a supposedly relativistically correct dispersion relation, and the various authors of such accounts will usually also add the 1/2 factor because they conveniently (but wrongly)
\[ \omega = \frac{\hbar k^2}{m} = \frac{\hbar c^2 k^2}{E} \]

So far, so good. In fact, we can easily verify this makes sense if we substitute the energy \( E \) using the Planck-Einstein relation \( E = \hbar \cdot \omega \) and assuming the wave velocity is equal to \( c \), which should be the case if we are talking about the same vacuum as the one through which Maxwell’s electromagnetic waves are supposed to be traveling:

\[ \omega = \frac{\hbar k^2}{m} = \frac{\hbar c^2 k^2}{E} \]

We know need to think about the question we started out with: one wave or many component waves? It is fairly obvious that if we think of many component waves, each with their own frequency, then we need to think about different values \( m_i \) or \( E_i \) for the mass and/or energy of the electron as well! How can we motivate or justify this? The electron mass or energy is known, isn’t it? This is where the uncertainty comes in: the electron may have some (classical) velocity or momentum for which we may not have a definite value. If so, we may assume different values for its (kinetic) energy and/or its (linear) momentum may be possible. We then effectively get various possible values for \( m, E \) and \( p \) which we may denote as \( m_i, E_i \) and \( p_i \) respectively. We can, then, effectively write our dispersion relation and, importantly, the condition for it to make physical sense as:

\[ \omega_i = \frac{\hbar k_i^2}{m_i} = \frac{\hbar c^2 k_i^2}{E_i} = \frac{\hbar c^2 k_i^2}{\hbar \omega_i} = \frac{c^2 k_i^2}{\omega_i} \implies \frac{\omega_i^2}{k_i^2} = c^2 \implies c = f \lambda_i \]

Of course, the \( c = f \lambda \) makes a lot of sense: we would not want the properties of the medium in which matter-particles move to be different from the medium through which electromagnetic waves are travelling: lightspeed should remain lightspeed, and waves – matter-waves included – should not be traveling faster.

In the next section, we will show how one relate the uncertainties in the (kinetic) energy and the (linear) momentum of our particle using the relativistically correct energy-momentum relation and also taking into account that linear momentum is a vector and, hence, we may have uncertainty in both its direction as well as its magnitude. Such explanations also provide for a geometric interpretation of the de Broglie wavelength. At this point, however, we should just note the key conclusions from our analysis so far:

1. If there is a matter-wave, then it must travel at the speed of light and not, as Louis de Broglie suggests, at some superluminal velocity.

2. If the matter-wave is a wave packet rather than a single wave with a precisely defined frequency and wavelength, then such wave packet will represent our limited knowledge about the momentum and/or the velocity of the electron. The uncertainty is, therefore, not inherent to Nature, but to our limited knowledge about the initial conditions.

We apologize if this sounds slightly ironic but we are actually astonished Louis de Broglie does not mind having to assume superluminal speeds for wave velocities, even if it is for phase rather than group velocities.

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23 We apologize if this sounds slightly ironic but we are actually astonished Louis de Broglie does not mind having to assume superluminal speeds for wave velocities, even if it is for phase rather than group velocities.
Let us now look at a moving electron in more detail.

The wavefunction and (special) relativity

Let us consider the idea of a particle traveling in the positive \( x \)-direction at constant speed \( v \). This idea implies a pointlike concept of position: we think the particle will be *somewhere* at some point in time. The *somewhere* in this expression does not mean that we think the particle itself is dimensionless or pointlike: we think it is *not*. It just implies that we can associate the ring current with some center of the oscillation. The oscillation itself has a physical radius, which we referred to as the Compton radius of the electron and which illustrates the quantization of space that results from the Planck-Einstein relation.

Two extreme situations may be envisaged: \( v = 0 \) or \( v = c \). However, let us consider the more general case inbetween. In our reference frame, we will have a position – a mathematical point in space, that is a function of time: \( x(t) = \nu \cdot t \). Let us now denote the position and time in the reference frame of the particle itself by \( x' \) and \( t' \). Of course, the position of the particle in its own reference frame will be equal to \( x'(t') = 0 \) for all \( t' \), and the position and time in the two reference frames will be related by Lorentz’s equations:

\[
x' = \frac{x - vt}{\sqrt{1 - \frac{v^2}{c^2}}} = \frac{vt - vt}{\sqrt{1 - \frac{v^2}{c^2}}} = 0
\]

\[
t' = \frac{t - \frac{vx}{c^2}}{\sqrt{1 - \frac{v^2}{c^2}}}
\]

Hence, if we denote the energy and the momentum of the electron in our reference frame as \( E \) and \( p = \gamma m_0 \nu \), then the argument of the (elementary) wavefunction \( a \cdot e^{i \theta} \) can be re-written as follows:

\[
\theta = \frac{1}{\hbar} (E_0 t - px) = \frac{1}{\hbar} \left( \frac{E_0}{\sqrt{1 - \frac{v^2}{c^2}}} t - \frac{E_0 \nu}{c^2} x \right) = \frac{1}{\hbar} E_0 \left( \frac{t}{\sqrt{1 - \frac{v^2}{c^2}}} - \frac{\frac{vx}{c^2}}{\sqrt{1 - \frac{v^2}{c^2}}} \right) = \frac{E_0}{\hbar} t'
\]

We have just shown that the argument of the wavefunction is relativistically invariant: \( E_0 \) is, obviously, the rest energy and, because \( p' = 0 \) in the reference frame of the electron, the argument of the wavefunction effectively reduces to \( E_0 t'/\hbar \) in the reference frame of the electron itself.

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24 We conveniently choose our \( x \)-axis so it coincides with the direction of travel. This does not have any impact on the generality of the argument.

25 We may, of course, also think of it as a position vector by relating this point to the chosen origin of the reference frame: a point can, effectively, only be defined in terms of other points.

26 These are the Lorentz equations in their simplest form. We may refer the reader to any textbook here but, as usual, we like Feynman’s lecture on it (chapters 15, 16 and 17 of the first volume of Feynman’s *Lectures on Physics*).

27 One can use either the general \( E = mc^2 \) or – if we would want to make it look somewhat fancier – the \( pc = E\nu/c \) relation. The reader can verify they amount to the same.
Note that, in the process, we also demonstrated the relativistic invariance of the Planck-Einstein relation! This is why we feel that the argument of the wavefunction (and the wavefunction itself) is more real – in a physical sense – than the various wave equations (Schrödinger, Dirac, or Klein-Gordon) for which it is some solution. Let us further explore this by trying to think of the physical meaning of the de Broglie wavelength $\lambda = h/p$. How should we think of it? What does it represent?

We have been interpreting the wavefunction as an implicit function: for each $x$, we have a $t$, and vice versa. There is, in other words, no uncertainty here: we think of our particle as being somewhere at any point in time, and the relation between the two is given by $x(t) = v \cdot t$. We will get some linear motion. If we look at the $\psi = a \cdot \cos(p \cdot x/h - E \cdot t/h) + i \cdot a \cdot \sin(p \cdot x/h - E \cdot t/h)$ once more, we can write $p \cdot x/h$ as $\Delta$ and think of it as a phase factor. We will, of course, be interested to know for what $x$ this phase factor $\Delta = p \cdot x/h = 2\pi$. Hence, we write:

$$\Delta = p \cdot x/h = 2\pi \iff x = 2\pi \cdot h/p = h/p = \lambda$$

What is it this $\lambda$? If we think of our Zitterbewegung traveling in a space, we may think of an image as the one below, and it is tempting to think the de Broglie wavelength must be the distance between the crests (or the troughs) of the wave.\(^\text{28}\)

\[\text{Figure 2: An interpretation of the de Broglie wavelength?}\]

However, that would be too easy: note that for $p = mv = 0$ (or $v \to 0$), we have a division by zero and we, therefore, get an infinite value for $\lambda = h/p$. We can also easily see that for $v \to c$, we get a $\lambda$ that is equal to the Compton wavelength $h/mc$. How should we interpret that? We may get some idea by playing some more with the relativistically correct equation for the argument of the wavefunction. Let us, for example, re-write the argument of the wavefunction as a function of time only:

$$\theta = \frac{1}{\hbar} (E_0 t - px) = \frac{1}{\hbar} \frac{E_0}{\sqrt{1 - \frac{v^2}{c^2}}} (t - \frac{v}{c^2} vt) = \frac{1}{\hbar} \frac{E_0}{\sqrt{1 - \frac{v^2}{c^2}}} \left(1 - \frac{v^2}{c^2}\right) t = \sqrt{1 - \frac{v^2}{c^2}} \cdot \frac{E_0}{\hbar} t$$

We recognize the inverse Lorentz factor here, which goes from 1 to 0 as $v$ goes from 0 to $c$, as shown below.

\(^{28}\) We have an oscillation in two dimensions here. Hence, we cannot really talk about crests or troughs, but the reader will get the general idea. We should also note that we should probably not think of the plane of oscillation as being perpendicular to the plane of motion: we think it is moving about in space itself as a result of past interactions or events (think of photons scattering of it, for example).
Figure 3: The inverse Lorentz factor as a function of (relative) velocity ($v/c$)

Note the shape of the function: it is a simple circular arc. This result should not surprise us, of course, as we also get it from the Lorentz formula:

$$t' = \frac{t - \frac{v^2}{c^2}x}{\sqrt{1 - \frac{v^2}{c^2}}} = \frac{t - \frac{v^2}{c^2}t}{\sqrt{1 - \frac{v^2}{c^2}}} = \sqrt{1 - \frac{v^2}{c^2}} \cdot t$$

This formula gives us the relation between the coordinate time and proper time which – by taking the derivative of one to the other – we can write in terms of the Lorentz factor:

$$\gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}} = \frac{1}{\sqrt{1 - \beta^2}} = \frac{dt}{d\tau}$$

We introduced a different symbol here: the time in our reference frame (t) is the coordinate time, and the time in the reference frame of the object itself (τ) is referred to as the proper time. Of course, τ is just $t'$, so why are we doing this? What does it all mean? We need to do these gymnastic because we want to introduce a not-so-intuitive but very important result: the Compton radius becomes a wavelength when v goes to c.\(^{29}\)

We will be very explicit here and go through a simple numerical example to think through that formula above. Let us assume that, for example, that we are able to speed up an electron to, say, about one tenth of the speed of light. Hence, the Lorentz factor will then be equal to $\gamma = 1.005$. This means we added 0.5% (about 2,500 eV) – to the rest energy $E_0$: $E_v = \gamma E_0 = 1.005 \cdot 0.511$ MeV $\approx 0.5135$ MeV. The relativistic momentum will then be equal to $m_v = (0.5135 \text{ eV}/c^2) \cdot (0.1 \cdot c) = 5.135$ eV/c. We get:

$$0 = \frac{E_0}{\hbar} t' = \frac{1}{\hbar} (E_v t - px) = \frac{1}{\hbar} \left( \frac{E_0}{\sqrt{1 - \frac{v^2}{c^2}}} t - \frac{E_0}{c^2 \sqrt{1 - \frac{v^2}{c^2}}} x \right) = 0.955 \frac{E_0}{\hbar} t$$

This is interesting because we can see these equations are not all that abstract: we effectively get an explanation for relativistic time dilation out of them. An equally interesting question is this: what

\(^{29}\)To be precise, the Compton radius multiplied by $2\pi$ becomes a wavelength, so we are talking the Compton circumference, or whatever you want to call it.
happens to the \textit{radius} of the oscillation for larger (classical) velocity of our particle? Does it change? It must. In the moving reference frame, we measure higher mass and, therefore, higher energy – as it includes the kinetic energy. The \( c^2 = a^2 \cdot \omega^2 \) identity must now be written as \( c^2 = a'^2 \cdot \omega'^2 \). Instead of the rest mass \( m_0 \) and rest energy \( E_0 \), we must now use \( m_v = \gamma m_0 \) and \( E_v = \gamma E_0 \) in the formulas for the Compton radius and the Einstein-Planck frequency\textsuperscript{30}, which we just write as \( m \) and \( E \) in the formula below:

\[
ma'^2\omega'^2 = m \frac{\hbar^2}{m^2 c^2} \frac{m^2 c^4}{\hbar^2} = mc^2
\]

This is easy to understand intuitively: we have the mass factor in the denominator of the formula for the Compton radius, so it must increase as the mass of our particle increases with speed. Conversely, the mass factor is present in the numerator of the \( zbw \) frequency, and this frequency must, therefore, increase with velocity. It is interesting to note that we have a simple (inverse) proportionality relation here. The idea is visualized in the illustration below\textsuperscript{31}: the \textit{radius} of the circulatory motion must effectively diminish as the electron gains speed.\textsuperscript{32}

**Figure 4:** The Compton radius must decrease with increasing velocity

Can the velocity go to \( c \)? In the limit, yes. This is very interesting because we can see that the circumference of the oscillation effectively turns into a \textit{linear} wavelength in the process!\textsuperscript{33} This rather remarkable geometric property related our \( zbw \) electron model with our photon model, which we will not talk about here, however.\textsuperscript{34} Let us quickly make some summary remarks here, before we proceed to what we wanted to present here: a geometric interpretation of the \textit{de Broglie} wavelength:

\textsuperscript{30} Again, the reader should note that both the formula for the Compton radius or wavelength as well as the Planck-Einstein relation are relativistically invariant.

\textsuperscript{31} We thank Prof. Dr. Giorgio Vassallo and his publisher to let us re-use this diagram. It originally appeared in an article by Francesco Celani, Giorgio Vassallo and Antonino Di Tommaso (\textit{Maxwell’s equations and Occam’s Razor}, November 2017).

\textsuperscript{32} Once again, however, we should warn the reader that he or she should imagine the plane of oscillation to rotate or oscillate itself. He should not think of it of being \textit{static} – unless we think of the electron moving in a magnetic field, in which case we should probably think of the plane of oscillation as being \textit{parallel} to the direction of propagation. We will let the reader think through the geometric implications of this.

\textsuperscript{33} We may, therefore, think of the Compton wavelength as a \textit{circular} wavelength: it is the length of a circumference rather than a linear feature!

\textsuperscript{34} We may refer the reader to our paper on \textit{Relativity, Light and Photons}. 

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1. The center of the *Zitterbewegung* was plain nothingness and we must, therefore, assume some two-dimensional oscillation makes the charge go round and round. This is, in fact, the biggest mystery of the model and we will, therefore, come back to it later. As for now, the reader should just note that the angular frequency of the *Zitterbewegung* rotation is given by the Planck-Einstein relation (\(\omega = E/\hbar\)) and that we get the *Zitterbewegung* radius (which is just the Compton radius \(a = r_C = \hbar/mc\)) by equating the \(E = mc^2\) and \(E = ma^2\omega^2\) equations. The energy and, therefore, the (equivalent) mass is in the oscillation and we, therefore, should associate the momentum \(p = E/c\) with the electron as a whole or, if we would really like to associate it with a single mathematical point in space, with the center of the oscillation – as opposed to the rotating massless charge.

2. We should note that the distinction between the pointlike *charge* and the electron is subtle but essential. The electron is the *Zitterbewegung* as a whole: the pointlike charge has no rest mass, but the electron as a whole does. In fact, that is the whole point of the whole exercise: we explain the rest mass of an electron by introducing a *rest matter oscillation*.

3. As Dirac duly notes, the model cannot be verified *directly* because of the extreme frequency (\(f_e = \omega_e/2\pi = E/\hbar \approx 0.123 \times 10^{-21}\) Hz) and the sub-atomic scale (\(a = r_C = \hbar/mc = 386 \times 10^{-15}\) m). However, it can be verified *indirectly* by phenomena such as Compton scattering, the interference of an electron with itself as it goes through a one- or double-slit experiment, and other indirect evidence. In addition, it is *logically* consistent as it generates the right values for the angular momentum (\(L = \hbar/2\)), the magnetic moment (\(\mu = (q_e/2m)\cdot\hbar\)), and other intrinsic properties of the electron. The model cannot be verified *directly* because of the extreme frequency (\(f_e = \omega_e/2\pi = E/\hbar \approx 0.123 \times 10^{-21}\) Hz) and the sub-atomic scale (\(a = r_C = \hbar/mc = 386 \times 10^{-15}\) m). However, it can be verified *indirectly* by phenomena such as Compton scattering, the interference of an electron with itself as it goes through a one- or double-slit experiment, and other indirect evidence. In addition, it is *logically* consistent as it generates the right values for the angular momentum (\(L = \hbar/2\)), the magnetic moment (\(\mu = (q_e/2m)\cdot\hbar\)), and other intrinsic properties of the electron.35

We are now ready to finally give you a geometric interpretation of the *de Broglie* wavelength.

**The geometric interpretation of the *de Broglie* wavelength**

We should refer the reader to Figure 4 to ensure an understanding of what happens when we think of an electron in motion. If the tangential velocity remains equal to \(c\), and the pointlike charge has to cover some horizontal distance as well, then the circumference of its rotational motion must decrease so it can cover the extra distance. Our formula for the *zbw* or Compton radius was this:

\[
a = \frac{\hbar}{mc} = \frac{\lambda_C}{2\pi}
\]

The \(\lambda_C\) is the *Compton* wavelength. We may think of it as a *circular* rather than a *linear* length: it is the circumference of the circular motion.36 How can it decrease? If the electron moves, it will have some kinetic energy, which we must add to the *rest energy*. Hence, the mass \(m\) in the denominator (\(mc\)) increases and, because \(\hbar\) and \(c\) are physical constants, \(a\) must decrease.37 How does that work with the...

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35 The two results that we gave also show we get the gyromagnetic factor (\(g = 2\)). We have also demonstrated that we can easily explain the anomaly in the magnetic moment of the electron by assuming a non-zero physical dimension for the pointlike charge (see our paper on *The Electron and Its Wavefunction*).

36 Hence, the \(C\) subscript stands for the \(C\) of Compton, not for the speed of light (\(c\)).

37 We advise the reader to always think about proportional and inversely proportional relations (\(y = kx\) versus \(y = x/k\)) throughout the *exposé* because these relations are *not* always intuitive. The *inverse* proportionality relation between the Compton radius and the mass of a particle is a case in point in this regard: a more *massive* particle has a *smaller* size!
frequency? The frequency is proportional to the energy \( E = \hbar \omega = \hbar f = \hbar / T \) so the frequency – in whatever way you want to measure it – must also increase. The cycle time \( T \) must, therefore, decrease.

We write:

\[
\theta = \omega t = \frac{E}{\hbar} t = \frac{\gamma E_0}{\hbar} t = 2\pi \cdot \frac{t}{T}
\]

Hence, our Archimedes’ screw gets stretched, so to speak. Let us think about what happens here. We get the following formula for the \( \lambda \) wavelength in Figure 2:

\[
\lambda = \nu \cdot T = \frac{\nu}{f} = \frac{\nu \cdot h}{E} = \nu \cdot \frac{h}{mc^2} = \frac{\nu \cdot h}{c^2} = \beta \cdot \lambda_c
\]

It is now easy to see that, if we let the velocity go to \( c \), the circumference of the oscillation will effectively become a linear wavelength. We can now relate this classical velocity \( (v) \) to the equally classical linear momentum of our particle and provide a geometric interpretation of the de Broglie wavelength, which we’ll denote by using a separate subscript: \( \lambda_p = h/p \). It is, obviously, different from the \( \lambda \) wavelength in Figure 2. In fact, we have three different wavelengths now: the Compton wavelength \( \lambda_c \) (which is a circumference, actually), that weird horizontal distance \( \lambda \), and the de Broglie wavelength \( \lambda_p \). It is easy to make sense of them by relating all three. Let us first re-write the de Broglie wavelength in terms of the Compton wavelength \( (\lambda_c = h/mc) \), its (relative) velocity \( \beta = \nu/c \), and the Lorentz factor \( \gamma \):

\[
\lambda_p = \frac{h}{p} = \frac{h}{mv} = \frac{hc}{mc^2} = \frac{h}{mc} \beta = \frac{\lambda_c}{\beta} = \frac{1}{\gamma \beta} \frac{h}{m_0 c} = \frac{1}{\gamma \beta} 2\pi a_0
\]

It is a curious function, but it helps us to see what happens to the de Broglie wavelength as \( m \) and \( v \) both increase as our electron picks up some momentum \( p = m \cdot v \). Its wavelength must actually decrease as its (linear) momentum goes from zero to some much larger value – possibly infinity as \( v \) goes to \( c \) – and the \( 1/\gamma \beta \) factor tells us how exactly. To help the reader, we inserted a simple graph (below) that shows how the \( 1/\gamma \beta \) factor comes down from infinity (+\( \infty \)) to zero as \( v \) goes from 0 to \( c \) or – what amounts to the same – if the relative velocity \( \beta = \nu/c \) goes from 0 to 1. The \( 1/\gamma \) factor – so that is the inverse Lorentz factor – is just a simple circular arc, while the \( 1/\beta \) function is just a regular inverse function \( (y = 1/x) \) over the domain \( \beta = \nu/c \), which goes from 0 to 1 as \( v \) goes from 0 to \( c \). Their product gives us the green curve which – as mentioned – comes down from +\( \infty \) to 0.

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38 This is why we think of the Compton wavelength as a circular wavelength. However, note that the idea of rotation does not disappear: it is what gives the electron angular momentum—regardless of its (linear) velocity! As mentioned above, these rather remarkable geometric properties relate our zbw electron model with our photon model, which we have detailed in another paper.
Figure 5: The $1/\gamma$, $1/\beta$ and $1/\gamma\beta$ graphs

This analysis yields the following:

1. The *de Broglie* wavelength will be equal to $\lambda_c = \hbar/mc$ for $v = c$:

$$\lambda_p = \frac{h}{p} = \frac{h}{mc} \cdot \frac{1}{\beta} = \lambda_c = \frac{h}{mc} \iff \beta = 1 \iff v = c$$

2. We can now relate both Compton as well as *de Broglie* wavelengths to our new wavelength $\lambda = \beta \cdot \lambda_c$ wavelength—which is that length between the crests or troughs of the wave.

We get the following two rather remarkable results:

$$\lambda_p \cdot \lambda = \lambda_p \cdot \beta \cdot \lambda_c = \frac{1}{\beta} \cdot \frac{h}{mc} \cdot \beta \cdot \frac{h}{mc} = \lambda_c^2$$

$$\frac{\lambda}{\lambda_p} = \frac{\beta \cdot \lambda_c}{\lambda_p} = p \cdot \frac{v}{c} \cdot \frac{h}{mc} = \frac{mv^2}{mc^2} = \beta^2$$

The product of the $\lambda = \beta \cdot \lambda_c$ wavelength and *de Broglie* wavelength is the square of the Compton wavelength, and their ratio is the square of the relative velocity $\beta = v/c$. – always! – and their ratio is equal to $1 – always!$

This is all very interesting but *not good enough* yet: the formulas do *not* give us the easy geometric interpretation of the *de Broglie* wavelength that we are looking for. We get such easy geometric interpretation only when using natural units.

If we re-define our distance, time and force units by equating $c$ and $h$ to 1, then the Compton wavelength (remember: it is a circumference, really) and the mass of our electron will have a simple inversely proportional relation:

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39 We should emphasize, once again, that our two-dimensional wave has no real crests or troughs: $\lambda$ is just the distance between two points whose argument is the same—except for a phase factor equal to $n \cdot 2\pi$ ($n = 1, 2, ...$).

40 Equating $c$ to 1 gives us natural distance and time units, and equating $h$ to 1 then also gives us a natural force unit—and, because of Newton’s law, a natural mass unit as well. Why? Because Newton’s $F = m \cdot a$ equation is relativistically correct: a force is that what gives some mass acceleration. Conversely, mass can be defined of the inertia to a change of its state of motion—because any change in motion involves a force and some acceleration. We, therefore, prefer to write $m$ as the proportionality factor between the force and the acceleration: $m = F/a$. This explains why time, distance and force units are closely related.
\[ \lambda_C = \frac{1}{\gamma m_0} = \frac{1}{m} \]

We get equally simple formulas for the de Broglie wavelength and our \( \lambda \) wavelength:

\[ \lambda_p = \frac{1}{\beta \gamma m_0} = \frac{1}{\beta m} \]

\[ \lambda = \beta \cdot \lambda_C = \frac{\beta}{\gamma m_0} = \frac{\beta}{m} \]

This is quite deep: we have three lengths here – defining all of the geometry of the model – and they all depend on the rest mass of our object and its relative velocity only. They are related through that equation we found above:

\[ \lambda_p \cdot \lambda = \lambda_C^2 = \frac{1}{m^2} \]

This is nothing but the latus rectum formula for an ellipse, which is illustrated below.\(^{41}\) The length of the chord – perpendicular to the major axis of an ellipse is referred to as the latus rectum. One half of that length is the actual radius of curvature of the osculating circles at the endpoints of the major axis.\(^{42}\) We then have the usual distances along the major and minor axis (\( a \) and \( b \)). Now, one can show that the following formula has to be true:

\[ a \cdot p = b^2 \]

![Figure 6: The latus rectum formula](image)

The reader can now easily verify that our three wavelengths obey the same latus rectum formula, which we think of as a rather remarkable result. We must now proceed and offer some final remarks on a far more difficult question.

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\(^{41}\) Source: Wikimedia Commons (By Ag2gaeh - Own work, CC BY-SA 4.0, [https://commons.wikimedia.org/w/index.php?curid=57428275](https://commons.wikimedia.org/w/index.php?curid=57428275)).

\(^{42}\) The endpoints are also known as the vertices of the ellipse. As for the concept of an osculating circle, that is the circle which, among all tangent circles at the given point, which approaches the curve most tightly. It was named circulus osculans – which is Latin for ‘kissing circle’ – by Gottfried Wilhelm Leibniz. Apart from being a polymath and a philosopher, he was also a great mathematician. In fact, he may be credited for inventing differential and integral calculus.
What is the nature of the centripetal force?

We think we have sufficiently demonstrated the theoretical attractiveness of the historical ring current model. This is why we shared it as widely as we could. We usually get positive comments. However, when we first submitted our thoughts to Prof. Dr. Alexander Burinskii, who is leading the research on possible Dirac-Kerr-Newman geometries of electrons\textsuperscript{43}, he wrote us this:

“I know many people who considered the electron as a toroidal photon\textsuperscript{44} and do it up to now. I also started from this model about 1969 and published an article in JETP in 1974 on it: “Microgeons with spin”. [However] There was [also] this key problem: what keeps [the pointlike charge] in its circular orbit?”

This question still puzzled us, and we do not have a definite answer to it. It is, in fact, the only remaining mystery in quantum physics. What can we say about it? Let us elaborate Burinskii’s point:

1. The centripetal force must, obviously, be electromagnetic because it only has a pointlike charge to grab onto, and comparisons with superconducting persistent currents are routinely made. However, such comparisons do not answer this pertinent question: in free space, there is nothing to effectively hold the pointlike charge in place and it must, therefore, spin away!

2. In addition, the analogy with superconducting persistent currents also does not give any unique Compton radius: the formulas work for any mass. It works, for example, for a muon-electron, and for a proton. The question then becomes: what makes an electron an electron—and what makes a muon a muon? Or a proton a proton?

For the time being, we must simply accept an electron is what it is. In other words, we must take both the elementary charge and the mass of the electron (and its more massive variant(s), as well as the mass of the proton) as given by Nature. In the longer run, however, we should probably abandon an explanation of the ring current model in terms of Maxwell’s equations in favor for what it is, effectively, the more mysterious explanation of a two-dimensional oscillation. However, we have not advanced very far in our thinking on these issues\textsuperscript{45} and we, therefore, welcome any suggestion that the reader of this paper might have in this regard.

END

\textsuperscript{43} We will let the reader \textit{google} the relevant literature on electron models based on Dirac-Kerr-Newman geometries here. The mentioned email exchange between Dr. Burinskii and the author of this paper goes back to 22 December 2018.

\textsuperscript{44} This was Dr. Burinskii’s terminology at the time. It does refer to the \textit{Zitterbewegung} electron: a pointlike charge with no mass in an oscillatory motion—orbiting at the speed of light around some center. Dr. Burinskii later wrote saying he does not like to refer to the pointlike charge as a toroidal \textit{photon} because a photon does not carry any charge. The pointlike charge inside of an electron does: that is why matter-particles are matter-particles and photons are photons. Matter-particles carry charge (we think of neutrons as carrying equal positive and negative charge).

\textsuperscript{45} For some speculative thoughts, we may refer the reader to previous references, such as our \textit{electron paper} or the Annex to our more general \textit{paper on classical quantum physics}. We calculate the rather enormous force inside of muon and proton in these papers and conclude they may justify the concept of a strong(er) force.